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An effective formulation for estimating wetland surface energy fluxes from weather data

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Abstract: In modelling evapotranspiration, the need for land surface variables including ground heat fluxes (G), surface temperature (T_s), surface relative humidity (RH_s) and surface resistance often present a challenge due to land heterogeneity and limited measurements. This study introduces a simple formulation rooted in the shared physical basis of the maximum entropy model (MaxEnt), the Relative Humidity at Equilibrium (ETRHEQ) method, and the Surface Flux Equilibrium (SFE) method, and it estimates sensible (H) and latent fluxes (LE) in wetlands without requiring land surface variables or site-specific calibration, except for an assumed vegetation height. Further, it effectively estimates LE from half-hourly to monthly scales in FLUXNET and AmeriFlux wetland sites. While its performance in estimating H is less satisfactory due to loosely constrained boundary conditions, it shows promising potential for simultaneously and precisely estimating LE , H , G , T_s , and RH_s from weather data in various ecosystems.

Key points:

1. The formulation is based on the principle of maximum Shannon information entropy production for turbulence fluxes.

22 2. The formulation does not require land surface variables or site-specific calibration; only an
23 assumed vegetation height is needed.

24 3. The formulation effectively estimates LE from half-hourly to monthly scales.

25 **Plain language summary:** This study introduces a new method to predict how much water and
26 heat wetlands transport to the atmosphere, a process that is usually complicated because it
27 involves a lot of detailed information about land properties that are hard to measure. This new
28 method does not need all those details, and instead just needs an estimate of how tall the plant
29 canopy is. This method works extremely well for predicting water release into the air over
30 periods ranging from half-hourly to monthly in FLUXNET and AmeriFlux wetland sites.
31 Although this method is not perfect at predicting heat release due to some assumptions that have
32 to be made about ground heat and surface temperature, it shows a lot of promise. With a bit of
33 fine-tuning, it could be used to accurately measure both water and heat exchanges in various
34 types of ecosystems, not just wetlands.

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37 **1. Introduction**

38 The partitioning of energy on the land surface of terrestrial ecosystems into ground heat (G),
39 sensible heat (H) and latent heat (LE) has long been recognized as a result of complex
40 interactions between atmospheric and land surface properties (Duveiller et al., 2018; Forzieri et
41 al., 2020; Williams and Torn, 2015; Wilson et al., 2002). At short temporal scales, it is impacted
42 by plant physiological activities and boundary layer properties, and over the long term, the
43 biogeochemical cycling, disturbance, and climate all have significant roles to play (Arneeth et al.,
44 2012; Green et al., 2017; Wilson et al., 2002). While the importance of land surface properties

45 cannot be overlooked, land surface variables are a challenge to parameterize due to land
46 heterogeneity and varied physiological responses of vegetation to changing environmental
47 conditions (Dickinson et al., 1991; Mueller and Seneviratne, 2014; Wang and Dickinson, 2012).

48 Recent studies proposed two methodologies, namely the Relative Humidity at Equilibrium
49 (ETRHEQ) method and the Surface Flux Equilibrium (SFE) method, to estimate surface energy
50 fluxes from near-surface atmospheric conditions (McColl et al., 2019; Salvucci and Gentine,
51 2013). ETRHEQ determines the optimal daily surface conductance that yields the most accurate
52 ET predictions based on minimum vertical variance of relative humidity (RH) (Salvucci and
53 Gentine, 2013), and SFE provides the solution of ETRHEQ at the steady state (McColl et al.,
54 2019). The two methods are justified by strong land-atmospheric coupling wherein land surface
55 properties are embedded in the near-surface atmospheric conditions (McColl and Rigden, 2020;
56 McColl et al., 2019). Conversely, the conditions of the near-surface atmosphere are also reflected
57 in land surface variables, which partly justifies another methodology called the maximum
58 entropy model (MaxEnt) that estimates surface energy fluxes using only the surface temperature
59 and surface relative humidity in addition to net radiation (Wang and Bras, 2011; Wang and Bras,
60 2009). Although the three models have shown success over a variety of ecosystems worldwide
61 (Chen et al., 2021; McColl and Rigden, 2020; Rigden and Salvucci, 2015; Yang et al., 2022),
62 each have their own limitations. ETRHEQ requires vegetation height and ground heat fluxes in
63 addition to 24-hour subdaily weather measurements, to estimate latent and sensible fluxes at the
64 daily scale (Rigden and Salvucci, 2015; Salvucci and Gentine, 2013). SFE, though it requires
65 less parameters (i.e., only net radiation, ground heat flux, air temperature and air specific
66 humidity), works for sites near or at the steady state and estimates energy fluxes at the daily or
67 larger temporal scales (Chen et al., 2021; Kim et al., 2023; McColl and Rigden, 2020). The

68 MaxEnt model is formulated based on minimizing the dissipation function of turbulent fluxes
 69 (which is equivalent to maximizing Shannon information entropy production of the turbulent
 70 fluxes (Dewar, 2005)) and the Monin-Obukhov similarity theory (MOST)'s extremum solution
 71 (Wang and Bras, 2009), but the justification of extremum solution still requires further
 72 examination (Wang and Bras, 2010; Wang et al., 2023).

73 Wang et al. (2023) investigated the linkage of the three models and found that minimizing the
 74 dissipation function of energy fluxes in MaxEnt is equivalent to minimizing the vertical variance
 75 of RH in ETRHEQ. The empirical success of the three models is explained by the fact that far-
 76 from-equilibrium ecosystems progress toward a steady state (i.e., the SFE state) by minimizing
 77 dissipation, and this tendency is manifested through the vertical variance of RH (Wang et al.,
 78 2023). In addition, Wang et al. (2023) demonstrated that the connection among the three models
 79 is independent of Monin-Obukhov similarity theory (MOST)'s extremum solution (Wang et al.,
 80 2023), and proposed a more general formulation describing the dissipation function (D) of
 81 energy fluxes for both non-steady and steady states, as:

$$82 \quad D = \frac{2G^2}{I_s} + \frac{2H^2}{I_a} + \frac{LE^2}{I_e} \quad (1)$$

$$83 \quad \text{with } I_s = \sqrt{I_d^2 + \theta I_w^2}, I_a = \rho c_p \sqrt{g_a}, \text{ and } I_e = \frac{\delta}{\gamma} RH_s I_a.$$

84 where I_s , I_a and I_e are the thermal inertia parameters for G, H and LE, respectively; The
 85 parameterization of I_s is provided in Huang et al. (2017) and Yang et al. (2022) in which I_d is the
 86 thermal inertia of dry soil; θ is the volumetric soil moisture; I_w is the thermal inertia of still
 87 liquid water; ρ is the density of air; c_p is the specific heat capacity of air; g_a is the aerodynamic
 88 conductance; δ is the slope of the relation between saturated specific humidity and temperature,

89 $\gamma = \frac{c_p}{\lambda}$ with λ being the latent heat of vaporization of water; and RH_s is the surface relative
90 humidity. The detailed formulation will be introduced in the next section.

91 The new formulation is denoted as MaxEnt-ETRHEQ, indicating the shared physical basis
92 underlying MaxEnt and ETRHEQ. It appears to require both atmospheric and land surface
93 variables at first glance. However, closer scrutiny revealed that land surface variables such as
94 surface temperature, surface relative humidity and soil moisture are interlinked in the calculation
95 of G, H and LE under energy closure. This interconnection renders the formulation self-
96 constrained. Consequently, the energy fluxes and the land surface variables can be analytically
97 determined by identifying the minimum value of D given suitable ranges of surface temperature
98 and relative humidity. Therefore, MaxEnt-ETRHEQ has potential to estimate surface energy
99 fluxes for various ecosystems, with minimal or no land surface information. But its effectiveness
100 is yet to be examined. Leveraging our proficiency and background in wetland ecosystems, we
101 demonstrate in this paper that MaxEnt-ETRHEQ is an effective formulation for estimating
102 energy fluxes for wetland ecosystems, especially for estimating LE from subdaily to monthly
103 scales, and it does not necessitates any land surface parameters; only an assumption regarding
104 vegetation height is required.

105

106 **2. Methods**

107 2.1 The formulation of MaxEnt-ETRHEQ

108 The main formula of MaxEnt-ETRHEQ is given as Eq. 1. The required input parameters are
109 atmospheric pressure (p), air temperature (T_a), wind speed (u), friction velocity (u^*), air relative
110 humidity (RH), net radiation (R_n), the height of the measurements of weather data (z) and

111 vegetation height (z_{veg}). Meanwhile, MaxEnt-ETRHEQ will automatically create two variables,
 112 surface temperature (T_s) and surface relative humidity (RH_s) within a pre-defined range for
 113 studied ecosystems (will be explained later).

114 The surface pressure (p_s) is calculated from the atmospheric pressure by rearranging the formulas
 115 used in ETRHEQ, as (Salvucci and Gentine, 2013):

$$116 \quad p_s = \frac{p}{\exp\left(\frac{-gz}{R_d T_a}\right)} \quad (2)$$

117 where p_s is the surface pressure (Pa), p is the atmospheric pressure (Pa), g is the gravitational
 118 constant ($9.8 \text{ m}\cdot\text{s}^{-2}$), z is the height of the measurements of weather data (m), R_d is the gas
 119 constant for dry air ($287 \text{ J}\cdot\text{kg}^{-1}\cdot\text{K}^{-1}$), and T_a is the air temperature (K).

120 Saturation vapor pressure (e^*) is calculated from integrated Clausius–Clapeyron relation, as
 121 (Salvucci and Gentine, 2013):

$$122 \quad e^*(T_a) = 611.2 \times \exp\left(\frac{17.67 \times (T_a - 273.15)}{T_a - 29.65}\right) \quad (3)$$

$$123 \quad e^*(T_s) = 611.2 \times \exp\left(\frac{17.67 \times (T_s - 273.15)}{T_s - 29.65}\right) \quad (4)$$

124 where $e^*(T_a)$ and $e^*(T_s)$ are saturation vapor pressure (Pa) at air temperature (T_a , K) and surface
 125 temperature (T_s , K), respectively.

126 Saturated specific humidity (q^*) is related to the saturation vapor pressure (e^*) through the
 127 following equations (Salvucci and Gentine, 2013):

$$128 \quad q^*(T_a) = \frac{\epsilon e^*(T_a)}{p - (1 - \epsilon)e^*(T_a)} \quad (5)$$

$$129 \quad q^*(T_s) = \frac{\epsilon e^*(T_s)}{p_s - (1 - \epsilon)e^*(T_s)} \quad (6)$$

130 where ϵ is the dimensionless ratio of the gas constant for dry air to water vapor (0.622).

131 Using Eq. 5 and 6, the slope of the relation between saturated specific humidity (q^*) and
132 temperature (T) can be linearly extrapolated following (Kim et al., 2021; McColl et al., 2019):

$$133 \quad \delta = \frac{q^*(T_s) - q^*(T_a)}{T_s - T_a} \quad (7)$$

134 where $q^*(T_s)$ and $q^*(T_a)$ are the surface and atmospheric saturated specific humidity ($\text{kg}\cdot\text{kg}^{-1}$),
135 respectively, and T_s and T_a are the surface and air temperature (K), respectively.

136 The sensible and latent heat fluxes are calculated using the flux gradient equations, as (Kim et
137 al., 2021):

$$138 \quad H = \rho c_p g_a (T_s - T_a) \quad (8)$$

$$139 \quad LE = \lambda \rho g_a (q_s - q_a) \quad (9)$$

140 where H and LE are the sensible and latent heats ($\text{W}\cdot\text{m}^{-2}$), ρ is the density of air ($\rho = \frac{p}{R_d T_a}$,
141 $\text{kg}\cdot\text{m}^{-3}$), c_p is the specific heat of air at constant pressure ($1004.7 \text{ J}\cdot\text{kg}^{-1}\cdot\text{C}^{-1}$), g_a is the
142 aerodynamic conductance accounting for atmospheric stability ($\text{m}\cdot\text{s}^{-1}$), λ is the latent heat of
143 vaporization ($2.502 \times 10^6 \text{ J}\cdot\text{kg}^{-1}$), q_s is the surface specific humidity ($q_s = RH_s \cdot q^*(T_s)$, $\text{kg}\cdot\text{kg}^{-1}$),
144 and q_a is the air specific humidity ($q_a = RH \cdot q^*(T_a)$, $\text{kg}\cdot\text{kg}^{-1}$).

145 The aerodynamic conductance under the neutral atmospheric condition (g_{a_n}) is given by Allen et
146 al. (1998), as:

$$147 \quad g_{a_n} = \frac{\kappa^2 u}{\ln\left(\frac{z-d}{z_{om}}\right) \ln\left(\frac{z-d}{z_{oh}}\right)} \quad (10)$$

148 with κ being the von Karman constant (0.41), u being the wind speed ($\text{m}\cdot\text{s}^{-2}$), z being the height
149 of height of the measurements of wind speed (m), d being the zero-plane displacement height

150 (m), z_{om} is the roughness length governing momentum transfer (m), and z_{oh} is the roughness
151 length governing transfer of heat and vapour (m)

152 When no vegetation is present in the study sites ($z_{veg}=0$ m), d is set as 0 m, with both z_{om} and z_{oh}
153 being set as 0.001 m; whereas in the presence of vegetation, d is set as 0.7 of z_{veg} , with z_{om} being
154 0.1 of z_{veg} , and z_{oh} being estimated using κB^{-1} approach, following Salvucci and Gentine (2013):

$$155 \quad \kappa B^{-1} = \ln\left(\frac{z_{om}}{z_{oh}}\right) \cong \kappa(6Re^{\frac{1}{4}} - 5) \quad (11)$$

156 where Re is the roughness Reynolds number ($Re = \frac{u^* z_{om}}{\nu}$, with u^* is the friction velocity ($m \cdot s^{-2}$)
157 and ν being the kinematic viscosity, as $1.45 \times 10^{-5} m^2 \cdot s^{-1}$).

158 To account for atmospheric stability, the actual aerodynamic conductance ($g_a, m \cdot s^{-1}$) is calculated
159 following Merlin et al. (2016), as:

$$160 \quad g_a = (1 + R_i)^\eta \cdot g_{a,n} \quad (12)$$

$$161 \quad R_i = \frac{\beta_{thermal} \times g z (T_s - T_a)}{T_a u^2} \quad (13)$$

162 where $\beta_{thermal}$ is the thermal expansion coefficient, and $\beta_{thermal}=5$ was used following Choudhury
163 et al. (1986) and Merlin et al. (2011); g is the gravitational constant ($9.8 m \cdot s^{-2}$), T_s is the surface
164 soil temperature (K), T_a is the air temperature (K). In Eq. (11), the coefficient η is set to 0.75 in
165 unstable conditions ($T_s > T_a$) and to 2 in stable conditions ($T_s < T_a$); u is the wind speed ($m \cdot s^{-1}$)
166 and z is the height (m) at which wind speed was measured.

167 The ground heat flux ($G, W \cdot m^{-2}$) is calculated using energy balance equation as:

$$168 \quad G = R_n - H - LE \quad (14)$$

169 where R_n is the net radiation ($\text{W}\cdot\text{m}^{-2}$), and H and LE are calculated based on Eq. 8 and 9.

170 The parameterization of thermal inertias (I_s , I_a and I_e) is provided in Eq. 1. To minimize the land
171 surface parameters needed in the MaxEnt-ETRHEQ formulation, I_s is set as a constant (1300
172 $\text{J}\cdot\text{m}^{-2}\cdot\text{K}^{-1}\cdot\text{s}^{-1/2}$, i.e., tiu) following Rigden and Salvucci (2017). It is postulated that such a
173 constant is acceptable, because: (1) Rigden and Salvucci (2015) stated that the optimal range of
174 I_s was between 300 and 1000 tiu for AmeriFlux sites, and as I_s increases with wetter soils, it
175 should be slightly higher than the optimal range; (2) Rigden and Salvucci (2017) used the
176 calibrated I_s of 1300 tiu for their study sites across united states; (3) the modelling results agree
177 well with the eddy covariance measurements (presented in the results section); and (4) using
178 measured soil moisture did not significantly improve the modelling performance (presented in
179 Table S4).

180 The last step is to specify appropriate ranges of G , T_s and RH_s . Without this specification, G
181 could become unrealistically large, which does not occur in the real world. After specifying the
182 ranges, the dissipation function D is computed for every set of input weather data and every
183 possible paring of G , T_s and RH_s . The selection of the optimal set of G , T_s and RH_s will be done
184 through finding the minimum D . Once these optimal values are found, H and LE are
185 concurrently determined through the calculations from Eq.1 to 14.

186 2.2 The boundary conditions for wetland ecosystems

187 The upper limit of G for wetland ecosystems is set as 0.20 of R_n based on the empirical
188 relationship between G and R_n used in GLEAM model: $G/R_n = 0.20$ for short vegetation (0.05
189 $\text{m} < Z_{\text{veg}} < 1 \text{ m}$) and $G/R_n = 0.15$ for tall vegetation ($Z_{\text{veg}} > 1 \text{ m}$) (Miralles et al., 2011). T_s and T_a
190 at the 2 m height above land surface may differ by several $^{\circ}\text{C}$, but the difference between
191 maximum T_a and maximum T_s may vary up to $30 \text{ }^{\circ}\text{C}$ (Good et al., 2017; Mildrexler et al., 2011),

192 so T_s is set to $T_a \pm 30$ °C. Typically, RH_s must be higher than RH for evapotranspiration to occur.
 193 As evapotranspiration progresses, RH tends to increase while RHs decreases until the ecosystem
 194 reaches the surface flux equilibrium state ($RH_{eq}=RH=RH_s$). This suggests that there exists a
 195 boundary for RH_s , which falls within the range of RH and RH_{eq} . To estimate RH_{eq} , the Priestley-
 196 Taylor equation for water body (i.e., the left of the equals sign of Eq. 15) (Priestley and Taylor,
 197 1972) is combined with the PM_{RH} equation under $RH=RH_s$ (the right of the equals sign of Eq.
 198 15) (Kim et al., 2021) to determine the maximum RH_{eq} , as:

$$199 \quad 1.26\lambda(R_n - G) \frac{\Delta}{\Delta+\gamma'} = \frac{RH_{eq}\Delta}{RH_{eq}\Delta+\gamma'} (R_n - G) \quad (15)$$

200 where 1.26 is the Priestley-Taylor coefficient for open water; λ is the latent heat of vaporization
 201 (2.502×10^6 J·kg⁻¹); R_n is the net radiation (W·m⁻²); G is the ground heat flux (W·m⁻²); Δ is the
 202 slope of the relation between saturation vapor pressure and temperature (Pa·°C⁻¹); γ' is the
 203 psychrometric constant ($\gamma' = \frac{p c_p}{\epsilon \lambda}$, with being is the air pressure (Pa), c_p being the specific heat of
 204 air at constant pressure (1004.7 J·kg⁻¹·°C⁻¹), ϵ being the dimensionless ratio of the gas constant
 205 for dry air to water vapor (0.622), and); and λ being the latent heat of vaporization (2.502×10^6
 206 J·kg⁻¹)); and RH_{eq} is the equilibrium RH of a saturated wetland ecosystem. Rearranging Eq. 15
 207 leads to the expression of RH_{eq} , as:

$$208 \quad RH_{eq} = \frac{1.26\lambda\gamma'}{\Delta(1-1.26\lambda)+\gamma'} \quad (16)$$

209 There are multiple ways to estimate Δ . In this study, the method provided in the FAO Penmen-
 210 Monteith equation is chosen to estimate Δ from T_a , as (Allen et al., 1998):

211
$$\Delta = \frac{1000 \times 4098 [0.6108 \exp(\frac{17.27(T_a - 273.15)}{(T_a - 273.15) + 237.3})]}{(T_a - 273.15 + 237.3)^2} \quad (17)$$

212 where 1000 is a unit conversion coefficient, T_a is the air temperature (K).

213 It is important to recognize that the range for G , T_s and RH_s can be refined in various ways. The
214 ranges defined above are just simple examples to determine the plausible ranges of these
215 parameters in wetland ecosystems, achieving more realistic results while reducing computation
216 time. The true ranges for G , T_s , and RH_s might be more constrained than these estimated values.
217 And many models, especially the models of G (e.g., the models listed in Purdy et al. (2016), can
218 be coupled with MaxEnt-ETRHEQ formulation. Exploring the potential enhancement of
219 MaxEnt-ETRHEQ's performance by integrating these models presents an intriguing subject for
220 future research.

221

222 **3. Data and model evaluation**

223 **3.1 Data**

224 All wetland sites classified as WET under the Vegetation IGBP category from the FLUXNET
225 2015 (Pastorello et al., 2020) and AmeriFlux (ameriflux.lbl.gov) FULLSET data products, shared
226 under the CC-BY-4.0 license, were chosen for this study. The characteristics of the sites include
227 latitude, longitude, elevation, mean measurement height, mean vegetation height, mean annual
228 temperature, mean annual precipitation, and the distance to the coast (Table S1 and Table S2).
229 Sites within 25 miles (~40 km) of the coast were removed, as ETRHEQ does not perform well in
230 coastal regions (Rigden and Salvucci, 2015). In addition, the sites without the measurements of
231 R_n and G were removed. The filter process results in 11 sites, including CZ-wet (Dušek et al.,
232 2016), DE-SfN (Klatt et al., 2016), DE-Zrk (Sachs et al., 2016), FI-Lom (Aurela et al., 2016),

233 US-Atq (Zona and Oechel, 2016a), US-BZB (Euskirchen, 2021a), US-BZF (Euskirchen, 2021b),
234 US-BZo (Euskirchen, 2022), US-ICs (Euskirchen et al., 2016), US-Ivo (Zona and Oechel,
235 2016b), and US-Los (Desai, 2016).

236 For every site, its fullset product encompasses five separate datasheets, containing measurements
237 of atmospheric variables and energy fluxes at half-hourly, daily, weekly, monthly, and annual
238 scales. At each temporal scale, u^* (“USTAR”), RH (“RH”), and R_n (“NETRAD”) as well as gap-
239 filled atmospheric measurements (denoted with the “_F” qualifier), including p (“PA_F”), T_a
240 (“TA_F”), u (“WS_F”), and VPD (“VPD_F”), and the energy fluxes with marginal distribution
241 sampling gap-filling method, which are G (“G_F_MDS”), H (“H_F”MDS”), and LE
242 (“LE_F_MDS”) were obtained. The names enclosed in double quotes within brackets in the
243 above sentence represent the variable names in the data products. RH at daily or larger scales
244 was not directly available, so it was estimated from VPD and T_a using the Clausius–Clapeyron
245 relation. Besides z and z_{veg} were provided in Table S1 and S2. For sites where z_{veg} is not
246 available, a value of 0.5 to z_{veg} was assigned.

247 The focus here was limited to temporal scales between half-hourly and monthly, due to a lack of
248 adequate sites and measurements for conducting a robust analysis at the yearly level. At the half-
249 hourly scale, data with poor quality (i.e., the quality flag (QC) >1) were removed. At coarser
250 temporal resolutions, i.e., from daily to monthly, only the measured data (QC=0) or the gap-filled
251 data where over 80% measured or good quality gap-filled (QC=1) records aggregated from finer
252 temporal resolutions were included, consistent with Kim et al. (2023). As a result, FI-Lom were
253 removed from daily to monthly scales due to the lack of the quality flag for G . At the monthly
254 scale, DE-sfN was also removed because only one measurement was available. In addition,
255 measurements were also removed if the surface energy imbalance was greater than $50 \text{ W}\cdot\text{m}^{-2}$

256 (McColl and Rigden, 2020; Rigden and Salvucci, 2015) or $R_n - G$ was negative (Kim et al.,
257 2023). The amount of data after all filters from half-hourly to monthly scales are presented in
258 Table S3. The model was also run for sites where soil moisture measurements were available
259 (i.e., US-BZB, US-BZF, US-BZo and US-ICs), to assess whether incorporating soil moisture
260 would enhance the model's performance, and the error statistics are provided in Table S4.

261 3.2 Model evaluation

262 The root-mean-square error (RMSE), slope and intercept (i.e., bias) of the fitted linear
263 relationship between modelled and measured energy fluxes, and the coefficient of determination
264 (R^2) were used as metrics to evaluate model performance. The evaluation was made of the
265 measurements without energy closure correction (specifically, “H_F_MDS” and “LE_F_MDS”
266 in the data product) and with correction using the energy balance closure correction factor on the
267 assumption that the Bowen ratio is correct (Pastorello et al., 2020) (the corrected energy fluxes
268 are “H_CORR” and “LE_CORR” in the data product), respectively. In addition, H_F_MDS and
269 LE_F_MDS were compared with H and LE calculated as the residual of the energy balance (i.e.,
270 $H_{re} = R_n - G_{F_MDS} - LE_{F_MDS}$, and $LE_{re} = R_n - G_{F_MDS} - H_{F_MDS}$) to assess the
271 how energy imbalance and the inherent uncertainty in the eddy covariance measurements affect
272 the evaluation of the performance of the model. If there is no energy closure problems in the
273 eddy covariance measurements, there would be a perfect fit between the measurements and the
274 residuals of the energy balance for each energy flux. This represents the highest level of
275 performance that can be expected from any model in comparison to eddy covariance
276 measurements, as explained in McColl and Rigden (2020). However, comparisons with other
277 models such as Penman-Monteith, Priestley-Taylor, MaxEnt, ETRHEQ, SFE, or MEP-SFE were
278 not conducted because MaxEnt-ETRHEQ is still in its early stage, and this paper is intended to

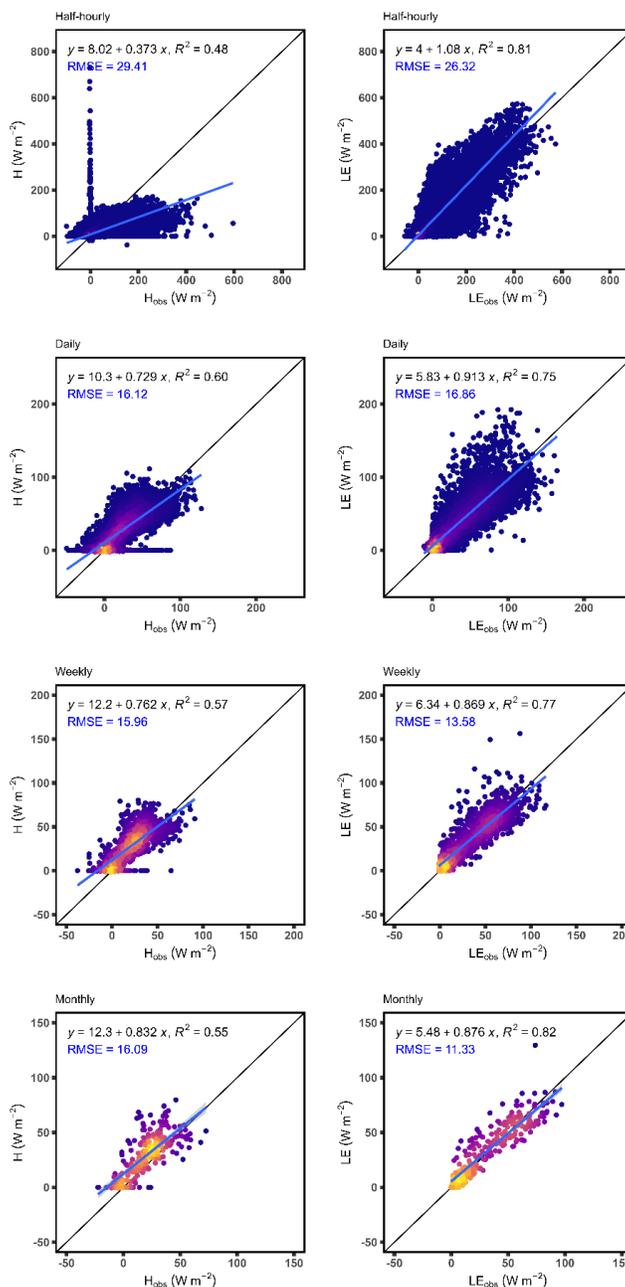
279 provide a possible way to utilize it for wetland ecosystems. Further, MaxEnt-ETRHEQ is unique
280 as it does not require G or T_s as inputs, unlike other models. However, inter-model comparisons
281 will be considered in future research.

282 All the analysis was conducted on R 4.3.0 (R Core Team 2023). The R scripts, which contain the
283 codes for calculating the distance from study sites to the coast, modelling the energy fluxes using
284 MaxEnt-ETRHEQ for each site, and creating the figures presented in this paper, are all available
285 at Wang (2024).

286

287 **4. Results**

288 MaxEnt-ETRHEQ provides highly accurate predictions for LE from half-hourly to monthly
289 scales (Figure 1), with slopes ranging from 0.86 to 1.08 and biases ranging from 4.00 to 6.34
290 $W \cdot m^{-2}$. When the energy balance residuals (e.g., H_{re} and LE_{re}) were used to compare with the
291 measurements (H_{F_MDS} and LE_{F_MDS}), their values of R^2 and the proximity of the slopes
292 to 1 show similar levels with MaxEnt-ETRHEQ, but their bias, which is around 14 to 16 $W \cdot m^{-2}$
293 and RMSE, which around 23 to 28 $W \cdot m^{-2}$ (Table 1), were slightly larger than those of MaxEnt-
294 ETRHEQ (bias: 4 to 7 $W \cdot m^{-2}$ and RMSE: 11 to 27 $W \cdot m^{-2}$). In this sense, the error statistics of
295 MaxEnt-ETRHEQ for estimating LE are slightly better than the errors statistics from the eddy
296 covariance measurements (Table 1).



297

298 **Figure 1. Modelled H and LE versus measurements (H_{obs} and LE_{obs}) without energy**
 299 **balance closure correction from half-hourly to monthly scales. The blue lines represent the**
 300 **fitted linear regressions. The black lines are 1:1 lines. The color of the points represents the**
 301 **density of the data ranging from low (purple) to medium (red) to high (yellow).**

302

303 On the other hand, the model does not predict H with the same performance as it predicts LE,
 304 especially at the half-hourly scale (Figure 1). But when the time scale becomes larger, the

305 performance on estimating H is improved (Figure 1 and Table 1). Overall, MaxEnt-ETRHEQ
306 tends to overestimate H when H is low and underestimate H when it is high (Figure 1). Given the
307 estimation of LE was quite accurate, the less satisfactory performance of MaxEnt-ETRHEQ for
308 H can be attributed to less accurately defined boundary conditions for G and T_s . In the current
309 model setting, G was limited to up to 20% percent of R_n based on the GLEAM model that was
310 designed for daily applications (Purdy et al., 2016). This explains why MaxEnt-ETRHEQ
311 performs better in estimating H at daily and larger time scales. But at most study sites, G often
312 exceeds 20% of R_n when R_n is exceptionally low (e.g., less than $50 \text{ W}\cdot\text{m}^{-2}$), and frequently falls
313 below 10% of R_n when R_n is high (greater than $400 \text{ W}\cdot\text{m}^{-2}$). Therefore, when R_n is low, G is
314 underestimated, leading to an overestimation of H, and vice versa.

315 At the half-hourly scale, there is a spike of estimated H when measured H is near zero (Figure 1).
316 The spike is only from the US-BZo site (Figure S1) that happened during the night when R_n ,
317 G_F_MDS , H_F_MDS and LE_F_MDS were all negative, and the absolute value of G was at
318 least 10 times larger than that of R_n . These energy fluxes suggest that US-BZo likely
319 encountered intense convective weather at these periods, characterized by air that was warmer
320 and more humid than the surface, accompanied by significant condensation. Under these weather
321 circumstances, the current MaxEnt-ETRHEQ formulations were unsuitable for H, LE, and g_a .
322 Determining the applicability of Eq. 1 in such conditions and devising revisions for the
323 calculations of H, LE, and g_a need to be addressed in future research.

324 When H and LE observations are adjusted to force energy balance closure (i.e., H_CORR and
325 LE_CORR), the performance of MaxEnt-ETRHEQ did not seem improve overall. For example,
326 the bias in H or LE estimations decrease when the energy fluxes are corrected at all temporal
327 scales, but the R^2 , RMSE and slopes deteriorate (Table 1). This is because the energy balance

328 closure correction results in higher H and LE for most of the study sites. While this adjustment
 329 could result in more accurate energy fluxes, it also has the potential for overcorrection as
 330 diagnosed in Mauder et al. (2018). Consequently, the actual performance of MaxEnt-ETRHEQ
 331 in estimating H and LE should be in between its performance when compared to uncorrected
 332 fluxes and its performance when evaluated against corrected fluxes.

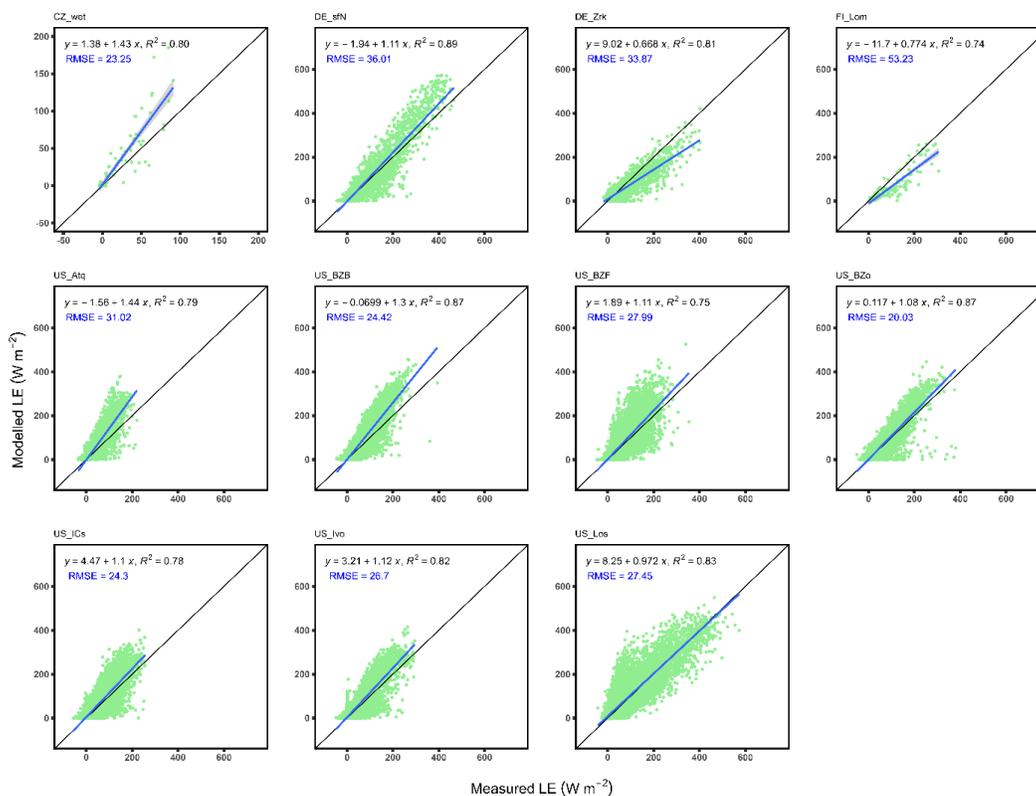
333 **Table 1. Summary of modelled fluxes against the energy balance corrected fluxes, and**
 334 **measured, uncorrected fluxes against the residuals of energy balance from half-hourly to**
 335 **monthly scales.**

Temporal scales	Variables		Slope	Intercept (bias)	R ²	RMSE
	x	y				
Half-hourly	H_CORR	Modelled H	0.27	7.62	0.47	43.87
	LE_CORR	Modelled LE	0.76	4.64	0.78	33.00
	H_F_MDS	H_re	0.90	12.80	0.69	27.87
	LE_F_MDS	LE_re	0.88	14.60	0.76	27.87
Daily	H_CORR	Modelled H	0.56	9.50	0.58	19.27
	LE_CORR	Modelled LE	0.72	5.59	0.74	20.66
	H_F_MDS	H_re	1.11	15.60	0.74	23.65
	LE_F_MDS	LE_re	1.08	15.10	0.82	23.65
Weekly	H_CORR	Modelled H	0.58	11.70	0.51	17.67
	LE_CORR	Modelled LE	0.70	4.95	0.78	18.23
	H_F_MDS	H_re	1.15	16.70	0.74	24.08
	LE_F_MDS	LE_re	1.13	15.40	0.84	24.08
Monthly	H_CORR	Modelled H	0.62	11.70	0.51	16.20
	LE_CORR	Modelled LE	0.68	4.23	0.81	18.27
	H_F_MDS	H_re	1.23	16.60	0.75	24.70
	LE_F_MDS	LE_re	1.18	15.10	0.86	24.70

336

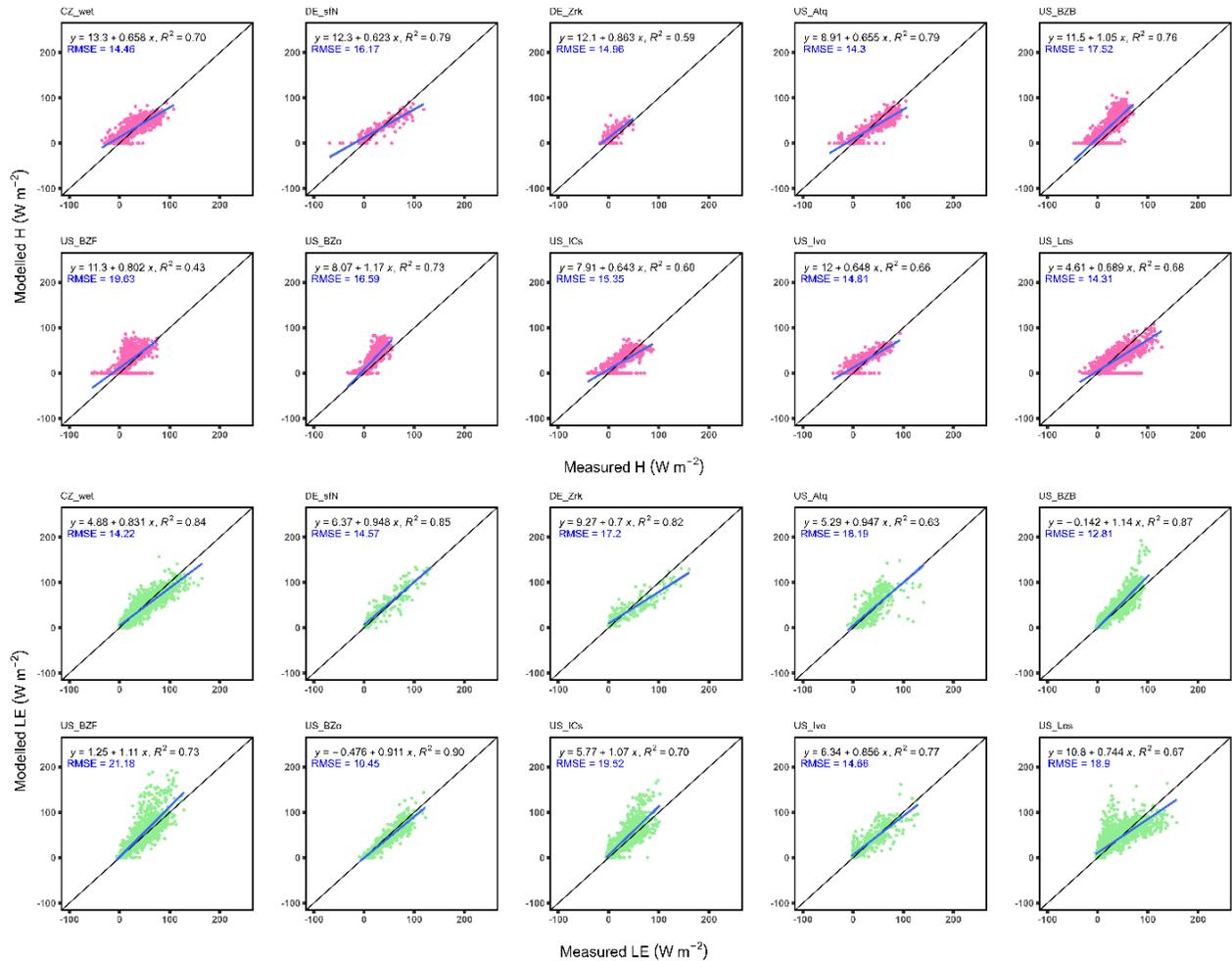
337 Finally, the performance of MaxEnt-ETRHEQ in estimating LE at individual sites throughout the
 338 temporal scales are also quite accurate. Figure 2 presents the half-hourly predictions of LE at
 339 each site, and shows that despite varying accuracy across different sites, MaxEnt-ETRHEQ
 340 demonstrates high precision in predicting half-hourly LE, with R² values between 0.74 to 0.89
 341 and RMSE ranging from 20.03 to 53.23 W·m⁻². However, the predictions of H at various

342 temporal scales were not satisfactory (Figure S1, Figure S2 and Figure S3). Nevertheless, when
 343 the time scale is at the daily, weekly or monthly, both H and LE estimations are improved (Figure
 344 3, Figure S2 and Figure S3). Considering that no site-specific calibration was made, and no T_s or
 345 G were used as inputs, the performance of MaxEnt-ETRHEQ at individual sites were excellent.



346

347 **Figure 2. Modelled LE versus measured LE without energy balance closure correction at**
 348 **the half-hourly scale at the study sites. The blue lines represent the fitted linear regressions.**
 349 **Black lines represent 1:1 lines.**



350

351 **Figure 3. Modelled H and LE versus measured H and LE without energy balance closure**
 352 **correction at the daily scale at the study sites. The blue lines represent the fitted linear**
 353 **regressions, and black lines d1:1 lines.**

354

355 **5. Model advantages and limitations**

356 The main advantage of MaxEnt-ETRHEQ is that it does not require land surface measurements
 357 like G and T_s , which outcompetes most evapotranspiration (ET) models. While it could be
 358 argued that ET from wetland ecosystems closely approximates potential ET, which can be easily
 359 calculated using the Priestley-Taylor or Penman-Monteith equations for saturated water surfaces,
 360 the computation of potential ET (PET) still necessitates at least G as input. Additionally,

361 wetlands may not consistently be in a state of saturation (Streich, 2019), and using these
362 equations could lead to substantial bias.

363 Moreover, MaxEnt-ETRHEQ is capable of providing estimates of LE at half-hourly intervals,
364 distinguishing it from most equilibrium-based models that require equilibration times that
365 typically extend beyond a daily timeframe, including the SFE model (McColl and Rigden, 2020)
366 and the SFE-MEP model (Kim et al., 2023). The highly accurate half-hourly LE estimates
367 provided by MaxEnt-ETRHEQ mean that the model is capable of precisely capturing the sub-
368 daily fluctuations of ET. Many land surface models have shown considerable inaccuracies in
369 sub-daily LE estimates, typically underestimating LE in the morning and overestimating it in the
370 afternoon, owing to insufficient parameterizations of stomatal conductance and plant hydraulics
371 (Matheny et al., 2014). MaxEnt-ETRHEQ and its underlying mechanism (i.e., the maximum
372 entropy production) may provide new perspectives to enhance the performance of these models.

373 However, MaxEnt-ETRHEQ is still in its early stages, as further efforts are required to
374 accurately refine the ranges of G , T_s and RH_s . However, that does not mean that these land
375 surface variables ought to be inputs for MaxEnt-ETRHEQ. Rather, identifying appropriate
376 boundary conditions for these variables should suffice. With growing evidence showing the
377 interactions between land surface variables like G , T_s , soil moisture, soil thermal inertia and
378 vegetation properties and near-surface atmospheric conditions (Bennett et al., 2008; Chu et al.,
379 2018; Purdy et al., 2016; Wang and Bras, 1999; Wang and Bou-Zeid, 2012), developing physical
380 models to describe these linkages and determining the limiting cases for G , T_s and RH_s are not
381 far off. Once these boundary conditions are defined properly, MaxEnt-ETRHEQ will be capable
382 of simultaneously estimating not only H , LE , and G , but also T_s and RH_s . Thus, it opens up a
383 promising avenue for future research.

384 In addition, it may be argued that MaxEnt-ETRHEQ relies on empirical parameters like the
385 parameterization of I_s and g_a . Indeed, most models for estimating surface energy fluxes are
386 largely based on empirical approaches, particularly in calculating parameters such as
387 displacement height, roughness length for momentum and heat transfer, and aerodynamic
388 conductance. Furthermore, when these models are scaled up for application over extensive areas,
389 the reliance on parameters that have been either assumed or previously calibrated becomes
390 inevitable. Therefore, the use of empirical parameterizations in MaxEnt-ETRHEQ should not be
391 viewed as shortcomings. Instead, it underscores the critical need for further research aimed at
392 refining these parameters to enhance the model's accuracy.

393

394 **6. Conclusion**

395 The goal of this paper is to demonstrate the effectiveness of a newly developed formulation
396 grounded in the principle of maximum Shannon information entropy production theory for
397 estimating surface energy fluxes in wetland ecosystems. The formulation requires neither land
398 surface variables nor site-specific calibration, except for a presumed vegetation height, and it
399 effectively estimates LE from half-hourly to monthly scales in the FLUXNET and AmeriFlux
400 wetland sites. While its estimation on H is less satisfactory due to roughly constrained boundary
401 conditions for G and T_s , the formulation holds promise for concurrently and accurately
402 estimating LE, H, G, T_s and RH_s for various ecosystems if limiting cases of G, T_s and RH_s are
403 properly established. Overall, the formulation contributes new insights into developing earth
404 systems models.

405

406 **Open research**

407 All datasets in this study, as well as the R scripts used for modeling and data visualization, are
408 publicly available. For access to the specific datasets used in this study, please refer to the
409 FLUXNET database (<http://www.fluxnet.org>) and the AmeriFlux network
410 (<http://ameriflux.lbl.gov>). For the data analysis, the R programming language version 4.3.0 (R
411 Core Team 2023) was employed. The codes can be accessed on Wang, Y. (2024). R scripts for
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414

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435

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