

Detection of slow slip events along the southern Peru - northern Chile subduction zone

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Abstract Detections of slow slip events (SSEs) are now common along most plate boundary fault systems globally. However, no such event has been described in the south Peru - north Chile subduction zone so far, except for the early preparatory phase of the 2014 Iquique earthquake. We use geodetic template matching on GNSS-derived time series of surface motion in Southern Peru - Northern Chile to extract SSEs hidden within geodetic noise. We detect 24 events with durations ranging from 17 to 36 days and magnitudes from M_w 5.4 to 6.2. Our events, analyzed from a moment-duration scaling perspective, reveal values consistent with observations reported in other subduction zones. We compare the distribution of SSEs with the distribution of coupling along the megathrust derived using Bayesian inference on GNSS- and InSAR-derived interseismic velocities. From this comparison, we obtain that most SSEs occur in regions of intermediate coupling where the megathrust transitions from locked to creeping or where geometrical complexities of the interplate region have been proposed. We finally discuss the potential role of fluids as a triggering mechanism for SSEs in the area.

Resumen Hoy en día, las detecciones de eventos lentos (SSEs, por sus siglas en inglés) son comunes a lo largo de la mayoría de los sistemas de fallas activas a una escala global. Sin embargo, hasta ahora, no se han reportado eventos de este tipo en la zona de subducción del sur del Perú y norte de Chile (10°S-24°S), exceptuando aquellos ocurridos durante la fase de preparación del terremoto de Iquique de 2014. En el presente trabajo, nosotros utilizamos una técnica conocida como "Template Matching" en series temporales de desplazamiento medido por datos GNSS (Global Navigation Satellite System, GNSS por sus siglas en inglés) en el sur del Perú y el norte de Chile, para extraer la firma de eventos lentos asísmicos ocultos en el ruido geodésico. Nosotros detectamos

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34 24 eventos asísmicos con duraciones de 17 a 36 días, y magnitudes de M_w 5.4 a 6.2. El análisis de
35 nuestros eventos utilizando leyes de escala momento-duración, revela valores consistentes con
36 observaciones realizadas en otras zonas de subducción. El momento sísmico liberado por estos
37 eventos es proporcional al cubo de su duración, lo que parece implicar una dinámica comparable
38 con la de los terremotos clásicos. Los eventos detectados en este trabajo están principalmente lo-
39 calizados en zonas donde el acoplamiento intersísmico presenta valores en transición (0.3 - 0.8 de
40 factor de acoplamiento), donde la zona de subducción transiciona de un estado bloqueado a uno
41 de deslizamiento continuo. Finalmente, nosotros discutimos el rol potencial que podrían jugar los
42 fluidos en el desencadenamiento de estos eventos lentos.

43 **Résumé** Depuis une vingtaine d'année, des événements de glissement asismiques ont été dé-
44 tectés le long de quasiment toutes les frontières de plaques au monde. Cependant, aucun n'a été
45 décrit pour l'instant le long de la zone de Subduction allant du Pérou au nord du Chili, si l'on omet le
46 glissement mesuré lors de la période d'activité ayant mené au séisme d'Iquique en 2014. Nous util-
47 isons une technique dite de Template matching sur des séries temporelles de déplacement mesuré
48 par GNSS dans le nord du Chili pour extraire la signature d'événements de glissement asismiques
49 cachés au sein du bruit géodésique. Nous détectons 24 événements asismiques avec des durées
50 allant de 17 à 36 jours pour des magnitudes équivalentes allant de M_w 5.4 à 6.2. Nos événements
51 ont des valeurs cohérentes avec les observations rapportées dans d'autres zones de subduction.
52 Il apparaît que ces événements asismiques sont essentiellement localisés dans des zones de cou-
53 plage intermédiaires où le megathrust est à mi-chemin entre un état bloqué et un état en glisse-
54 ment permanent. Nous discutons finalement de l'influence éventuelle de fluides profonds dans le
55 déclenchement de ces événements asismiques.

56 **Non-technical summary** Earthquakes correspond to a sudden release of elastic energy
57 stored in the crust as a response to the relative motion of tectonic plates. However, this release
58 of energy is not always sudden and accompanied by destructive seismic waves. It sometimes hap-
59 pens slowly during aseismic, slow slip events. It has been shown that SSEs can be associated with
60 the nucleation, propagation, and termination of big earthquakes. SSEs have been detected along
61 many subduction zones in the world but not in northern Chile, yet. Here, we use a template match-
62 ing method to scan GNSS observations of ground motion to detect and characterize slow slip events
63 along the southern Peru - northern Chile subduction zone. We find 24 aseismic events at depths
64 comparable with that of SSEs in other subduction zones, as well as in regions that slip aseismically
65 persistently. We discuss how our findings relate to past earthquake ruptures, the geometry of the
66 subduction zone, and fluids circulating at depth. Our results show the importance of implementing
67 methods to extract small aseismic signals in noisy data, key observations for a better understand-
68 ing of fault mechanics.

1 Introduction

Overwhelming evidence suggest that the Elastic Rebound Theory proposed by Reid (1910) after the 1906 California earthquake associated with the stick-slip behavior of frictional interface (Brace and Byerlee, 1966) is insufficient to explain the slip behavior along active faults. Geodetic measurements of surface motion have revealed the presence of aseismic, slow slip along all types of active faults. After the first descriptions in the mid-20th century from direct observations of damage to human-made structures crossing the San Andreas (Louderback, 1942; Steinbrugge et al., 1960) and North Anatolian (Ambraseys, 1970) faults, aseismic slip has been directly observed, or inferred, from geodetic measurements at different stages of the earthquake cycle. For instance, afterslip corresponds to the diffusion of slow slip during the post-seismic period accommodating a co-seismic stress perturbation (e.g., Heki et al., 1997; Bürgmann et al., 2001; Hsu et al., 2002, 2006). Creep, on the other hand, often refers to steady aseismic slip during the interseismic period (Steinbrugge et al., 1960; Ambraseys, 1970; Jolivet et al., 2015b). In addition, interseismic transients (i.e., slow slip events or SSEs) during this interseismic period were discovered in the 2000s along subduction zones. SSEs often locate in the deeper portion of the seismogenic zone (e.g., Hirose et al., 1999; Dragert et al., 2001), but some of these SSEs are associated with seismic signals that occur within the seismogenic zone, and may contribute to reducing geodetic coupling (Mazzotti et al., 2000; Bürgmann et al., 2005; Loveless and Meade, 2010; Radiguet et al., 2012; Béjar-Pizarro et al., 2013; Villegas-Lanza et al., 2016; Métois et al., 2016; Michel et al., 2019a; Jolivet et al., 2020; van Rijnsingen et al., 2021; Loverly et al., 2024). This along-dip segmentation differs from one subduction zone to the other (Nishikawa et al., 2019) and we note more occurrences of SSEs along young, warm subduction zones (i.e., Nankai, Mexico, Cascadia), than old and cold ones. Finally, slow slip appears to be an important ingredient of the preparation phase of earthquakes (e.g., Ruegg et al., 2001; Ruiz et al., 2014; Radiguet et al., 2016; Socquet et al., 2017; Voss et al., 2018). More recently, it has been proposed that a significant fraction of observed geodetic displacement in seismically active regions results from the occurrence of slow slip events (Jolivet and Frank, 2020, and reference therein), suggesting a burst-like, episodic behavior of aseismic slip at all time scales from seconds to decades in places as varied as Mexico (Frank, 2016; Rousset et al., 2017; Frank and Brodsky, 2019), Cascadia (Michel et al., 2019a; Ducellier et al., 2022; Itoh et al., 2022), along the San Andreas Fault (Khoshmanesh and Shirzaei, 2018; Rousset et al., 2019; Michel et al., 2022), the Haiyuan fault in Tibet (Jolivet et al., 2015a; Li et al., 2021), on the Alto Tiberina and Pollino fault systems in Italy (Gualandi et al., 2017; Cheloni et al., 2017; Essing and Poli, 2022), or Japan (Nishimura et al., 2013; Takagi et al., 2019; Nishikawa et al., 2019; Uchida et al., 2020). All observations suggest the importance of accounting for aseismic slip in our understanding of earthquake cycle dynamics. However, the underlying physics controlling aseismic slip is still debated, mainly due to the lack of good, dense observational databases.

Nowadays, observations of aseismic slip in subduction zones are frequently documented over a wide range of slip amplitudes and at different stages of the earthquake cycle (Avouac, 2015; Obara and Kato, 2016; Bürgmann, 2018; Kato and Ben-Zion, 2021, and references therein). Regular slow slip events have been documented mainly along warm subduction zones such as Cascadia, Nankai (southwest Japan), Mexico, or New Zealand (e.g., Graham et al., 2016; Nishikawa et al., 2019; Wallace, 2020; Michel et al., 2022, and references therein). Instead, observations of slow slip events in cold subduction zones such as off-shore Japan or Chile are sparse or indirect, through seismic swarms, repeaters, or slow earthquakes (Kato et al., 2012; Kato and Nakagawa, 2014; Gardonio et al., 2018; Nishikawa et al.,

2019), and rarely with geodetic observations (Hino et al., 2014; Ruiz et al., 2014; Socquet et al., 2017; Boudin et al., 2021). Geodetic displacement corresponding to such slow slip events are usually of mm to cm-scale amplitude and require the development of novel and systematic methods to extract SSEs from noisy time series of geodetic data (Frank, 2016; Rousset et al., 2017; Michel et al., 2019a; Uchida et al., 2020; Itoh et al., 2022).

We focus on the South Peru- North Chile subduction zone. The region is seismically active, with two historical earthquakes in 1868 (southern Peru), and 1877 (northern Chile), both tsunamigenic earthquakes of magnitude ~ 8.5 (Kausel, 1986; Comte and Pardo, 1991; Vigny and Klein, 2022) (Figure 1). Since these two events, the region has experienced several large earthquakes ($M_w > 7.5$) (Ruiz and Madariaga, 2018) accompanied by an important background seismic activity (Jara et al., 2017; Sippl et al., 2018, 2023) (Figure 1). In addition, coupling is highly variable along the subduction interface. Coupled regions overlap with the inferred rupture extent of the 2001 M_w 8.1 Arequipa and 2014 M_w 8.1 Iquique earthquakes (Schurr et al., 2014; Métois et al., 2016; Villegas-Lanza et al., 2016; Jolivet et al., 2020). A large coupled section is inferred where the 1877 earthquake is thought to have ruptured (Jolivet et al., 2020; Vigny and Klein, 2022). In addition, two low-coupling regions are observed. In southern Peru, low coupling coincides with the subduction of the Nazca ridge ($\sim 15^\circ$) (Villegas-Lanza et al., 2016; Loverly et al., 2024). In northern Chile, a reduction in coupling is inferred offshore Iquique and below the Mejillones peninsula ($\sim 21^\circ$) (Béjar-Pizarro et al., 2013; Métois et al., 2016; Jolivet et al., 2020).

In addition to low coupling, aseismic slip has been observed in South Peru and North Chile. Afterslip has been reported following large earthquakes, including the 1995 M_w 8.1 Antofagasta (Chlieh et al., 2004; Pritchard and Simons, 2006), the 2001 M_w 8.1 Arequipa (Ruegg et al., 2001; Melbourne, 2002), the 2007 M_w 8.0 Pisco (Perfettini et al., 2010; Remy et al., 2016), the 2007 M_w 7.7 Tocopilla (Béjar-Pizarro et al., 2010) and the 2014 M_w 8.1 Iquique earthquakes (Hoffmann et al., 2018) (Figure 1). Geodetic transients interpreted as the signature of aseismic slip occurred in the days to months preceding the M_w 8.4 Arequipa earthquake in 2001, before one of its largest aftershock, and preceding the Iquique earthquake in 2014 (e.g., Ruegg et al., 2001; Melbourne, 2002; Ruiz et al., 2014; Schurr et al., 2014; Socquet et al., 2017). Aseismic slip is considered responsible for a significant fraction of such geodetic transients (Twardzik et al., 2022). There is therefore plenty of evidence of occurrences of aseismic slip in this broad region but, despite intense efforts to instrument the area, no obvious spontaneous slow slip events have been detected during the interseismic period.

A change in the interseismic surface velocity field was observed following the M_w 7.5 intermediate-depth Tarapaca earthquake over a decade (Peyrat et al., 2006; Peyrat and Favreau, 2010) (Figure 1), an observation interpreted as the signature of a decoupling of the subduction interface (Ruiz et al., 2014; Jara et al., 2017). Comparable changes in surface velocity field, observed following the 2010 Maule earthquake, have also been observed in the regions affected by the 2015 Illapel (Ruiz et al., 2016) and 2016 Chiloé (Ruiz et al., 2017; Melnick et al., 2017) earthquakes. Such shifts in surface velocity may be linked to postseismic viscoelastic processes acting over long distances (Bouchon et al., 2018) in contrast to the localized behavior observed after the Tarapaca earthquake (Jara et al., 2017). Over the same period, we observed a significant increase in background seismicity (Jara et al., 2017), as well as an apparent synchronization of intermediate-depth and shallow seismic activities (Bouchon et al., 2016; Jara et al., 2017). Changes in background seismicity rates have been associated with the occurrence of aseismic slip events and fluid migration (Marsan et al.,

143 2013; [Reverso et al., 2016](#); [Marsan et al., 2017](#)). The synchronization of the seismicity is interpreted as related to
144 aseismic slip events occurring along the subduction interface due to a broader slab deformation ([Bouchon et al.,](#)
145 [2016](#)). These indirect observations suggest aseismic transients may occur in South Peru - North Chile during the
146 interseismic period.

147 We aim to detect small, short-term aseismic slip events in this region and discuss their occurrence and location
148 with respect to the interseismic coupling pattern and past seismic crises. We explore GNSS time series, searching for
149 small transients, using a geodetic template matching approach ([Rousset et al., 2017](#)). We use GNSS and InSAR data to
150 infer an updated distribution of interseismic coupling using a Bayesian framework following the approach of [Jolivet](#)
151 [et al. \(2020\)](#), comparing the detected aseismic events with the coupling model, along with geophysical information
152 available in the region (seismicity, Vp/Vs ratio, gravity models). We finally discuss potential mechanisms explaining
153 the occurrence of aseismic events in the area.

154 **2 Data, Methods and Results**

155 **2.1 GNSS processing and time series analysis**

156 We process data from 119 continuous GNSS (cGNSS) sites in the central Andes region (Figure S1a) and worldwide
157 (Figure S1b), using a double difference approach with the GAMIT/GLOBK software ([Herring et al., 2015](#)). 67 cGNSS
158 sites are in the South Peru - North Chile region (Figure S1a and Figure 2, brown arrows), installed and maintained
159 by the Integrate Plate boundary Observatory Chile (IPOC) ([Klotz et al., 2017](#)), the Laboratoire International Asso-
160 cié “Montessus de Ballore” (LIA-MB) ([Klein et al., 2022](#)), the Central Andean Tectonic Observatory (CANTO, Caltech)
161 ([Simons et al., 2010](#)), the Instituto Geofísico del Perú (IPG) ([Jara et al., 2017](#); [Socquet et al., 2017](#)), the Institut des
162 Sciences de la Terre (ISTerre) ([Jara et al., 2017](#); [Socquet et al., 2017](#)), and the Centro Sismológico Nacional of Chile
163 (CSN) ([Báez et al., 2018](#)). The remaining 52 stations are part of the International GNSS Service (IGS) ([Teunissen and](#)
164 [Montenbruck, 2017](#)) global network. We separate these stations into three subnetworks (two locals and one global)
165 with 33 overlapping stations, where the local separation depends on the station data span: one local network with
166 data from 2000-2014 and the other including data from 2007-2014. Global network processing includes 99 stations
167 over the 2000 - 2014 period, with 22 stations in South America (Figure S1b). We use the GAMIT 10.6 software ([Herring](#)
168 [et al., 2015](#)), choosing the ionosphere-free combinations and fixing the ambiguities to integer values. We use precise
169 orbits from the IGS, precise earth-orientation parameters (EOPs) from the International Earth Rotation and Refer-
170 ence System Service (IERS) bulletin B, IGS tables to describe the phase centers of the antennas, FES2004 ocean-tidal
171 loading corrections, and atmospheric loading corrections (tidal and non-tidal). We estimate one tropospheric zenith
172 delay every two hours and one pair of horizontal tropospheric gradients per 24h session using the Vienna Mapping
173 Function (VMF1) ([Boehm et al., 2006](#)). We use the GLOBK software to combine daily solutions and the PYACS software
174 ([Nocquet, 2018](#)) to derive position time series in the ITRF 2008 reference frame ([Altamimi et al., 2011](#)). Finally, time
175 series are referenced to fixed South-America considering the Euler pole solution proposed by [Nocquet et al. \(2014\)](#).

176 We fit the time series with a parametric function of time for each component (N, E, and U) ([Bevis and Brown,](#)
177 [2014](#)). Each time series $x(t)$, function of time t , is modeled as

$$x(t) = x_R + v(t - t_R) + \sum_{j=1}^{n_j} b_j H(t - t_j) + \sum_{k=1}^{n_F} [s_k \sin(\omega_k t) + c_k \cos(\omega_k t)] + \sum_{i=i}^{n_T} a_i \log(1 + t_i/\Delta T), \quad (1)$$

where x_R is a reference position at a time t_R and v is the interseismic velocity for each component. H is a Heaviside function applied each time t_j an earthquake (or antenna change) offsets the time series. The combination of sin and cos functions describes seasonal oscillations (with annual and semi-annual periods), while the logarithmic function models the transient, post-seismic signal following large earthquakes ($M_w \geq 7.5$) with a relaxation time ΔT . For a given station, we consider a Heaviside function for all earthquakes of magnitude larger than 6 with an epicenter to station distance lower than $d(M) = 10^{\frac{M}{2}-0.8}$, as proposed by the Nevada Geodetic Laboratory (www.geodesy.unr.edu). We only include a post-seismic term for earthquakes of magnitude larger than 7.5. All inferred parameters for each component and each cGPS site are in Supplementary Information, Tables S1-S38. Figures S2-S17 compare the data and model at each station. We then estimate and remove a common-mode error by stacking all the time series (Bock and Melgar, 2016; Socquet et al., 2017; Jara et al., 2017). This procedure enables us to get residual time series (Figures S18-S19) as well as an interseismic velocity field (Table S1-S2). We use the obtained residual time series to search for geodetic transients compatible with slip on the megathrust and use the geodetic velocity field to update the last published coupling map (Jolivet et al., 2020).

2.2 Fault Geometry and Green's Functions

Coupling map estimation and geodetic template matching methods need a fault geometry and Green's functions calculation, as described below. In both cases, we define the geometry of the megathrust using Slab 2.0 (Hayes et al., 2018) as a reference, but with different meshing strategies. For the coupling case, we use triangles with 10 km-long sides along the coast and 25 km-long sides, both at the trench and depth, between latitudes 17°S-25°S. In the northern part (10°S-17°S), we adapt the size to the GNSS station density, considering a constant 50 km-long triangle side. In contrast, in the geodetic template matching case, we use triangles with 10 km-long sides along the coast and 25 km-long sides in the entire region. Then, we consider slip on the fault as the linear interpolation of slip values at the mesh nodes. Finally, we compute the Green's functions assuming a stratified elastic medium derived from Husen et al. (1999) using the EDKS software (Zhu and Rivera, 2002).

2.3 Coupling map for Southern Peru - Northern Chile

We update the distribution of coupling from Jolivet et al. (2020) in order to compare short- (i.e., days to months) and long-term (i.e., years to decades) aseismic deformation in the region. We use the GNSS velocity fields from Métois et al. (2016) (data span 1996 - 2013) and Villegas-Lanza et al. (2016) (data span 2008-2013), that we complement with our GNSS velocity field (Figure 2a, data span 2000-2016). Additionally, we use the line of sight (LOS) velocity map from Jolivet et al. (2020), derived from the processing of Envisat data covering the period 2003 - 2010 (Figure 2b).

We use the backslip approach to estimate the distribution of coupling (Savage, 1983). A coupling of 1 (resp. 0) corresponds to a fully locked megathrust (resp. a megathrust that slips at plate rate). We consider plate motion estimated by UNAVCO (www.unavco.org) under the ITRF 2014 model (Altamimi et al., 2016) to estimate the convergence

211 rate, angle, and rake on each node of the fault mesh. The backslip rate is evaluated by subtracting the sliver move-
 212 ment proposed by Métois et al. (2016) in Chile (11 mm/yr) and by (Villegas-Lanza et al., 2016) in Peru (5.5 mm/yr) to
 213 the convergence rate. In the Arica bend (16°S - 18°S), at the boundary of the Chilean and Peruvian slivers, we build
 214 a gradient to make a smooth transition between the two slivers. We solve for the distribution of models that satisfy
 215 the geodetic data.

216 The forward problem is written as $\mathbf{d} = \mathbf{G}\mathbf{m}$, with \mathbf{d} the geodetic data (GNSS and InSAR velocities), \mathbf{m} the vector
 217 of parameters to solve for and \mathbf{G} the Green's functions (Section 2.2). Parameters include coupling at each mesh node
 218 and geometric transformations akin to those in Jolivet et al. (2020). We adopt a probabilistic approach to estimate
 219 the parameters in order to evaluate the associated uncertainties. The *a posteriori* Probability Density Function (PDF)
 220 of a model \mathbf{m} given a dataset \mathbf{d} , $p(\mathbf{m}|\mathbf{d})$, writes as

$$221 \quad p(\mathbf{m}|\mathbf{d}) \propto p(\mathbf{m})p(\mathbf{d}|\mathbf{m}), \quad (2)$$

222 where $p(\mathbf{m})$ is the *a priori* model PDF and $p(\mathbf{d}|\mathbf{m})$ is the data likelihood. The *a priori* PDF describes our knowledge
 223 of coupling along the megathrust before collecting geodetic data. We define the *a priori* PDF at each node for the
 224 coupling factor as follows:

$$225 \quad \mathbf{X} \sim \begin{cases} \mathcal{N}(\mu_c, \sigma_c^2) & \text{if } -0.1 \leq \mathbf{X} \leq 1.1 \\ 0 & \text{otherwise} \end{cases} \quad (3)$$

226 where μ_c and σ_c are the mean and standard deviation of a normal distribution. We select the bounds of [-0.1, 1.1] to
 227 ensure an accurate sampling for the full range of coupling values between 0 and 1 (Dal Zilio et al., 2020a; Jolivet et al.,
 228 2020). We know the megathrust is decoupled below 60 km depth from geodetic (Chlieh et al., 2004; Béjar-Pizarro
 229 et al., 2013; Jolivet et al., 2020) and seismological evidence (Comte et al., 2016). Thus, we apply an *a priori* condition
 230 based on the depth of each node. If a node is deeper than 60 km, the *a priori* mean (μ_c) is set to 0 and the standard
 231 deviation (σ_c) to 0.1. In cases where a node is shallower than 60 km, we assign an *a priori* mean (μ_c) of 0.5 and a
 232 standard deviation (σ_c) of 0.5.

233 We adopt a Gaussian formulation for the data likelihood, $p(\mathbf{d}|\mathbf{m})$, which writes as

$$234 \quad p(\mathbf{d}|\mathbf{m}) = \frac{1}{\sqrt{2}\mathbf{C}_\chi} \exp\left\{-\frac{1}{2}(\mathbf{G}\mathbf{m} - \mathbf{d})^T \mathbf{C}_\chi^{-1}(\mathbf{G}\mathbf{m} - \mathbf{d})\right\}, \quad (4)$$

235 where \mathbf{C}_χ is the misfit covariance matrix (Duputel et al., 2014) defined as $\mathbf{C}_\chi = \mathbf{C}_p + \mathbf{C}_d$, where \mathbf{C}_d is the data
 236 covariance matrix (data uncertainties), while \mathbf{C}_p is the prediction error covariance matrix, representing uncertainties
 237 on the assumed elastic model (P and S wave velocities and density). We assume a 10% error on the elastic parameters
 238 following Jolivet et al. (2020).

239 We explore the model space using Altar (altar.readthedocs.io) to sample the *a posteriori* PDF of the coupling factor,
 240 generating 250000 models. Altar is based on the Cascading Adaptive Transitional Metropolis in Parallel (CATMIP)
 241 algorithm (Minson et al., 2013; Duputel et al., 2014; Jolivet et al., 2015b). These models enable us to perform statistics,
 242 derive the mean model for the interseismic coupling (Figure 3), and collect information about the model resolution

(see Supporting Information for model GNSS and InSAR residuals, Figure S20-S23, as well as Standard Deviation, Mode, Skewness, and Kurtosis, Figure S24).

The mean coupling model (Figure 3a), is close to previously published models in the region (e.g., Chlieh et al., 2011; Béjar-Pizarro et al., 2013; Métois et al., 2016; Villegas-Lanza et al., 2016; Jolivet et al., 2020; Loverly et al., 2024), especially considering the along-strike segmentation. Our model differs from previously published models in the coupling intensity at locked patches, as well as the depth of these coupled patches. In Peru, we observe three patches with interseismic coupling that varies between 0.5-0.75 (Figure 3a). Previous models report similar patches, although totally locked (coupling factor ~ 1) (Chlieh et al., 2011; Villegas-Lanza et al., 2016; Loverly et al., 2024). Unfortunately, the density of GNSS stations in this region is not anywhere near that in Chile, hence the large standard deviations in the Peruvian region (Figure S25). Analyzing the moments of the *a posteriori* PDF, including standard deviation, skewness and kurtosis confirms this (Figure S24). Similarly, these moments show that the resolution at the trench over the entire region is low. Additionally, our model varies from those constrained only by GPS data in Chile (e.g., Métois et al., 2016). The InSAR data helps constraining interseismic coupling at depth (Béjar-Pizarro et al., 2013; Jolivet et al., 2020) and the strong *a priori* coupling damps potential large variations at depth, which we consider not physical.

2.4 Detection of aseismic slip events with geodetic template matching

2.4.1 Methodology

We use a geodetic template matching approach to detect potential aseismic slip events on the residual GNSS time series (Section 2.1). We summarize here the method presented in detail by Rousset et al. (2019). We search for the spatio-temporal signature of slip events in surface displacement time series by cross-correlating synthetic templates with our GNSS residual time series, in velocity. These templates correspond to the surface displacement caused by slip on dislocations located on the subduction megathrust embedded in a stratified, semi-infinite elastic medium. We calculate such templates (\mathbf{w}) by convolving the Green's functions (Section 2.2) with a time-dependent slip evolution $\mathbf{s}(t)$ defined as

$$\mathbf{s}(t) = \frac{1}{2} \left[1 - \cos \left(\frac{\pi t}{T} \right) \right], \quad (5)$$

where T is the duration of a synthetic event. Following Rousset et al. (2019), we derive for each template the weighted correlation function for each fault node, defined as

$$\mathbf{C}_f(t) = \frac{\sum_{i=1}^{2N} |\mathbf{G}_i| \mathbf{C}_i(t)}{\sum_{i=1}^{2N} |\mathbf{G}_i|}, \quad (6)$$

where \mathbf{G} is the Green's functions and \mathbf{C}_i is the correlation between the time series and the synthetic template at a given fault node i given by

$$\mathbf{C}_i(t) = \frac{\sum_{k=1}^T \dot{\mathbf{w}}_i(t_k) \dot{\mathbf{d}}_i(t_k + \tau)}{\sqrt{\sum_{k=1}^T \dot{\mathbf{w}}_i^2(t_k) \sum_{k=1}^T \dot{\mathbf{d}}_i^2(t_k + \tau)}}, \quad (7)$$

where $\dot{\mathbf{w}}$ and $\dot{\mathbf{d}}$ are the time derivatives of the template in terms of displacement (i.e., the template's velocity with duration T), and the time derivatives of the GNSS time series, respectively. τ denotes a moving time variable that enables the temporal matching search between templates and observations. We then search for peaks in $\mathbf{C}_f(t)$ corresponding to candidate slip events. As can be seen in the Supporting Information (see Fig. S31b, red and black lines), in the case of synthetic events, the correlation peaks in C_f arise from the geodetic noise using as many GNSS stations as possible.

For each candidate slip event, we stack the time series of displacement weighted by Green's functions around the time of detection (see Supporting Information Figure S31b, for an example of stacks on synthetic time series, purple and yellow lines). Such weighting accounts for displacement amplitude and direction, increasing the signal-to-noise ratio (Rousset et al., 2017). Stacks are computed over a period of 180 days, centered on each potential occurrence. On each stack, we estimate two linear trends, before and after the candidate occurrence, and the time dependent slip evolution of Eq. 5 to the weighted stack in order to determine the amplitude, the start and end date of each detected transient. We apply a non-linear regression to determine the posterior Probability Density Function of the model parameters given a stack of time series following Tarantola (2005). Effectively, we use an MCMC algorithm to derive 30,000 samples from the posterior PDF and evaluate the mean and standard deviation of the duration and magnitude of each candidate slow slip event.

In order to curate the potential detections from artefacts, we perform a sensitivity and resolution analysis, to determine the minimum magnitude of a slip event that can be detected for each fault node. Although the method above has been extensively described by Rousset et al. (2019), the novelty of our approach relies on the evaluation of uncertainties through a Bayesian exploration of all important parameters.

2.4.2 GNSS network sensitivity and resolution

We analyze the sensitivity of our approach by testing its ability to detect, locate, and estimate the source parameters (magnitude and duration) of synthetic aseismic slip events. We first evaluate the parameters characterizing the noise affecting each GNSS time series of displacement by building synthetic time series of noise on which we perform the tests. In order to generate synthetic noise, we model each component of the residual time series (Eq. 1) as a combination of white and colored noise (Williams, 2003), such as,

$$\mathbf{P}(f) = P_0 (\mathbf{f}^{-\alpha} + f_0^{-\alpha}), \quad (8)$$

where \mathbf{P} is the power spectrum as function of temporal frequency \mathbf{f} , P_0 and f_0 are normalization constants, and α is the spectral index. We explore P_0 , f_0 , and α using Bayesian inference to estimate their mean and standard deviation at each station component (see the Supporting Information for further details and an example of the power spectrum and the probability density function (PDF) of parameters at the UAPE station in Figures S26 - S27, as well as

305 Tables S39 - S42 for all the network noise parameters inferred). We use these inferred noise parameters to build 1000
306 synthetic time series of displacement at each GNSS station. We use these synthetic time series to estimate thresholds
307 of detection for each fault node.

308 The number of GNSS stations in the study area has evolved during the observation period. We, therefore, must
309 consider three periods independently depending on the number of active stations: 2000 - 2003 (four stations), 2004
310 - 2007 (20 stations), and 2008 - 2014 (55 stations). We first determine which stations are able to capture a slow slip
311 event on a given node. For each period and fault node, we correlate the 1000 synthetic time series of noise with a
312 template of a duration of 40 days and slip equivalent to a magnitude M_w 6.0. We evaluate the standard deviation of
313 the resulting weighted correlation functions, σ_t , as a minimal threshold to be exceeded (i.e., when dealing with time
314 series that might include slip events, a peak of correlation higher than $3\sigma_t$ is a positive detection).

315 Once this threshold has been defined, we compute the weighted correlation function for 1000 time series of noise
316 to which we have added the signal of synthetic transients with different duration (10, 20, and 30 days) and magnitudes
317 (5.0 - 7.0 M_w , every 0.1 of magnitude). In case of a detection, we stack the displacement time series around the detec-
318 tion time. We consider a synthetic event has been correctly detected and located if we can recover four quantities,
319 including the slip event location, timing, duration, and magnitude. If the estimated location is within 150 km from
320 the true location, if the estimated timing and duration are within five days of the actual ones, and if the estimated
321 magnitude is within 0.25 of the actual one, we consider the detection to be valid. This procedure enables us to deter-
322 mine the minimum magnitude that can be detected over each of the three observation periods and build resolution
323 maps for each period investigated (see Supporting Information, Figures S29-S30). For instance, in the Iquique region
324 ($\sim 19^\circ\text{S} - 71^\circ\text{W}$), the minimal magnitude M_w ranges from 6.6 to 6.8 from 2000 to 2003, decreases to 6.1-6.3 from 2004
325 to 2007 and again down to 5.9 to 6.1 from 2008 to 2014. Thus, as expected, we observe a significant improvement in
326 detection sensitivity when the number of stations in a given region increases.

327 **2.5 Application to GNSS time series**

328 After exploring the network sensitivity to detect aseismic slip events, we search for transients in the residual time
329 series obtained after subtracting the trajectory model described earlier. We fix the duration T of the template to 40
330 days and the slip to an event equivalent to M_w 6.0 (see Supporting Information, Figures S58-S59 for a test in the dura-
331 tion template sensitivity). By doing so, we detect 733 candidate slip events in the stacked correlation functions. Since
332 some of these candidates may correspond to the same candidate slip event, we retain maximum occurrences within
333 a radius of 150 km (i.e., if two maxima affect nodes separated by a distance higher than 150 km, they are considered as
334 independent occurrences). After this selection step, we are left with 59 candidate slip events in the region. We eval-
335 uate their durations and magnitudes and compare these with our resolution maps. We keep candidates for which
336 the obtained magnitude is higher than the minimum detectable magnitude for the corresponding node (Figure 4),
337 leaving us with 24 validated slip events.

338 The duration of the slip events ranges from 17 to 36 days with magnitudes from M_w 5.4 to 6.2 and depths from
339 20 to 66 km. Figure 3 shows the location of the detected slip events along with four examples of weighted stacks.
340 Figures 5 and 6 show two examples of stacks and correlation functions, along with the time series used to build the
341 stacks and the map view of correlation peaks (see Supplementary information Tables S43 for the event parameters

342 estimated with their uncertainties, and Figures S33 - S43 to see the data employed in the modeling, the data stack,
343 and the model).

344 Following the methodology proposed by [Nishimura et al. \(2013\)](#), validated events are categorized into two types:
345 probable and possible. This classification is achieved by comparing the displacement fields derived directly from
346 observations with those generated by synthetic events of estimated magnitudes. Note that the magnitudes are es-
347 timated on the correlation stack and not directly on the measured displacements. A disagreement between the dis-
348 placements corresponding to the detected magnitude on the detected node and the observed displacement would
349 suggest our assumptions do not hold. Observed displacements are determined directly on the GNSS time series by
350 estimating a linear trend along with a time-dependent slip evolution (Eq. 5). To estimate the displacement field for a
351 detected magnitude, the slip corresponding to that magnitude is applied at the inferred location of each event. Fig-
352 ures 5 (b) and (c) illustrate examples of these estimates, with the actual displacements shown in magenta, while the
353 displacements predicted from the magnitudes of each event are shown in green for Events #10 and #12 (see Supple-
354 mentary Information, Figures S33 - S43 for the rest of the events). Upon analysis, we find that the agreement between
355 observed and modeled ground motion is acceptable for 10 of our events, leading us to classify these as probable (A
356 events, Table S43). Meanwhile, we observe a weaker agreement for 14 events which we hence categorize as possible
357 (B events, Table S43).

358 Since our template matching approach only considers GNSS observations, we must ensure that the detected slip
359 events (A and B) are mostly aseismic. We cross-check the 24 positive detections with the seismic catalog provided
360 by the ISC ([International Seismological Centre, 2016](#)). We randomly generate 10000 synthetic locations for each slip
361 event considering a normally distributed location uncertainty based on our resolution tests and estimate the sum
362 of the seismic moment of all earthquakes occurring within at least a $2\text{-}\sigma$ radius of the detected slip event. We then
363 compare this estimate of the seismic moment to the estimated aseismic one. All the detected slip events have an
364 equivalent magnitude at least twice larger than the seismic magnitude (aseismic/seismic ratio for each event and
365 further details on ratio estimation are in Supplementary Information, Table S43). Figures 5 and 6 (d) present the
366 location of the two events detailed in Figures 5 and 6 (a) together with the seismicity that coincides with the occur-
367 rence of the slip event. These two events occur during the preparation phase of the 2014 Iquique earthquake (Event
368 #12, Figure 1) and during the interseismic phase (Event #10). The combination of synthetic tests and the seismic vs.
369 aseismic moment analysis confirms we detected 24 aseismic slip events (A and B) along southern Peru - northern
370 Chile subduction zone over the period 2006 - 2014.

371 **3 Discussion**

372 **3.1 Aseismic slip events and scaling laws**

373 Aseismic slip events are now frequently observed along most subduction zones in the world, but the underlying
374 physics is still debated. Among the points of debate, the comparison between slow slip and earthquakes should allow
375 to point out whether comparable physics are involved. [Ide et al. \(2007\)](#) have proposed that, while the seismic moment
376 of earthquakes is proportional to the cube of their duration, the moment of slow earthquakes, from tremors and low-
377 frequency earthquakes to slow slip events, is proportional to the duration. Considering that simple considerations

about size and stress drop led to the emergence of the observed scaling for earthquakes, the difference in moment-duration scaling should involve a fundamental difference between the mechanics of slow slip and that of earthquakes. [Peng and Gomberg \(2010\)](#) argued that the apparent moment duration scaling of slow earthquakes proposed by [Ide et al. \(2007\)](#) was only due to a lack of observations, suggesting both rapid and slow slip were driven by the same mechanism, namely a slip instability with variable speed and stress drop propagating along a weakened fault surface. In addition, [Gomberg et al. \(2016\)](#) proposed that seismic moment scales either with the duration or the cube of the duration depending on whether the rupture was elongated and pulse-like or mostly crack-like. [Michel et al. \(2019b\)](#) confirmed that the moment of slow slip events in Cascadia scales with the cube of their duration although being elongated and pulse-like. These observations agree with recent studies of aseismic slip and tremors in Japan ([Takagi et al., 2019](#); [Supino et al., 2020](#)) and Mexico ([Frank and Brodsky, 2019](#)), as well as numerical modeling using dynamic simulations of frictional sliding ([Dal Zilio et al., 2020b](#)). Such numerical and observational evidence suggests that SSEs might exhibit comparable scaling as classical earthquakes, only with lower rupture speeds and stress drops.

We evaluate the scaling between moment and duration for the aseismic slip events we have detected. We estimate that the moment, M , is such as $M \propto T^{4.99 \pm 0.48}$, with T the duration for the 24 detected SSEs (refer to Figures 7, S45, and S46 in the Supporting Information for an in-depth explanation of the scaling estimation procedure). This scaling relationship remains consistent when analyzing events A ($M \propto T^{5.05 \pm 0.59}$, see Figures S47 and S48) and B ($M \propto T^{4.89 \pm 0.52}$, illustrated in Figures S49 and S50) independently. Our events seem to align with a moment-duration scaling T^3 . However, as extensively discussed by [Ide and Beroza \(2023\)](#), uncertainties associated with the estimation of event duration might influence significantly our results. Consequently, it is challenging to definitively conclude that our findings adhere to the moment-duration T^3 scaling. That said, our detections are situated within the range of moment-duration observed in other subduction zones such as Cascadia, Japan, or Mexico ([Ide and Beroza, 2023](#), and references therein). Building on this observation, we adopt the methodology outlined by [Gomberg et al. \(2016\)](#) to deduce the source properties of our events. We infer that the rupture velocities of our detections range between 0.5 and 10 km/day, accompanied by a stress drop of 0.1 MPa (see the Supporting Information for detailed information on the parameter estimation process). Although our method does not allow to detect events that would propagate, we observe our SSEs are more compatible with crack-like, unbounded ruptures than pulse-like, bounded ones. As a conclusion, our findings along southern Peru - northern Chile region align with SSEs observations from other subduction zones.

3.2 Aseismic slip and interseismic coupling distribution

Our coupling estimate corresponds to an average behavior over a decade, without accounting for potential slow slip events hidden within the noise. The slow slip events we detect hence correspond to fluctuations around this average. We compare the map of coupling to the location of our 24 aseismic events to explore how such fluctuations distribute with respect to locked and creeping asperities along the megathrust (Figure 3). We compare the distribution of coupling where our events are located to a distribution coupling at randomly picked locations (Figure 8, see the Supporting Information for a detailed explanation of the calculation of the PDF for coupling and detected events). The distributions differ but mostly when considering only events in northern Chile, where our estimate of coupling is much more robust. Detected slow slip events occur mostly in regions of intermediate coupling. This observation

415 is not as clear for the Peruvian region, probably because of the sparsity of the data used here, although the same ten-
416 dency is suggested on Figure 8. This result aligns with Frank (2016) findings in the Mexico subduction zone, where a
417 database of slow slip events seems to compensate the lack of slip deficit in transition zones with respect to coupled
418 regions of the megathrust. Materna et al. (2019) describe a comparable behavior over longer periods where cou-
419 pling variations seem to occur in regions of transitional coupling (Michel et al., 2019a). In addition, events offshore
420 Peru tend to cluster spatially around locked asperities, areas that are generally of intermediate coupling (Figure 9).
421 In general, slow slip events occur in transitional regions between seismic asperities and freely slipping areas. This
422 is consistent with model predictions from rate-and-state friction in which slow slip events are expected to occur at
423 the transition between seismic, rate-weakening and creeping, rate-strengthening asperities (e.g., Liu and Rice, 2005,
424 2007; Perfettini and Ampuero, 2008).

425 The average depth of the detected slow slip events is 33 km (Figure 8, see the Supporting Information for a detailed
426 explanation of the PDF calculation). Separating the events, by region, yields an average depth of 37 km for Peru and
427 30 km for northern Chile with comparable standard deviations (19 and 10 km respectively, Figure 8). This result
428 remains consistent when conducting separate analyses of events A and B (refer to Figures S56-S57 in the Supporting
429 Information). Lay (2015) separates the subduction megathrust along depth into four domains (A, B, C, and D). Domain
430 A, located between the trench and a depth of about 15 km, hosts either tsunami earthquakes or aseismic deformation.
431 Domain B, between approximately 15 and 30 km depth, hosts large megathrust earthquakes. Domain C, between
432 approximately 30 and 50 km depth, hosts intermediate sized earthquakes. At greater depths, Domain D, between
433 50 and 70 km, hosts slow slip events, tremors, and very low-frequency earthquakes. Our slow slip events mainly
434 occur in Domains C and D. It is understood that small, velocity weakening asperities in Domain C are embedded in
435 conditionally stable regions of the megathrust, prone to host slow slip events. Domain D is dominated by aseismic
436 sliding and potential slip rate variations could explain deeper detections. Therefore, the depth distribution of our
437 events matches regions where slow slip events are expected in a subduction zone context.

438 Our resolution tests (Figures S24, S29-S30) suggest that it is impossible to capture aseismic slip near the trench,
439 in domain A, with the current GNSS network. However, large, shallow slow slip events have been observed in Japan
440 (Nishimura, 2014; Nishikawa et al., 2019) and New Zealand (Wallace, 2020). Seafloor geodesy might help to detect
441 the occurrence of such large events and potentially for small, cm-scale ones comparable to our aseismic slip events
442 (Araki et al., 2017). Additionally, stress-shadow induces apparent coupling in velocity-weakening regions, especially
443 late in the interseismic period (Hetland and Simons, 2010; Lindsey et al., 2021). For this reason, we also cannot rule
444 out the potential occurrence of aseismic slip event near the trench.

445 In addition to the depth-dependent segmentation, we observe an along-strike segmentation in the distribution
446 of SSEs. In particular, we observe a lack of events within the rupture area of the 1877 earthquake, within the Are-
447 quipa rupture area and other detections gather around locked asperities, like in the doughnut model for seismicity
448 (Kanamori, 1981; Schurr et al., 2020). Such configuration is comparable to that of the Japan trench where the as-
449 perity that ruptured during the Tohoku earthquake in 2011 overwhelms the simple depth-dependent distribution of
450 behavior from Lay (2015). In particular, Nishikawa et al. (2019) propose that, unlike the Nankai subduction interface
451 which exhibits a depth-dependent segmentation due to a young, warm slab, the megathrust beneath Tohoku is not

452 segmented at depth into four distinct domains. In our area of interest, the subducting slab is older than the Nankai
453 slab and probably colder (Müller et al., 2008), which would explain why the behavior we unravel is not completely
454 consistent with that of Lay (2015) and potentially closer to that of the Japan trench.

455 As an additional level of complexity, three events coincide with the subduction of the Nazca ridge (14°S, Figures 3
456 and 9a), six events are located beneath the Mejillones Peninsula (23°S, Figures 3 and 9d), and three events are within
457 the Arica bend (17°S - 19°S, Figures 3 and 9b and c). These morphological structures are anomalies compared to the
458 model proposed by Lay (2015) as they are considered as barriers to the propagation of large earthquakes (Armijo and
459 Thiele, 1990; Comte and Pardo, 1991; Béjar-Pizarro et al., 2010; Villegas-Lanza et al., 2016; Poli et al., 2017). In these
460 regions, the depth of our detected slow slip events does not match the depth-dependency described by Lay (2015).
461 We can speculate that local geometrical complexities may lead to the occurrence of slow slip events (Romanet et al.,
462 2018) in the case of the subduction of the Nazca Ridge or that the apparent low coupling is the result of multiple slow
463 slip events (Jolivet et al., 2020) in the case of the Arica Bend.

464 **3.3 Aseismic slip events before and after large earthquakes**

465 Among all the detected slow slip events, only events #7, and #12 (Figure 3, S36 and 5) do not occur during the steady
466 interseismic period. Event #7 locates in the region struck by the Iquique earthquake in 2014 (Figure 9c, and S36)
467 during the post-seismic relaxation that followed the mainshock (Meng et al., 2015; Hoffmann et al., 2018; Shrivastava
468 et al., 2019) (M_w 6.1 and duration of 28 days in June 2014). Such slow slip events embedded within a post-seismic
469 sequence have already been observed following the Illapel earthquake (Tissandier et al., 2023) and in a completely
470 different setting, following the 2004 Parkfield earthquake, along the San Andreas Fault (Michel et al., 2022).

471 Aseismic slip has been recognized as an important element of the earthquake preparation phase (Obara and
472 Kato, 2016; McLaskey, 2019; Kato and Ben-Zion, 2021, and references therein). An 8-month-long slow slip event was
473 reported before the Iquique earthquake in 2014 (Socquet et al., 2017), and event #12 coincides with one of the regions
474 of the megathrust that slipped aseismically during that preparation phase (Figure 9 c). In addition, event #12 occurred
475 where and when intermediate-depth and shallow seismicity synchronized before the Iquique earthquake (Bouchon
476 et al., 2016; Jara et al., 2017) (M_w 6.0 and duration of 30 days in January 2014). Such synchronization of seismicity
477 began in January 2014, lasted for one month, and is interpreted as evidence of a slow, slab-wide deformation process
478 prior to megathrust earthquakes (Bouchon et al., 2016). Furthermore, event #12 is coincident with the transient
479 event reported by Boudin et al. (2021) using a long-base tiltmeter. Our epicentral location differs by ~50km from
480 the one reported by (Boudin et al., 2021), a difference that can be explained by different modeling strategies and/or
481 uncertainties. We propose that event #12 is linked to the 8-month aseismic slip transient observed preceding the
482 2014 Iquique earthquake. Such detection suggests the growing instability preceding the Iquique earthquake exhibits a
483 complex spatio-temporal behavior that hides within the noise of the data, in agreement with the hypothesis proposed
484 by Jolivet and Frank (2020) and Twardzik et al. (2022).

485 **3.4 Aseismic slip and fluids**

486 Fluids may also play a role in the occurrence of aseismic slip events (Avouac, 2015; Harris, 2017; Jolivet and Frank,
487 2020, and references therein). Pore pressure affects fault normal stress, hence modify the probability of a slip in-

488 stability as well as the nucleation size (Liu and Rice, 2007; Avouac, 2015; Bayart et al., 2016; Harris, 2017; Bürgmann,
489 2018; Jolivet and Frank, 2020; Behr and Bürgmann, 2021). An increase in pore pressure within the fault zone leads
490 to a decrease in normal stress, which promotes slip but increases nucleation size, promoting slow slip. We compare
491 our detections to the distribution of the Vp/Vs ratio and to gravity-inferred structural models in the region. We use
492 the Vp/Vs ratio inferred by Comte et al. (2016) for the events located in Northern Chile. Statistically, the 17 aseismic
493 events in northern Chile are not related to a specific Vp/Vs value (Figure 8, see the Supporting Information for a de-
494 tailed explanation of the PDF calculation). In particular, no slow slip events are found to collocate with high Vp/Vs
495 ratios ($Vp/Vs > 1.8$) (Comte et al., 2016) (Figure S44).

496 We also compare the location of our aseismic events to a 3-D density model in the region (Tassara and Echaurren,
497 2012). Figure 10 shows the location of aseismic events along ten different trench-perpendicular cross sections. The
498 slow slip events are primarily located along the contact between the slab and the overriding lithospheric mantle
499 (Figure 10, see Figure S51 for an analysis of depth uncertainties). This mantle corner is principally hydrated by the
500 dehydration of the subducting slab due to water releasing metamorphic reactions (Peacock, 2001; Rüpke et al., 2004;
501 Comte et al., 2016; Wang et al., 2019; Contreras-Reyes et al., 2021). The fact that our aseismic slip events tend to cluster
502 at depths corresponding to the lithospheric mantle along the megathrust, and not deeper, might imply that fluids may
503 be trapped and accumulate below the continental Moho, an hypothesis that would require further investigations.

504 **4 Conclusions**

505 We have systematically analyzed GNSS time series in the region, searching for the occurrence of aseismic slip events
506 with a template matching approach. We find 24 events in the period 2006 - 2014, with durations of 17 - 36 days,
507 magnitudes of M_w 5.4 - 6.2, and located at depths of 20-66 km. These events are mostly aseismic and are observed at all
508 stages of the earthquake cycle, including during post-seismic periods (afterslip, one event), earthquake preparation
509 phase (one event), and interseismic period (22 events). We compare those slow slip occurrence to a wide range of
510 possible models of interseismic coupling based on GNSS and InSAR velocity fields and infer a distribution of coupling
511 along the megathrust.

512 By conducting a moment-duration scaling analysis, we find that our observations are consistent with values re-
513 ported in subduction zones globally. We do not find particular correlations with published seismic velocity structures
514 but find that slow slip events cluster around past ruptures and locked asperities, where the megathrust transitions
515 from sliding to locked. Additionally, our events are located in regions of intermediate coupling values and mean
516 depths of 33 km, which match regions where slow slip events occur in the context of subduction zones.

517 Some of these events occur on the subduction interface deeper than than the continental MOHO, i.e. where the
518 slab is in contact with the mantle wedge corner where fluids are supposedly trapped. This points toward the influence
519 of fluids as it may explain both their spontaneous triggering and their long duration. However, as some events are
520 found at shallower depth, the involvement of fluids might not be the only explanation. Other mechanisms such as
521 geometrical complexities might be involved but more evidence are required.

522 The main outcome of this study is that we found numerous aseismic slip events in a place where none were found
523 during the interseismic period before. As a consequence, aseismic slip events may be found elsewhere in subduction

524 zone contexts where experts did not find any event, pending dedicated noise analysis methods. We provide here one
525 piece of evidence supporting the hypothesis proposed by Jolivet and Frank (2020) which states that slow slip happens
526 everywhere and at all times.

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541 Data and code availability

542 GNSS data are processed using GAMIT software (Herring et al., 2015) (<http://geoweb.mit.edu/gg/>), while the reference
543 frame is defined using PYACS (Nocquet, 2018) (<https://github.com/JMNocquet/pyacs36>). GNSS time series used in this
544 work can be found at: <https://doi.org/10.5281/zenodo.7898656>. The modeling has been performed using the Classic
545 Slip Inversion library (Jolivet et al., 2015b) (CSI, <https://github.com/jolivetr/csi>) and AltTar (Minson et al., 2013) (<https://github.com/AltTarFramework/altar>). All plots are made using Matplotlib (Hunter, 2007) and Cartopy (Office, 2010)
546 Python packages.
547

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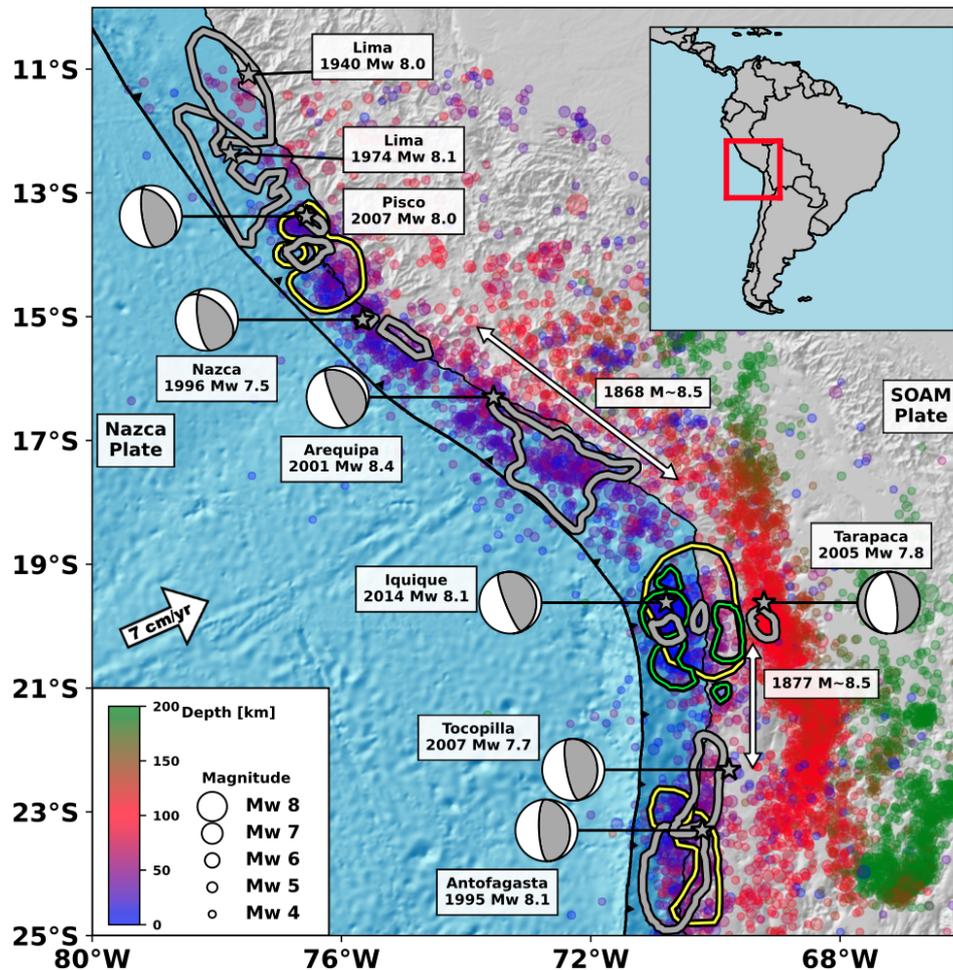


Figure 1 Seismotectonic map of the South Peru - North Chile subduction zone. White arrows show the extent of historical earthquakes (Comte and Pardo, 1991; Vigny and Klein, 2022). Gray contours are the rupture area of instrumental earthquakes with $M > 7.5$, with corresponding epicenters (gray stars) and focal mechanisms (if available) (Dorbath et al., 1990; Beck and Ruff, 1989; Hartzell and Langer, 1993; Delouis et al., 1997; Chlieh et al., 2004; Pritchard et al., 2007; Dziewonski et al., 1981; Ekström et al., 2012; Peyrat and Favreau, 2010; Sladen et al., 2010; Béjar-Pizarro et al., 2010; Duputel et al., 2015; Jara et al., 2018). Yellow lines are the 0.1 m afterslip contours available in the region (Chlieh et al., 2004; Béjar-Pizarro et al., 2010; Remy et al., 2016; Hoffmann et al., 2018), whereas the green ones are the pre-seismic slip reported for Iquique earthquake by Socquet et al. (2017). Colored dots are earthquakes with $M > 4.0$ from the International Seismological Centre (International Seismological Centre, 2016) over the period 1990 - 2016, color-coded by depth and scaled by magnitude. Large white arrow shows convergence direction and rate from Métois et al. (2016). SOAM: South America plate.

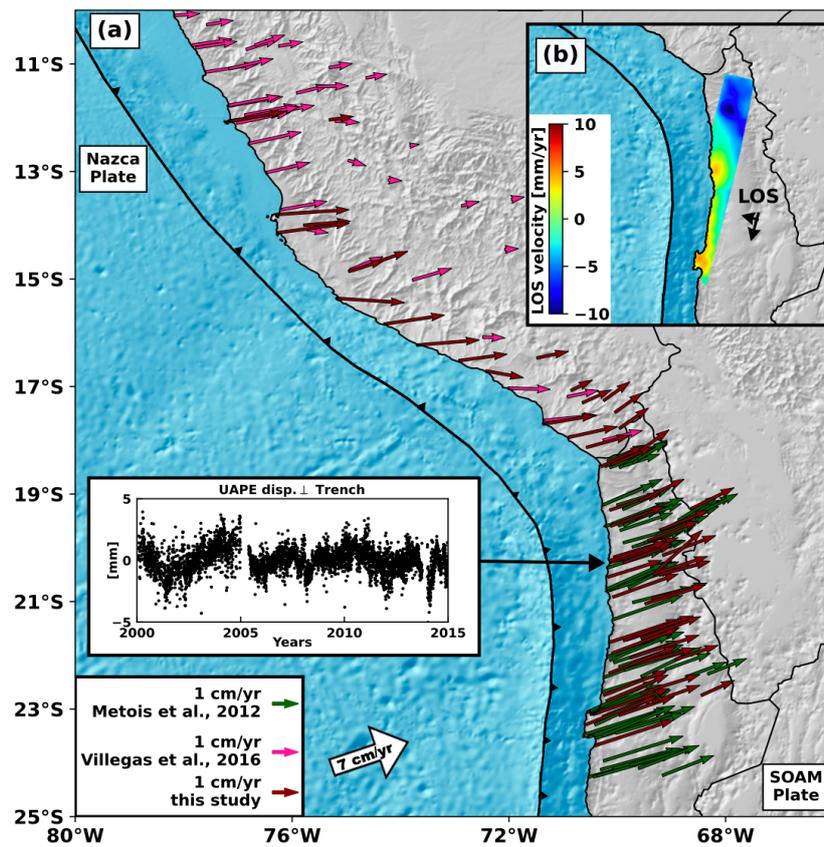


Figure 2 Geodetic data. (a) Colored dark green and pink arrows are the GNSS interseismic velocities from Métóis et al. (2016) and Villegas-Lanza et al. (2016), respectively, while brown arrows are the continuous GNSS processed in this study. The inset shows the residual trench perpendicular displacement time series for GNSS station UAPE. (b) Line-of-sight (LOS) interseismic ground velocity from track 96 (Envisat data) from (Jolivet and Simons, 2018; Jolivet et al., 2020). Black arrows indicate the flight direction of the satellite and its line of sight (LOS).

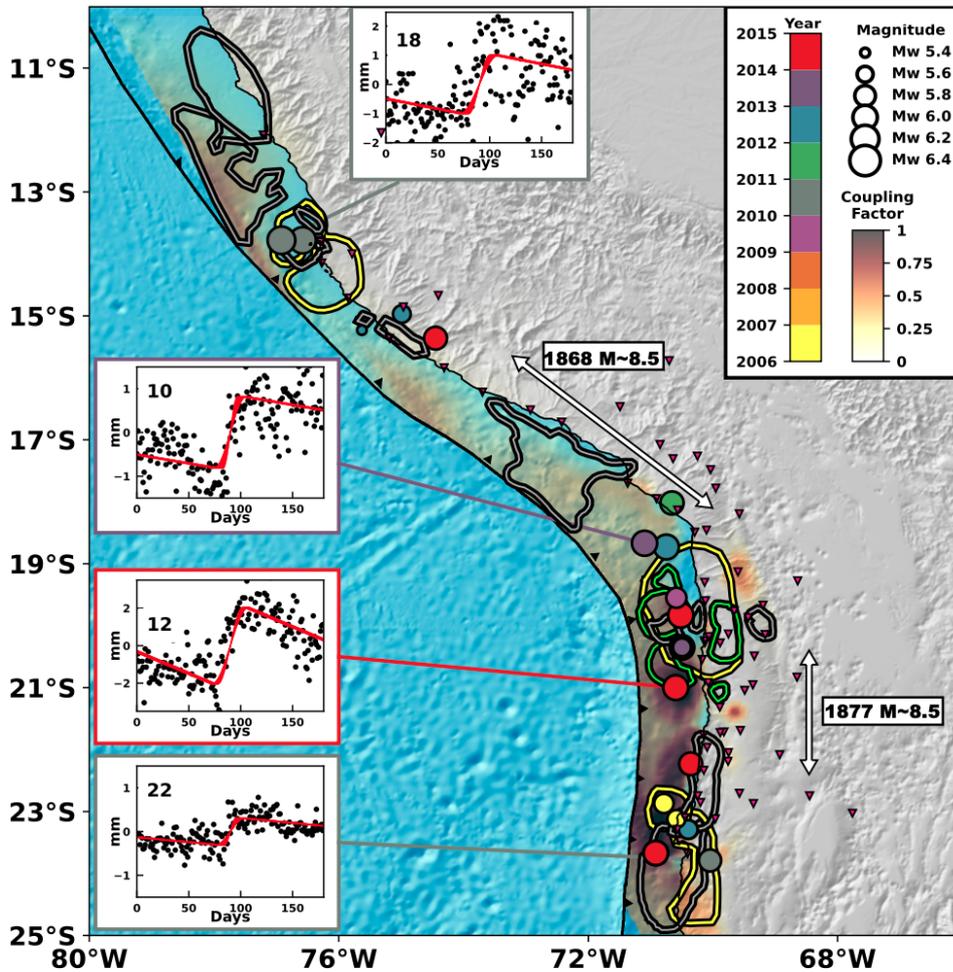


Figure 3 Location of detected aseismic slip events. Markers are color-coded by time of occurrence and scaled by magnitude. Four examples of weighted stacked correlations are shown with the event id number. Red line is the best fit model used to evaluate the event magnitude and duration, considering their estimated σ . Background color from white to dark through yellow and red is the mean coupling distribution. Black red areas (coupling factor ~ 1) are locked regions, while transparent areas (coupling factor ~ 0) are regions that slip aseismically at a rate equal to the plate convergence rate. Gray contours show instrumental ruptures. Yellow contours are afterslip regions, whereas green ones indicate slip inferred during the period preceding the Iquique earthquake. White arrows are the historical rupture extents.

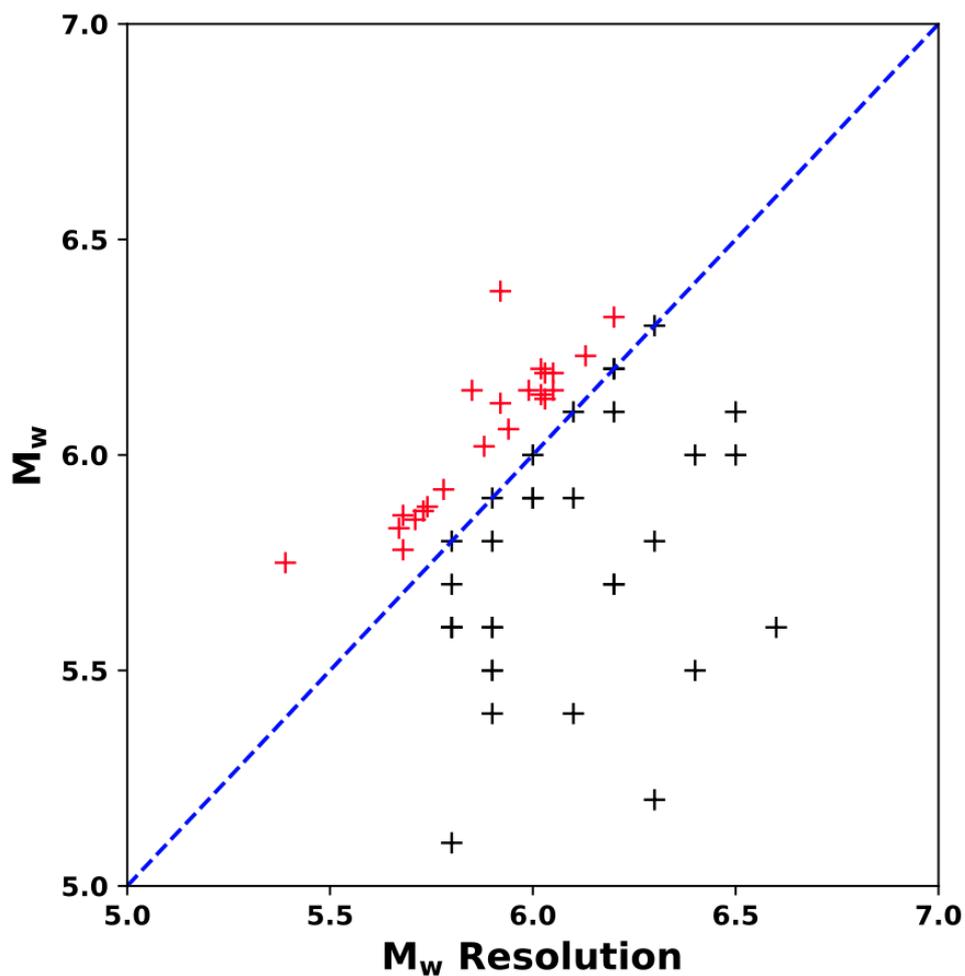


Figure 4 Event magnitude as a function of the resolution magnitude of the node where the event is located. Red crosses are events that passed the resolution test. Dashed blue line is the 1:1 line that separates validated from excluded events.

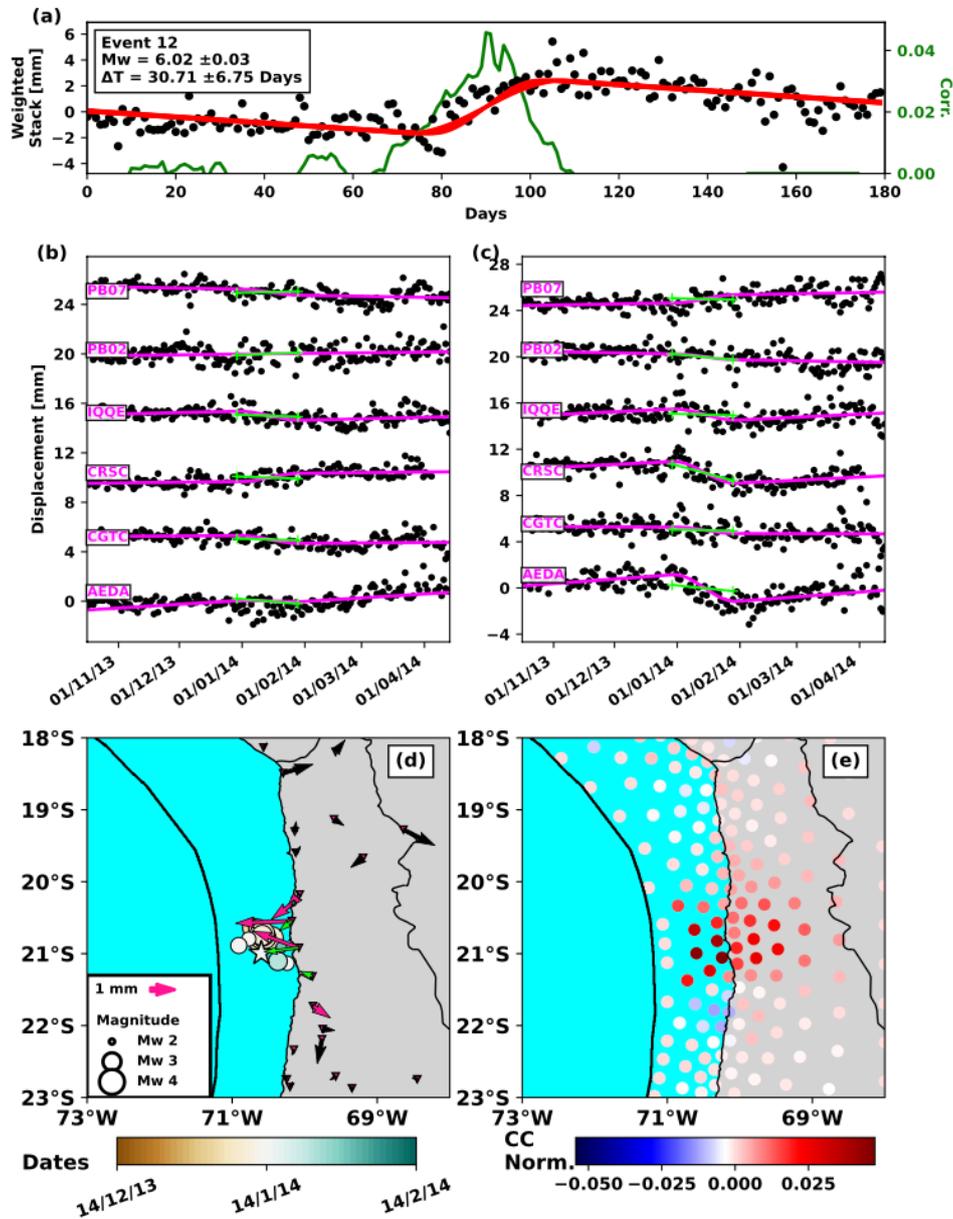


Figure 5 Example of detected aseismic slip event #12 in the vicinity of the 2014 Iquique earthquake, its locations, and associated seismicity. Figure (a) features the weighted stack for the event #12, with the red line representing the preferred model used to estimate event duration and magnitude, as indicated at the top left. The dark green line denotes the correlation function where event detection is made. Figures (b) and (c) display the displacement time series for the North and East components, respectively. Displacement data from six stations contributing to the weighted stack are shown. The pink lines indicate the best-fitting model for each displacement time series, which incorporates a linear trend and a transient, in accordance with Eq. 5. Meanwhile, the green lines represent the displacements for the estimated magnitude of each event. Figure (d) illustrates the event location (marked by white star), with dots indicating seismicity before and after the event (spanning half of the event's duration for each period), scaled by magnitude and color-coded by date. Inverted triangles mark the GNSS station locations. Pink arrows denote the GNSS-derived displacements from observations used to estimate the weighted stack during the detected slow slip event, whereas black arrows indicate displacements not used in the estimation. The green arrows show displacements resulting from dislocations for the estimated magnitudes at each event location (white star). Figure (e) displays the map view of the correlation peak within the correlation function (illustrated in dark green in Figures a) for the event, pinpointing the moment when the detection is made.

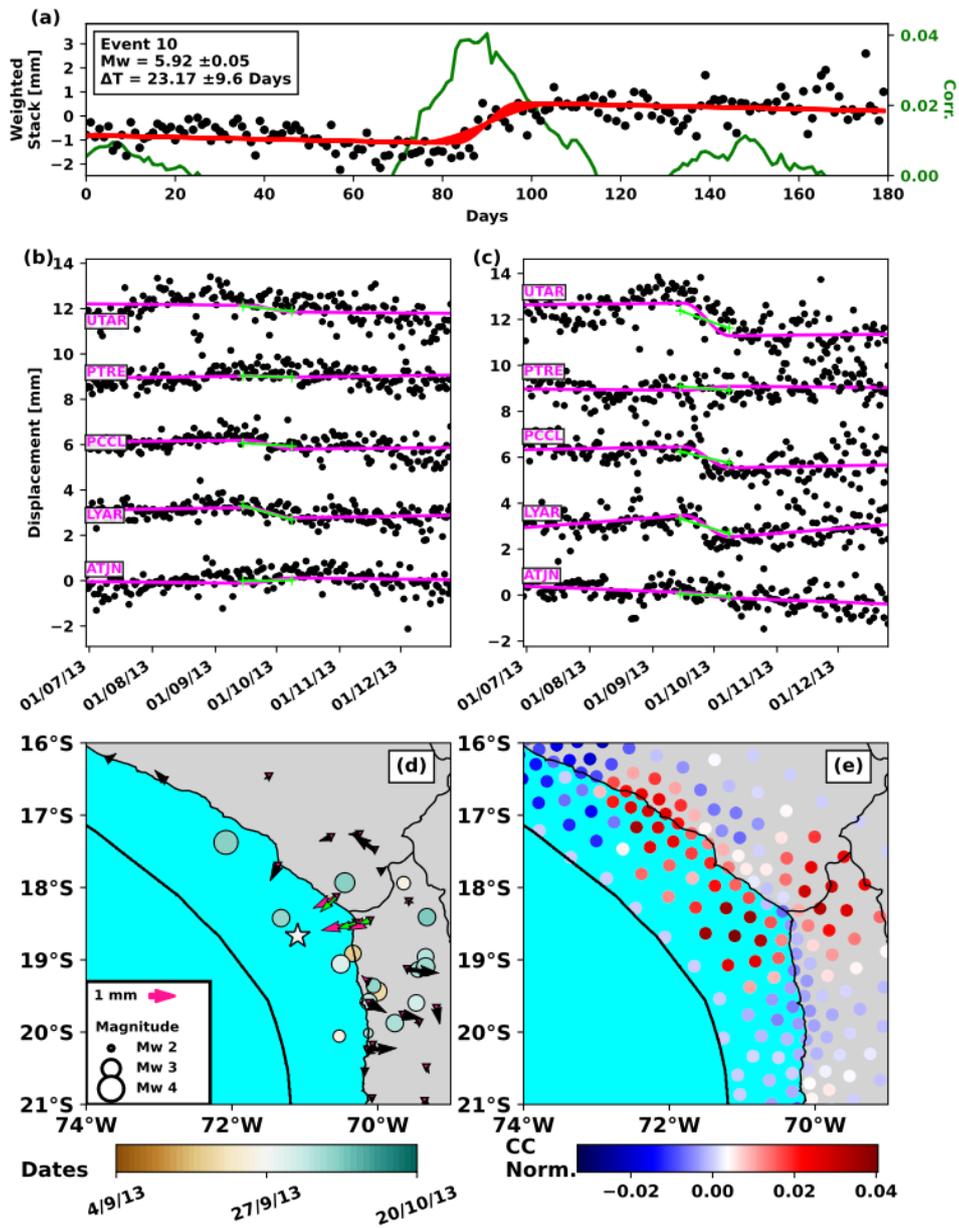


Figure 6 Same caption as Figure 5, but for event #10.

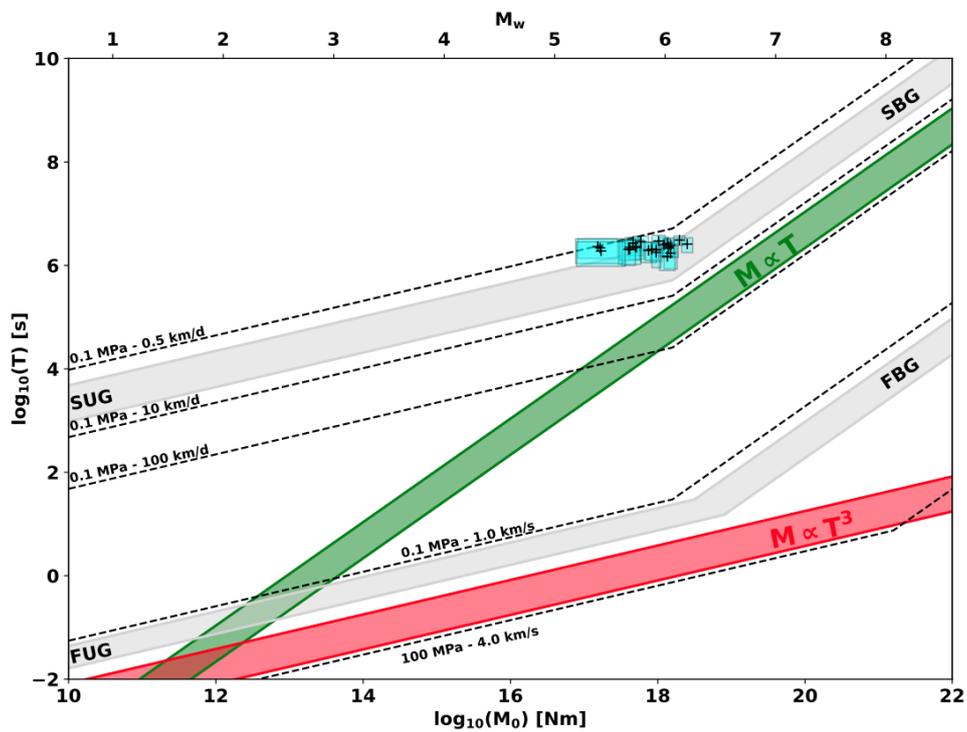


Figure 7 Seismic moment versus duration for our aseismic slip events following the scaling law proposed by [Gomberg et al. \(2016\)](#). Slow bounded/unbounded (SBG, SUG) and fast bounded/unbounded (FBG, FUG) regions are shown by light gray areas. Dashed lines are the theoretical relationship between moment and duration for a few selected stress-drop and rupture velocity values. The $M \propto T$ scaling is shown in green. The $M \propto T^3$ scaling is shown in red.

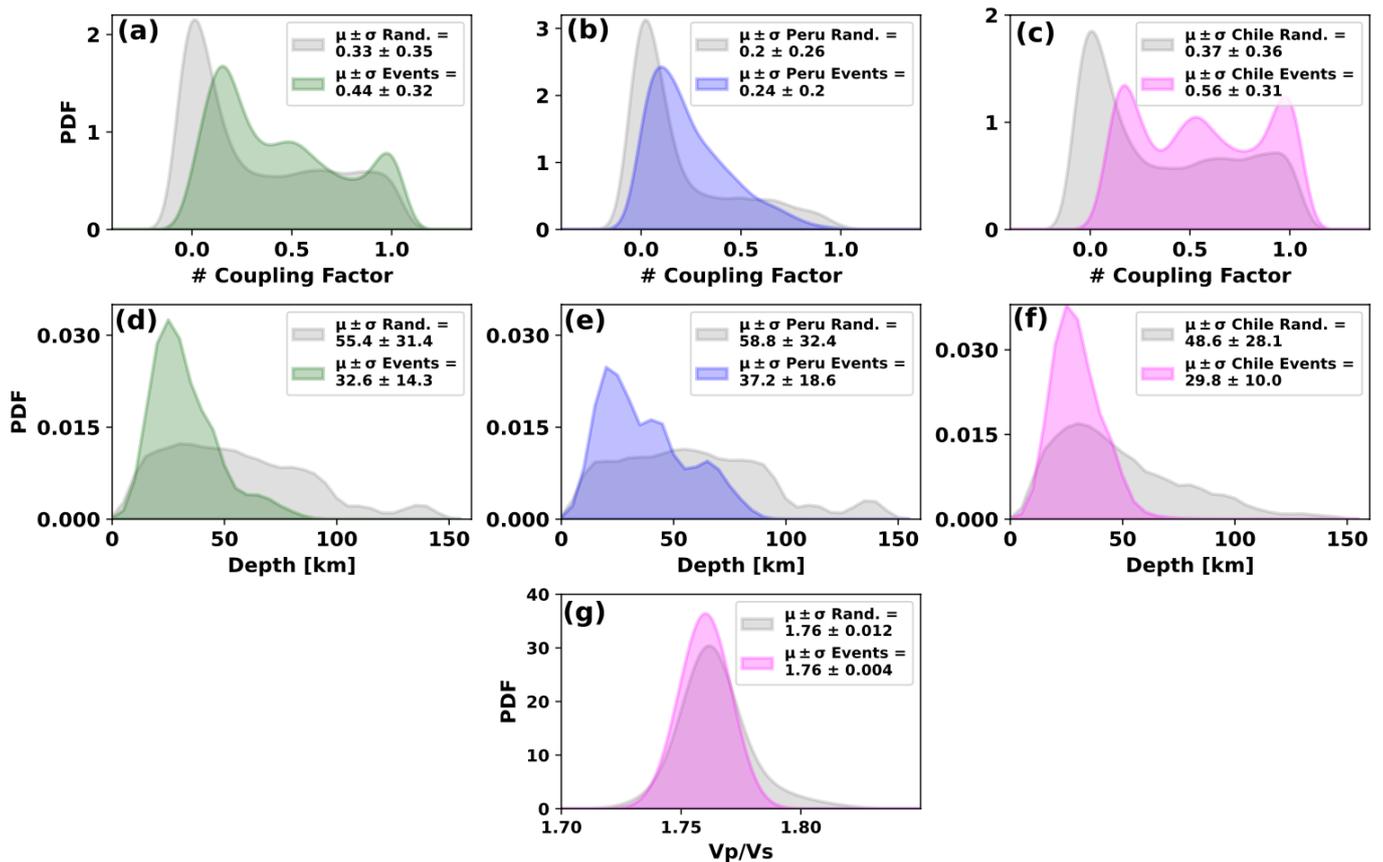


Figure 8 Coupling, depth, and Vp/Vs ratio of the detected aseismic slip events. (a) Probability Density Functions (PDF) of 1000 coupling models for 24 random picks (gray) and PDF of coupling where 24 aseismic slip events are detected (green), with respective mean (μ) and standard deviations (σ). (b) and (c) are the same as (a) for the Peru region only (gray: random, blue: SSEs) and northern Chile only (gray: random, magenta: events), respectively. (d) PDF of the depths of 24 random events (gray) and aseismic slip events detected in the region (green). (e) and (f) Same as (d) but for Peru (gray: random, blue: events) and Chile (gray, magenta) regions. (g) PDF of the Vp/Vs ratio for the Chilean region (gray, 17 random events), and detected aseismic events in Chile (magenta).

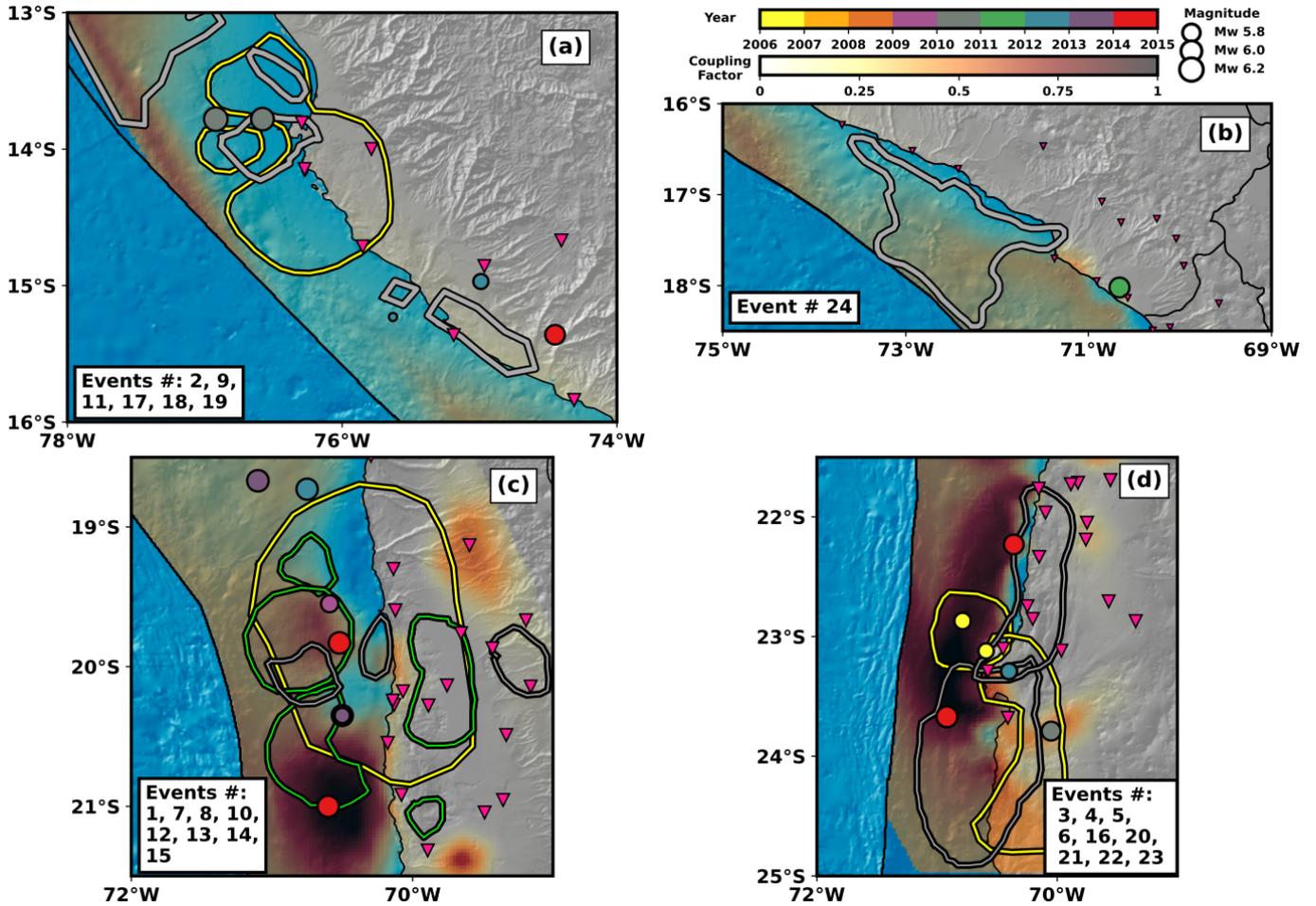


Figure 9 Zoom over a selection of regions of interest. Gray contours are instrumental ruptures. Yellow contours show reported afterslip. Our aseismic slip events are color-coded by time and scaled by magnitude. Background color shows our Bayesian inference of coupling. Inverted pink triangles are the GNSS stations used in this study. (a) Region struck by the Pisco (2007) and Nazca (1996) earthquakes. Our detections seem to cluster around asperities broken during earthquakes or afterslip regions. (b) Region struck by the Arequipa (2001) earthquake. (d) Region struck by the Iquique earthquake in 2014. Green contours show the pre-seismic slip reported by Socquet et al. (2017). Events occur around locked interseismic patches or low-coupled regions. (d) Region struck by the Antofagasta (1995) and Tocopilla (2007) earthquakes. Events surround broken asperities or locked interseismic patches, with a cluster beneath Mejillones Peninsula, potentially associated with earthquake afterslip. For citations of instrumental ruptures and afterslip, please refer to Figure 1

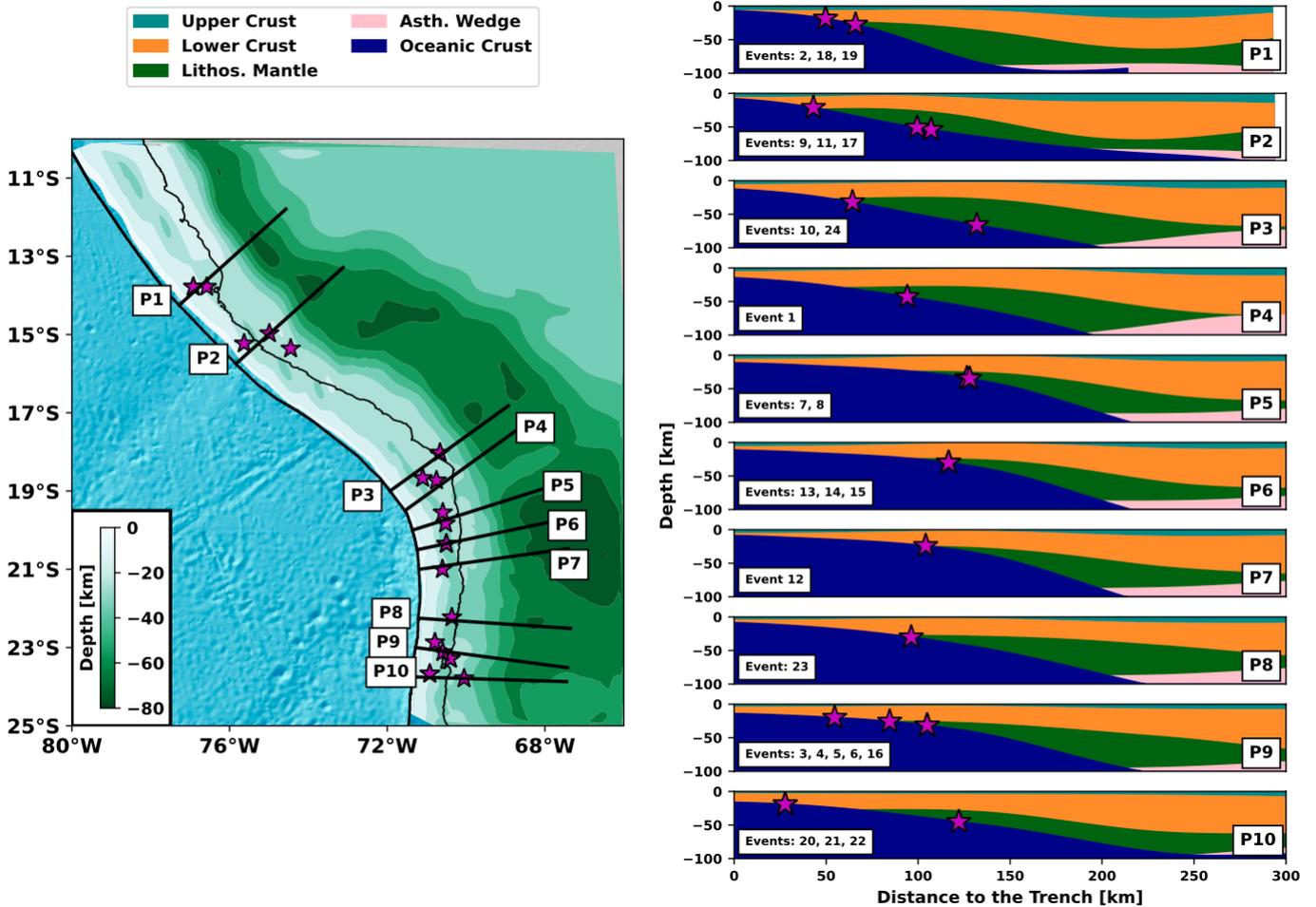


Figure 10 Map view of the depth of the continental Moho discontinuity from gravity-derived structural models by [Tassara and Echaurren \(2012\)](#). Magenta stars are the location of our 24 aseismic events. Black lines indicate the location of the profiles shown on the right. Colors indicate the structure at depth (upper and lower crusts, lithospheric mantle, asthenospheric wedge, and oceanic crust). White box indicates the id of events occurring along each profile.