

Constraints for the martian crust structure from Rayleigh waves ellipticity of large seismic events

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Key Points:

- Rayleigh waves ellipticity was measured at long periods (15-35 s) at the InSight landing site using large seismic events, including S1222a.
- A 4-layer crust model, including a shallow low-velocity layer, explains the local-scale observations.
- Low crustal velocities are derived for the InSight site, which may be due to high porosity or heavy alteration at local scale.

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Abstract

For the first time, we measured the ellipticity of direct Rayleigh waves at long periods (15 - 35 s) on Mars using the recordings of three large seismic martian events, including S1222a, the largest event recorded by the InSight mission. These measurements, together with P-to-s receiver functions and P-wave reflection times, were utilized for performing a joint inversion of the local crust structure at the InSight landing site. Our inversion results are compatible with previously reported intra-crustal discontinuities around 10 and 20 km depths, whereas the preferred resulting models show a strong discontinuity at ~ 37 km, which is interpreted as the crust-mantle interface. We propose the presence of a top shallow low-velocity layer of 2-3 km thickness. Compared to nearby regions, lower seismic wave velocities are derived for the local crust, thus suggesting a higher porosity or alteration of the whole local crust.

Plain Language Summary

As never before on Mars, we measured the characteristics of the seismic waves traveling along the martian surface that carry information about the crust structure at the InSight site. We combined these measurements with two other local-scale independent observations to derive a consolidated model for the crust underneath the InSight lander. Our results suggest a martian crust with 4 layers and, particularly, one thin layer of about 2 km thickness close to the surface. The crust-mantle discontinuity was found at ~ 37 km depth, where the sharpest change in seismic wave velocity is observed. Overall, the seismic wave velocities of the local martian crust at the InSight site are lower than those derived in other regions on Mars, which suggest a higher porosity or local alteration.

1 Introduction

The structure and properties of the martian crust are important for understanding the evolution of Mars and rocky planets. In this context, the InSight mission (Banerdt et al., 2020), which landed on Elysium Planitia in November 2018, aims to study the interior of Mars by means of one seismological station placed on the martian surface (Lognonné et al., 2019).

After more than three years of InSight data collection, the martian crust has been studied by a thorough analysis of the InSight seismological data. In particular, early analysis of P-to-s receiver functions (RFs) allowed to derive a shallow layer with low seismic velocities in the first 8-11 km of the martian crust at the landing site (Lognonné et al., 2020). Further investigation including more marsquakes suggests a layered martian crust beneath the InSight lander with two discontinuities at 8 ± 2 km and 20 ± 5 km, and a bimodal estimation of the crust thickness of either 20 ± 5 or 39 ± 8 km (Knapmeyer-Endrun et al., 2021). Lately, the analysis of further seismic events and the identification of other secondary phases (PPs, Sp and three new crustal multiples) favor a 3-layer crustal model (Kim et al., 2021; Joshi et al., 2023; Durán et al., 2022), for which the crust-mantle boundary would be located at ~ 40 km depth. The shallow discontinuity around 8 km depth has been confirmed by the detection of SsPp phases (J. Li, Beghein, Davis, et al., 2022) and SH-wave reflections (J. Li, Beghein, Wookey, et al., 2022). Besides, the identification of both PP and SS precursors indicates that the second discontinuity observed at the InSight site (at ~ 20 km depth) is also present around 4000 km away from the station and has been interpreted as a regional or global feature (J. Li, Beghein, McLennan, et al., 2022). On a global scale, Wiczorek et al. (2022) have estimated an average thickness of the martian crust between 30 and 72 km and a maximum permissible crustal density between 2850 and 3100 kg m⁻³. Further properties and characteristics of the martian crust at regional scales have been derived from the analysis of the InSight seismic data using different approaches (e.g., Kim, Banerdt, et al., 2022; Beghein et al., 2022; Kim, Stähler, et al., 2022; J. Li, Beghein, Lognonné, et al., 2022). Nevertheless, the el-

72 lipticity analysis of direct Rayleigh waves is still a pending task that could provide fur-
 73 ther constraints on the structure of the martian crust at the InSight landing site, as pro-
 74 posed by Panning et al. (2017).

75 Besides RFs, the analysis and inversion of the ellipticity of direct Rayleigh waves
 76 is a single-station technique that, when measured at long periods (>10 s), allows inves-
 77 tigating the local underground structure down to crustal depths (Tanimoto & Rivera,
 78 2008; Yano et al., 2009). In this sense, the ellipticity has been shown to be invariant to
 79 the seismic source, the wave propagation path or the epicentral distance and rather de-
 80 pend on the local structure directly beneath the receiver site (e.g., Ferreira & Woodhouse,
 81 2007). Even though this methodology was proposed in the 1970s, it has gained popu-
 82 larity only in recent years and with successful applications on Earth (e.g., Lin et al., 2012;
 83 G. Li et al., 2016; Berbellini et al., 2017), but it has not been applied on Mars so far due
 84 to the lack of observation of direct Rayleigh waves and, therefore, the characteristic el-
 85 lipticity of the InSight site at long periods remains unknown.

86 On May 4, 2022, the InSight lander recorded the S1222a event, which had a mo-
 87 ment magnitude estimated as M_W 4.7 and thus corresponds to the largest seismic event
 88 ever recorded on Mars (Kawamura et al., 2022). This event exhibits clear surface waves
 89 (both Love and Rayleigh) that can be used for studying lithosphere properties along the
 90 source-receiver path (e.g., Beghein et al., 2022; J. Li, Beghein, Lognonné, et al., 2022;
 91 Kim, Stähler, et al., 2022), but is also a unique opportunity to investigate the local crust
 92 structure through ellipticity analysis of Rayleigh waves. Similarly, the events S1000a and
 93 S1094b are two other good quality seismic events caused by the impact of meteorites on
 94 the martian surface (Posiolova et al., 2022), and for which Rayleigh waves have been ob-
 95 served (Kim, Banerdt, et al., 2022). These three large events are thus excellent candi-
 96 dates to perform a Rayleigh waves ellipticity (also referred to as ε) analysis.

97 The analysis of ε data can be used as an independent observation to characterize
 98 the martian crust at the InSight landing site. In fact, even though the inversion of ε data-
 99 only can be performed to map the mechanical properties (especially shear-wave veloc-
 100 ity, V_s) of the local subsurface, there is a great non-uniqueness of the solutions. In a sim-
 101 ilar way, the inversion of RFs is affected by the non-uniqueness phenomenon and, although
 102 this observation is highly sensitive to discontinuities, the addition of ε data can provide
 103 constraints to long-wavelength velocity features (Chong et al., 2016). As both data sets
 104 map the local structure at similar local scales, a joint inversion of ε and P-to-s RFs data
 105 can be performed. Similarly, the autocorrelation times obtained from the analysis of seis-
 106 mic ambient noise, such as those by Schimmel et al. (2021), are sensitive to the local crustal
 107 structure and therefore can be used as an independent data set to constrain the inver-
 108 sion.

109 In this work, we analyze the ellipticity and phase shift of the Rayleigh waves re-
 110 lated to the large seismic events S1000a, S1094b and especially S1222a, to obtain fur-
 111 ther constraints on the structure of the Martian crust. In particular, we perform a joint
 112 inversion of Rayleigh wave ellipticity, P-to-s RFs and autocorrelation times, aiming to
 113 obtain a consolidated local crustal model.

114 2 Data processing and methods

115 2.1 Rayleigh wave ellipticity

116 In order to retrieve the characteristic ε at the InSight site, we performed a time-
 117 domain analysis in a similar way as has been performed on Earth (e.g., Ferreira & Wood-
 118 house, 2007; Berbellini et al., 2016; Ferreira et al., 2020), following the steps shown in
 119 Fig. 1. First, the time windows where the Rayleigh waves arrive are roughly estimated
 120 by manual inspection, following the identification by Kim, Banerdt, et al. (2022) for S1000a
 121 and S1094b and the time estimates by the Marsquake Service (MQS) for S1222a (Kawa-

122 mura et al., 2022), and assuming they correspond to the fundamental mode (${}_0R_1$). These
 123 arrival time windows are shown in Fig. 1a, where the arrival of a suspected first over-
 124 tone of the Rayleigh wave (${}_1R_1$) for S1222a is also outlined, but it is not used in this work
 125 due to the low amplitude of the signal on the radial component.

126 The raw data (InSight Mars SEIS Data Service, 2019) are restituted to ground mo-
 127 tion (displacement) and transferred into the ZNE coordinates system. The data are then
 128 rotated into a ZRT coordinates system by using the back azimuth (BAZ) of each event.
 129 Particularly, as S1000a and S1094b have been identified as impacts with source locations
 130 confirmed by orbital images (Posiolova et al., 2022), their BAZs are well constrained at
 131 around 36° and 51° , respectively. For S1222a, even though there are estimates for the
 132 location of the source (e.g., Kawamura et al., 2022), there is no ground-truth location
 133 as for S1000a and S1094b. Accordingly, we applied an independent methodology to es-
 134 timate the BAZ of S1222a, based on the analysis of the Rayleigh waves. Our estimate
 135 is based on the fact that the Rayleigh waves propagate in the ZR plane and, thus, we
 136 look for the maximization of the ground-motion amplitudes on the Z and R components
 137 while minimizing the ground motion on the T component. Further details on this method-
 138 ology are provided in section S1 (Supp. Material). In the case of S1222a, our estimates
 139 for the BAZ are in the range 114° - 137° , with a median value of 129° , which is used in
 140 this work. Our BAZ is in the range of back azimuths obtained from the analysis of multi-
 141 orbit Rayleigh waves (Panning et al., 2022) and differs by $\sim 30^\circ$ from the BAZ derived
 142 from body waves by the MQS (Kawamura et al., 2022). This difference might be due to
 143 off-great-circle propagation of the Rayleigh waves, which is caused by strong lateral het-
 144 erogeneities along the source-receiver path (Evernden, 1953; Ji et al., 2005), particularly
 145 because they likely travel along the Martian dichotomy (Panning et al., 2022; Kim, Stähler,
 146 et al., 2022). In this region, clear topographic differences are observed between the north-
 147 ern and southern hemispheres. After rotating into the ZRT coordinates system, the Z
 148 and R components are bandpass-filtered, with a 30% bandwidth, around different cen-
 149 tral periods T_0 (Fig. 1b). In this way, every T_0 will be associated with one ellipticity value.

150 Subsequently, we obtained the optimal phase shift (ϕ) by maximizing the cross-correlation
 151 between the Z and R components. Examples of resulting ϕ and optimally shifted wave-
 152 forms are shown in Fig. 1c for the three events studied here. Once optimally shifted, we
 153 computed the cross-correlation factor $\nu(t)$, ranging from -1 to 1, between the shifted Z
 154 and R components (Fig. 1d). The normalized envelope ($\eta(t)$), also shown in Fig. 1d and
 155 given by the multiplication of the Z and R envelopes, is calculated in order to find the
 156 specific time window where the energy is maximized. In order to discard low-amplitude
 157 signals with high cross-correlation (or vice versa), a characteristic function $\chi(t) = \nu(t)\eta(t)$
 158 is calculated (Fig. 1e). The final ellipticity for the central period T_0 , computed as the
 159 R/Z quotient, corresponds to the average of the ellipticities in the time window where
 160 the characteristic function χ is larger than 0.8, as shown in Fig. 1f.

161 These processing steps are performed on the observed fundamental mode of the di-
 162 rect Rayleigh waves (${}_0R_1$) of the events S1000a, S1094b and S1222a. In order to com-
 163 bine these measurements, even though radial and azimuthal anisotropy may be present,
 164 we consider that these would affect ellipticity to a lesser extent than the underlying crustal
 165 structure. One ε value is obtained for each event and each central period T_0 . To get the
 166 characteristic ellipticity curve at the InSight landing site, only pairs of event- T_0 with ac-
 167 tual Rayleigh wave signal are considered, which is achieved by choosing data time win-
 168 dows where the signal-to-noise ratio (SNR) is larger than a given threshold in both Z
 169 and R components. We use $\text{SNR}_Z \geq 10$ and $\text{SNR}_R \geq 5$ (given the higher noise levels
 170 on the horizontals). The final Rayleigh ellipticity curve for the InSight site is then cal-
 171 culated as the smoothed median of the valid ε values, obtained by means of a Savitzky-
 172 Golay filter (Savitzky & Golay, 1964).

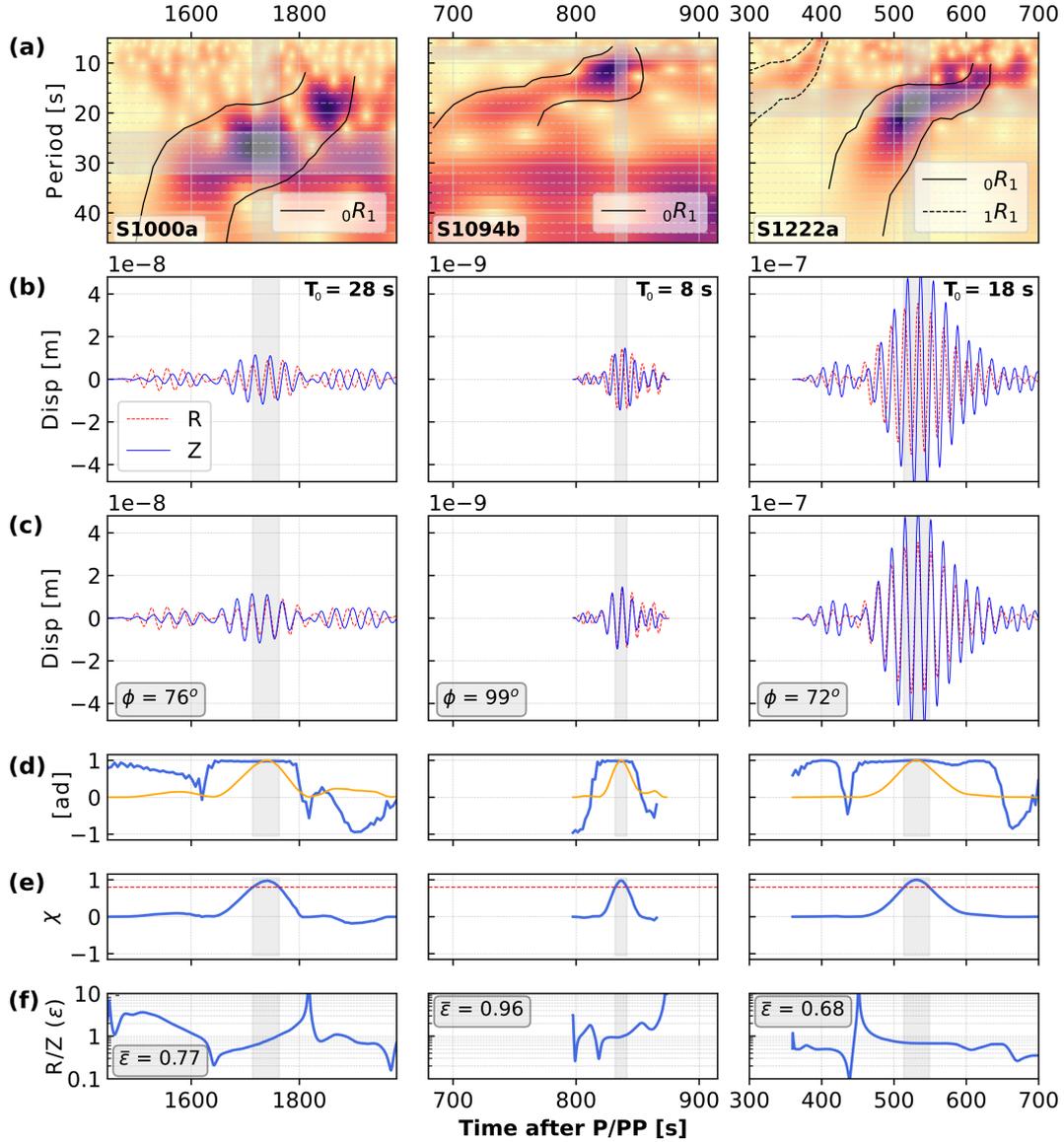


Figure 1. Extraction of Rayleigh wave ellipticity from large marsquakes. From left to right, columns correspond to events S1000a, S1094b and S1222a, respectively. From top to bottom: (a) S-transform spectrogram of the Z component, where the fundamental mode and first overtone of the Rayleigh waves are highlighted; (b) unshifted vertical (Z, solid blue) and radial (R, dashed red) components bandpass-filtered with a 30% bandwidth around a given central period T_0 , where the Rayleigh wave is clearly visible for each event. The spectral range shown in each case is also denoted by the horizontal gray area in (a); (c) shifted Z and R waveforms using the optimal phase shift for the given period; (d) cross-correlation factor between shifted Z and R components (blue) and the envelope of $ZR(t)$ (yellow); (e) characteristic function $\chi(t)$, where the red dashed line denotes a minimum threshold $\chi = 0.8$; (f) ellipticity $\epsilon(t)$. For each event and each central period, the average ellipticity ($\bar{\epsilon}$) is computed over the time window where $\chi \geq 0.8$ (gray box).

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2.2 Joint inversion of ellipticity, P-to-s RFs and autocorrelation times

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As the trade-off between the layer velocities and the depth of the discontinuities is a well-known feature of the ellipticity inversion, the incorporation of independent measurements such as receiver functions and autocorrelation times can help to further constrain the possible models (Chong et al., 2016). Because of this, we retrieved the 1D crustal structure underneath the InSight landing site by jointly inverting the ε measurements (this work) with the mean P-to-s RF by Joshi et al. (2023), which was obtained by stacking radial-component receiver functions for eight high-quality marsquakes mainly originating in the Cerberus Fossae region. Further details on the processing of the RF data can be found in Joshi et al. (2023). Additionally, we included the autocorrelation times obtained by Schimmel et al. (2021), which are understood as the two-way travel times of P waves bouncing back from the first (~ 6.15 s) and second (~ 10.6 s) strong discontinuity within the crust.

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For every model, the RF forward calculation is performed via a Python implementation of the code by Shibutani et al. (1996), using a fixed ray parameter $p^* = 6.9$ s/deg, suitable for events in the Cerberus Fossae region (Joshi et al., 2023). In the RF, four peaks at around 0, 2.4, 4.8 and 7.2 s are observed, which are related to the direct P phase and three converted phases P_{1s} , P_{2s} , P_{3s} , where the subscript indicates the interface where each phase is generated. The synthetic ellipticity curves are obtained via the *gpell* toolbox included in the Geopsy package (Wathelet, 2005). Lastly, the autocorrelation times are computed assuming two-way vertical travel paths of the P waves.

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The inversion scheme utilized is the Neighbourhood Algorithm (NA), first introduced by Sambridge (1999), that has been widely used for multiple geophysical applications and, particularly, for subsurface characterization (e.g., Hobiger et al., 2013; Ferreira et al., 2020; Carrasco et al., 2022). We applied a modified version of the NA, where the parameters space can be constrained by physical conditions and prior information (Wathelet, 2008). Among other constraints, we used the main outcomes of the transdimensional inversion of P-to-s RFs by Joshi et al. (2023), who found confirmation for a 3-layer crust model and provide constraints for the depths of the largest discontinuities. The NA is based on a fully non-linear, self-adaptive Monte Carlo approach, which efficiently explores the parameters space in order to find the model \mathbf{m} that minimizes the misfit function Ψ between the observed and theoretical data. For each explored model \mathbf{m} , Ψ is defined as:

$$\Psi(\mathbf{m}) = W_e \psi_e(\mathbf{m}) + W_{rf} \psi_{rf}(\mathbf{m}) + W_a \psi_a(\mathbf{m}), \quad (1)$$

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where the weights of the ellipticity (W_e), the receiver functions (W_{rf}) and autocorrelation times (W_a) data sets were fixed to 0.45, 0.45 and 0.1, respectively, to explain mainly the ε and RFs measurements, with a small constraint given by the autocorrelation times. The $\psi_i(\mathbf{m})$ terms correspond to the individual misfits of each data set, given by the fit between the observed and synthetic calculations for each model \mathbf{m} . Further details can be found in section S6 (Supp. Material).

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3 Results and discussion

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3.1 Ellipticity and phase-shift observations

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The resulting ε and ϕ values are shown in Fig. 2a and b, respectively. Besides the limited number of events analyzed, there is a lack of valid ellipticity measurements due to the low-amplitude energy on the radial component.

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The valid ε values, i.e., those for which the imposed SNR criteria are fulfilled, define an ellipticity curve between 15 and 35 s and two points at 8 and 9 s (Fig. 2). For

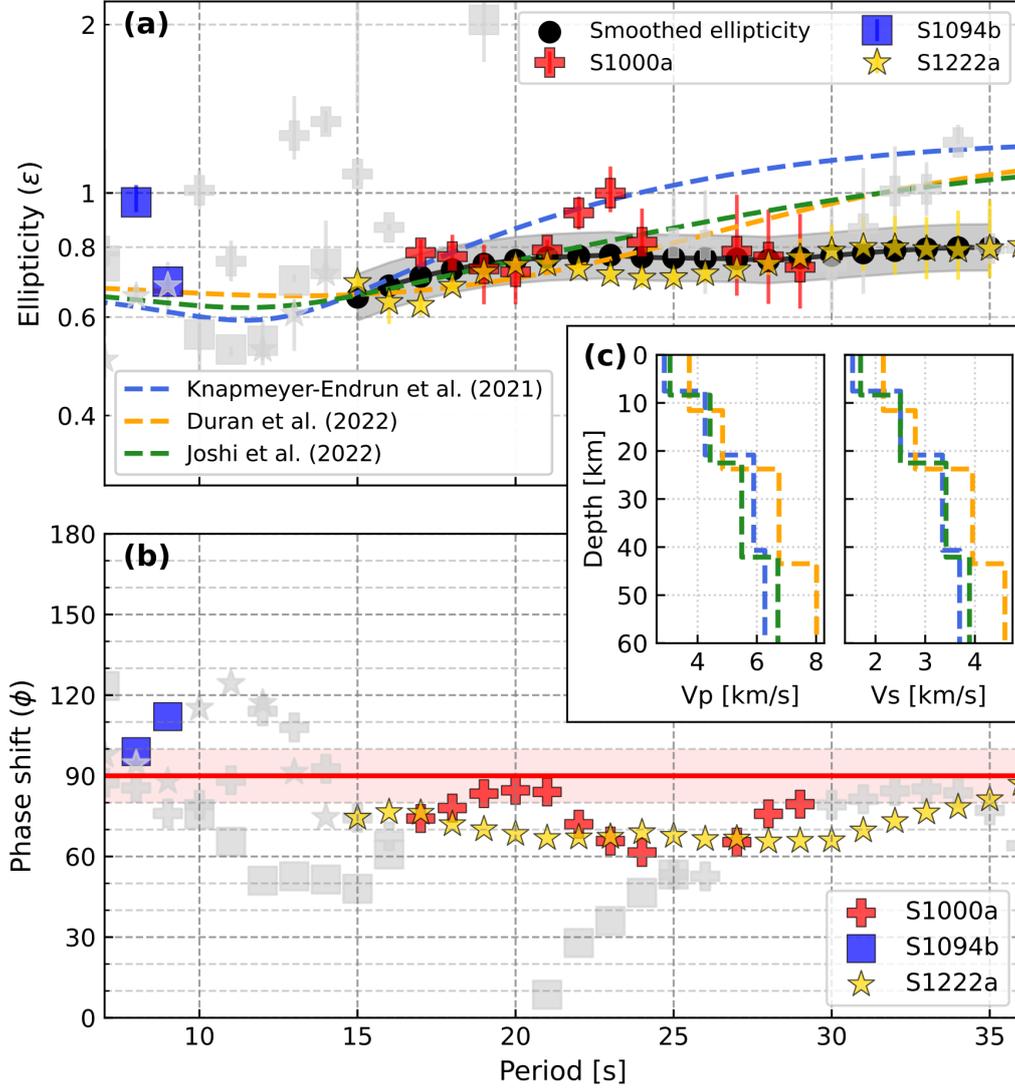


Figure 2. Rayleigh wave ellipticity and phase shift results. (a) Rayleigh wave ellipticity (ϵ) retrieved at different central periods T_0 for the three large events S1000a (red cross), S1094b (blue square) and S1222a (yellow star). Colored markers correspond to periods where the SNR criteria are fulfilled. Background light gray markers did not meet the minimum SNR threshold. The final ellipticity curve for the InSight site is denoted by the black circles, and a 15% uncertainty range has been imposed (gray area). Synthetic ellipticity curves are illustrated for characteristic crustal models from Knapmeyer-Endrun et al. (2021), Durán et al. (2022) and Joshi et al. (2023). (b) Phase shift ϕ between Z and R components of ${}_0R_1$ for the same events and with the same color code as in (a). The light red area corresponds to ϕ in the range $90 \pm 10^\circ$. (c) V_p and V_s models for the InSight crust used for synthetic calculation of Rayleigh wave ellipticity in (a).

219 the purpose of this paper, we will focus on the period range between 15 and 35 s. In this
 220 range, the horizontal contribution of the tilt induced by Rayleigh waves can be neglected.
 221 Due to the lack of further observations, a common 15% uncertainty is utilized, which prop-
 222 erly covers our observations. The resulting ellipticity curve is rather flat, especially be-
 223 tween 20 and 35 s, with a slight decrease towards shorter periods. No strong peaks or
 224 troughs are observed in this period range.

225 On Earth, the ellipticity curve of one specific site is obtained from the recordings
 226 of a large number of earthquakes. In our work, even though three events were initially
 227 investigated, the ellipticity curve is mainly ruled by S1222a and therefore the lack of ob-
 228 servations can be disadvantageous for further analysis and interpretation. It has been
 229 shown that ellipticity data can have a large spread (Berbellini et al., 2016; Attanayake
 230 et al., 2017), which has been associated with complex Rayleigh wave propagation effects
 231 (Sexton et al., 1977; Pedersen et al., 2015).

232 For a better assessment, the same ellipticity analysis was performed for one sta-
 233 tion on Earth (station BQ.DREG in Germany, see section S2 in Supp. Material). These
 234 measurements suggest that well-recorded events can show some variation with respect
 235 to the characteristic smoothed ellipticity curve of the site, but the broad shape of the
 236 curve remains similar. In particular, some events show different slopes and trends at spe-
 237 cific period ranges, with respect to what is observed for the median ellipticity (Fig. S2
 238 and S3), but the general trend is the same. Thus, we conclude that small oscillations on
 239 the ellipticity curve should not be over-interpreted. It is also worth noting that the un-
 240 certainty range of these Earth observations is about 20%, close to the error range uti-
 241 lized here. Besides, in the InSight case, the ${}_0R_1$ of the event S1222a has a high SNR in
 242 both Z and R components, so the final smoothed ε is expected to be a good represen-
 243 tation of the local structure.

244 We calculated synthetic ellipticity values for representative three-layer crust mod-
 245 els (see section S4), based on the inversion results by Knapmeyer-Endrun et al. (2021),
 246 Durán et al. (2022) and Joshi et al. (2023). These models were derived from the inver-
 247 sion of P-to-s RFs from different sets of martian seismic events and, in the case of Joshi
 248 et al. (2023), combined with apparent S-wave velocities, whereas Durán et al. (2022) also
 249 included teleseismic data. Even though the corresponding ellipticity curves have a good
 250 fitting at periods between 15 and 21 s, all these models overestimate ε at longer peri-
 251 ods (Fig. 2a). Furthermore, the calculated times of the P-wave reflections for these mod-
 252 els (Table S1) do not fit the times from the seismic noise autocorrelations analysis (Schim-
 253 mel et al., 2021), which further encourages a joint inversion of the three data sets.

254 Besides the ellipticity values, the resulting Z-R phase shifts are another interest-
 255 ing observation. Mainly ruled by the S1222a event, our results suggest ϕ values between
 256 60° and 70° (Fig. 2b), which are persistently lower than the theoretical phase shift for
 257 Rayleigh waves under isotropic and homogeneous conditions (90°). Considering the pos-
 258 sible uncertainties, we tested different values of back azimuth for S1222a and the result-
 259 ing ϕ are still persistently lower than 90° (see Fig. S5).

260 On Earth, variations of ϕ are usually in the range 80 - 100° , but larger deviations
 261 up to $\pm 30^\circ$ have been also found and have been proposed to be related to small-scale
 262 heterogeneities (e.g., Ferreira & Woodhouse, 2007). This phenomenon has been less well-
 263 studied than ellipticity, so the interpretation of this observation is challenging. Accord-
 264 ing to our observations on Earth for station DREG, even though a typical 90° phase shift
 265 is measured for most of the earthquakes, low ϕ (60 - 75°) were measured for some earth-
 266 quakes at specific period ranges (see Fig. S4). These observations suggest that anomalous
 267 ϕ might be related to the Rayleigh-waves path, or even the source, rather than the
 268 characteristics of the receiver site. However, Tanuma & Man (2008) also proposed that
 269 variations of ϕ with respect to the isotropic theoretical value depend on the perturba-
 270 tion of the initial stress conditions of the medium. Under these considerations, the low

271 ϕ measured at the InSight landing site might be due to a combination of path or source
 272 effects with the initial stress conditions at this site, but not directly related to the me-
 273 chanical characteristics of the crust at the InSight landing site. In any case, low ϕ val-
 274 ues have previously been observed on Earth to occur in some cases. Further analysis of
 275 this phenomenon is beyond the scope of this work.

276 **3.2 Local crust structure at the InSight landing site**

277 In order to derive a consolidated crustal structure for the InSight landing site, the
 278 joint-inversion of the ellipticity measurements (this work), P-to-s RFs (Joshi et al., 2023)
 279 and P-wave reflection times (Schimmel et al., 2021) was performed. In this regard, it is
 280 worth noting that the inversion of ellipticity data alone does not properly constrain the
 281 local crustal structure. In fact, the ellipticity inversion using different settings between
 282 one and four layers over a half-space can equally well explain the ε observations (Fig.
 283 S6).

284 The initial joint-inversion results suggest that these local measurements can be well
 285 explained by a three-layer model, with an overall good fitting of the three data sets (see
 286 Fig. S8). However, the amplitude of the P-phase (first peak on the RFs), which is not
 287 normalized as it provides information on the near-surface velocities (e.g., Ammon, 1991),
 288 is clearly overestimated. As shown by synthetic forward modeling (Fig. S9), lower am-
 289 plitudes for this arrival can be achieved by having lower velocities in the first layer. Be-
 290 cause of this, we performed a second inversion where we allowed a shallow low-velocity
 291 layer (see Table S3 for details on the parameter space).

292 The resulting models from this second inversion are preferred as they explain the
 293 three independent data sets reasonably well (Fig. 3a-d), and especially the amplitude
 294 of the direct P-phase arrival is well retrieved. The predicted ellipticity curves have a bet-
 295 ter fitting at long periods (between 25 and 35 s) than between 15 and 25 s, where the
 296 short-scale curvature cannot be fully modeled (Fig. 3a). This slight discrepancy may be
 297 related to the lack of data for proper averaging and smoothing of the ellipticity curve,
 298 as can be done on Earth (e.g., Fig. S2). As discussed above, we consider the long-wavelength
 299 shape of the ellipticity curve to be the most reliable observation and, therefore, our in-
 300 version results provide an acceptable prediction of it. On the other hand, the autocor-
 301 relation times are well predicted within their uncertainty range (Fig. 3b).

302 The arrival times and amplitudes of the main phases on the RF are also well pre-
 303 dicted by our best-resulting models (Fig. 3b), but some discrepancies are observed for
 304 the amplitude of the second peak (~ 2.2 s). This mismatch can be related to a larger
 305 amplitude variability of the P_{1s} phase (see Fig. 3 in Joshi et al., 2023). In any case, the
 306 two arrival times and the amplitudes of the major peaks are well retrieved.

307 Our preferred models show that the crust at the InSight landing site can be un-
 308 derstood as a four-layer crust, including a shallow low-velocity layer (sLVL) close to the
 309 surface. The sLVL presents low velocities between 1.5 and 1.8 km/s and extends down
 310 to 2-3 km depth. Its origin is unclear, but previous works on the shallowest structure
 311 of Mars have proposed the presence of sedimentary and heavily fractured basalt layers
 312 in the first hundreds of meters (Hobiger et al., 2021; Carrasco et al., 2022), where shear-
 313 wave velocities below 1 km/s were derived. Furthermore, Pan et al. (2020) suggested the
 314 presence of a sedimentary layer extending from 200-300 m up to kilometers in depth, whereas
 315 Warner et al. (2022) showed evidence for heavily fractured basalt units in the surround-
 316 ing area. Thus, we interpret this shallow discontinuity to be the end of this large-scale
 317 either sedimentary or heavily fractured basalt layer. A similar observation of such a shal-
 318 low layer, based on the analysis of high-frequency RFs of S1222a, has been also recently
 319 proposed by Shi et al. (2022).

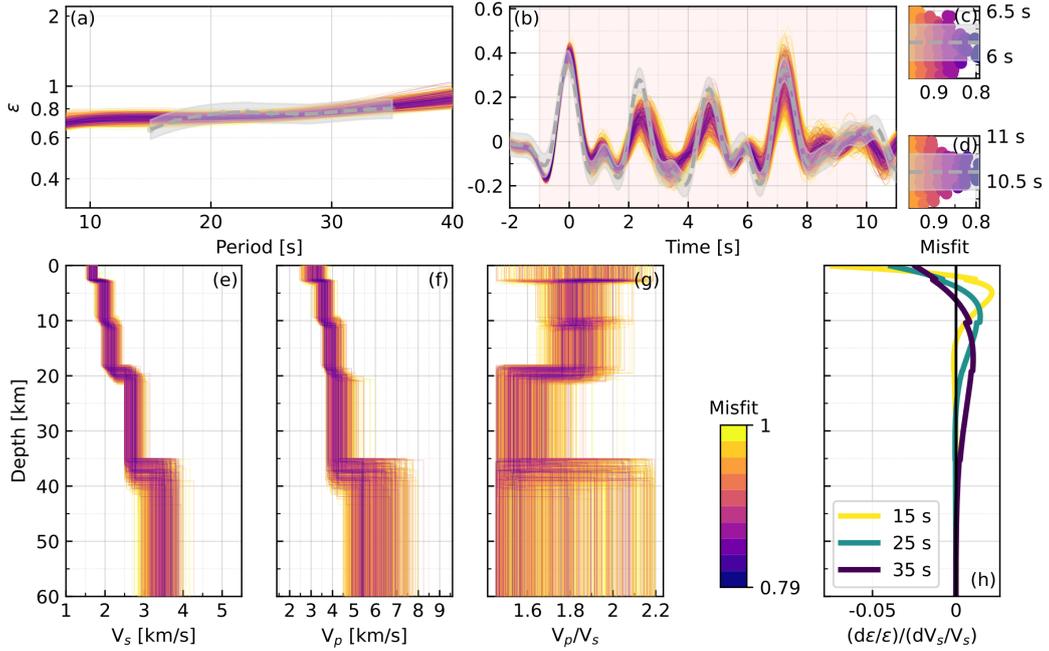


Figure 3. Joint-inversion results and resulting local martian crustal models. The measurements and the synthetic modeling of the 200 best-fitting crustal models are shown for (a) Rayleigh wave ellipticity, (b) P-to-s RFs from Joshi et al. (2023), and (c) first and (d) second P-wave reflection times by Schimmel et al. (2021). The light gray dashed line and area correspond to the measured data and their uncertainty, respectively, whereas the synthetic data are colored by their misfit. The corresponding crustal models are shown in terms of their (e) S-wave velocity V_s , (f) P-wave velocity V_p and (g) V_p/V_s , as a function of depth. (h) Depth sensitivity kernels for a characteristic resulting model calculated at three different periods: 15 s (yellow), 25 s (green) and 35 s (dark blue). Light red box in (b) shows the time window where the data was fitted.

320 Below this shallow layer, our models have a similar structure and the discontinu-
 321 ities at ~ 10 and ~ 20 km depth are in good agreement with previously reported mod-
 322 els for the local crust (e.g., Knapmeyer-Endrun et al., 2021; Durán et al., 2022; Joshi et
 323 al., 2023). A third discontinuity is derived at around 37 km depth and it is interpreted
 324 as the crust-mantle boundary, where the sharpest velocity contrast is observed and where
 325 the ellipticity data has low sensitivity (Fig. 3h). Nevertheless, the ellipticity data can
 326 properly constrain the velocities at shallow depths and, given the identification of the
 327 P_{3s} phase on the RF, the joint-inversion can provide reliable depths for this discontinu-
 328 ity.

329 Our preferred models present overall P-wave and S-wave velocities lower than 5 and
 330 3 km/s, respectively, for the whole crust. Particularly, the shear-wave velocities are lower
 331 than those in nearby regions on Mars. For example, Kim, Banerdt, et al. (2022) found
 332 V_s around 3.2 km/s, between 5 and 30 km depth, for a path between Amazonis Plani-
 333 tia and the InSight site, in the lowlands. Similarly, V_s between 3.05-3.17 km/s were de-
 334 rived for the upper crust (first 20 km) of the path between S1222a and the lander (~ 2190
 335 km distance), along the dichotomy region (J. Li, Beghein, Lognonné, et al., 2022). At
 336 a local scale, our proposed shear-wave velocities are also lower than those previously re-
 337 ported for the InSight site (Knapmeyer-Endrun et al., 2021; Durán et al., 2022; Joshi
 338 et al., 2023). Therefore, following the crustal layering interpretation by Wieczorek et al.
 339 (2022), and considering that porosity plays an important role in decreasing the veloc-
 340 ity of the seismic waves (e.g., Heap, 2019), these lower velocities might be related to an
 341 even larger porosity or stronger alteration of the whole crust at the InSight site.

342 4 Conclusions

343 For the first time, the ellipticity of the fundamental mode of direct Rayleigh waves
 344 was measured on Mars for long periods, especially thanks to S1222a: the largest event
 345 recorded on Mars so far. The obtained ellipticity curve, between 15 and 35 s, is a good
 346 representation of the characteristic ellipticity at the InSight landing site. Anomalously
 347 low phase shifts were measured, which seem to be unrelated to the local crust structure.
 348 In order to consolidate the local crustal model, we performed a joint inversion of local-
 349 scale observations (ellipticity, P-to-s RFs and P-wave reflection times). Our results sug-
 350 gest a four-layer crustal model, with a shallow low-velocity layer of about 2-3 km thick-
 351 ness at the InSight landing site. Other discontinuities around 10, 20 and 37 km depth
 352 correlate well with previously reported models. We propose lower P- and S-wave veloc-
 353 ities for the corresponding crustal layers, which might be related to a higher porosity or
 354 alteration of the martian crust at the InSight landing site.

355 Acknowledgments

356 The authors acknowledge NASA, CNES, their partner agencies and institutions (UKSA,
 357 SSO, DLR, JPL, IPGP-CNRS, ETHZ, IC and MPS-MPG), and the flight operations team
 358 at JPL, SISMOC, MSDS, IRIS-DMC and PDS for providing SEED SEIS data. Thanks
 359 to Dr. J. Li for his thoughtful comments about the manuscript. French co-authors ac-
 360 knowledge the French Space Agency CNES and ANR (ANR-19-CE31-0008). PL, ES, ZX
 361 acknowledge IDEX Paris Cité (ANR-18-IDEX-0001). The research was carried out in part
 362 at the Jet Propulsion Laboratory, California Institute of Technology, under a contract
 363 with the National Aeronautics and Space Administration (80NM0018D0004). Open ac-
 364 cess funding enabled and organized by Projekt DEAL. This is InSight contribution ICN
 365 299.

Open Research

The event information of the martian seismic events can be found in the InSight seismic event catalog version 13 (InSight Marsquake Service, 2023). The waveform data and station metadata are available from IGP’s MSDC as well as from IRIS MDC (InSight Mars SEIS Data Service, 2019). Seismic data for station DREG (BQ network, 2016) are publicly available via EIDA (<http://eida.gfz-potsdam.de/webdc3/>). The GCMT catalog is available here <https://www.globalcmt.org/CMTsearch.html>, whereas the events analyzed for the Earth case and the codes utilized for the processing of the events can be found under https://github.com/scarrascom/rayleigh_ellipticity.

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