

Constraints for the Martian crustal structure from Rayleigh waves ellipticity of large seismic events

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Key Points:

- Rayleigh waves ellipticity was measured between periods 15-35 s at the InSight landing site using large seismic events, including S1222a.
- A 4-layer crust, including a shallow low-velocity layer, is required to explain the ellipticity, receiver functions and P-wave lag times.
- Low crustal velocities are derived for the InSight site, which may be due to high porosity or heavy alteration at local scale.

Abstract

For the first time, we measured the ellipticity of direct Rayleigh waves at intermediate periods (15 - 35 s) on Mars using the recordings of three large seismic Martian events, including S1222a, the largest event recorded by the InSight mission. These measurements, together with P-to-s receiver functions and P-wave reflection times, were utilized for performing a joint inversion of the local crustal structure at the InSight landing site. Our inversion results are compatible with previously reported intra-crustal discontinuities around 10 and 20 km depths, whereas the preferred models show a strong discontinuity at ~ 37 km, which is interpreted as the crust-mantle interface. Additionally, we support the presence of a shallow low-velocity layer of 2-3 km thickness. Compared to nearby regions, lower seismic wave velocities are derived for the crust, suggesting a higher porosity or alteration of the whole local crust.

Plain Language Summary

As never before on Mars, we measured the characteristics of seismic waves traveling along the Martian surface that carry information about the crustal structure at the InSight site. We combined these measurements with two other local-scale independent observations to derive a consolidated model for the crust underneath the InSight lander. Our results suggest a Martian crust with 4 layers and, particularly, one thin layer of about 2 km thickness close to the surface. The crust-mantle discontinuity was found at ~ 37 km depth, where the sharpest change in seismic wave velocity is observed. Overall, the seismic wave velocities of the local Martian crust at the InSight site are lower than those derived in other regions on Mars, which suggests a higher porosity or local alteration.

1 Introduction

The structure and properties of the Martian crust are important for understanding the evolution of Mars and rocky planets. In this context, the InSight mission (Banerdt et al., 2020), which landed on Elysium Planitia in November 2018, aims to study the interior of Mars by means of one seismological station placed on the Martian surface (Lognonné et al., 2019).

After more than three years of data collection, the Martian crust has been studied by a thorough analysis of the InSight seismological data. Locally at the InSight landing site, early analysis of P-to-s receiver functions (RFs) allowed to derive a shallow layer with low seismic velocities in the first 8-11 km of the Martian crust (Lognonné et al., 2020). Further investigation including more marsquakes suggests a layered local Martian crust with two discontinuities at 8 ± 2 km and 20 ± 5 km, and a bimodal estimation of the crust thickness of either 20 ± 5 or 39 ± 8 km (Knapmeyer-Endrun et al., 2021). Subsequently, the analysis of further seismic events and the identification of other secondary phases (PPs, Sp and three new crustal multiples) favor a 3-layer crustal model (Kim et al., 2021; Joshi et al., 2023), for which the crust-mantle boundary would be located at ~ 40 km depth. The shallow discontinuity around 8 km depth has been confirmed by the detection of SsPp phases (J. Li, Beghein, Davis, et al., 2022) and SH-wave reflections (J. Li, Beghein, Wookey, et al., 2022). Recently, Shi et al. (2023) proposed a shallow discontinuity at around 2 km depth, based on the analysis of high-frequency RFs. Further properties and characteristics of the Martian crust at regional scales (e.g., Kim, Banerdt, et al., 2022; Beghein et al., 2022; Kim, Stähler, et al., 2022; J. Li, Beghein, Lognonné, et al., 2022), as well as at global scale and away from the lander (e.g. J. Li, Beghein, McLennan, et al., 2022; Wiczorek et al., 2022; Durán et al., 2022), have been derived from the analysis of the InSight seismic data using different approaches. Table S1 in the Supplementary Material details some representative models for the Martian crust.

71 Besides RFs, the analysis and inversion of the ellipticity of direct Rayleigh waves
 72 is a single-station technique that, when measured at intermediate and long periods (>10
 73 s), allows investigating the local underground structure down to crustal depths (Tani-
 74 moto & Rivera, 2008; Yano et al., 2009). The ellipticity, understood as the ratio between
 75 the radial (R) and the vertical (Z) ground-motion of the Rayleigh wave (Fig. 1a), has
 76 been shown to be invariant to the seismic source, the wave propagation path or the epi-
 77 central distance and rather depend on the local structure directly beneath the receiver
 78 (e.g., Ferreira & Woodhouse, 2007). Even though this methodology was proposed in the
 79 1970s, it has gained popularity only in recent years and with successful applications on
 80 Earth (e.g., Lin et al., 2012; G. Li et al., 2016; Berbellini et al., 2017), but it has not been
 81 applied on Mars so far due to the lack of observation of direct Rayleigh waves. Hence,
 82 the ellipticity analysis of direct Rayleigh waves is still a pending task that could provide
 83 further constraints on the structure of the Martian crust at the InSight landing site, as
 84 proposed by Panning et al. (2017).

85 On May 4, 2022, the InSight lander recorded the S1222a event, which had a mo-
 86 ment magnitude estimated as M_W 4.7 and thus corresponds to the largest seismic event
 87 ever recorded on Mars (Kawamura et al., 2022). This event exhibits clear surface waves
 88 (both Love and Rayleigh) that can be used for studying lithospheric properties along the
 89 source-receiver path (e.g., Beghein et al., 2022; J. Li, Beghein, Lognonné, et al., 2022;
 90 Kim, Stähler, et al., 2022; Xu et al., 2023), but is also a unique opportunity to investi-
 91 gate the local crustal structure through ellipticity analysis of Rayleigh waves. Similarly,
 92 the events S1000a and S1094b are two other good-quality seismic events, which are re-
 93 lated to the impact of meteorites on the Martian surface, 126° and 58° away from the
 94 InSight lander, and had magnitudes estimated as M_W 4.1 and 4.0, respectively (Posi-
 95 olova et al., 2022). For these two events, Rayleigh waves have also been observed (Kim,
 96 Banerdt, et al., 2022). These three large events, whose source locations are shown in Fig.
 97 1b, are thus excellent candidates to perform a Rayleigh wave ellipticity (also referred to
 98 as ε) analysis.

99 The analysis of ε data can be used as an independent observation to characterize
 100 the Martian crust at the InSight landing site. The inversion of RFs is affected by the non-
 101 uniqueness phenomenon and, although these data are highly sensitive to discontinuities,
 102 the addition of ε data can provide constraints on long-wavelength velocity features (Chong
 103 et al., 2016). As both data sets map the local structure at similar local scales, a joint
 104 inversion of ε and P-to-s RFs can be performed. Similarly, the autocorrelation times (here-
 105 after T_a) obtained from the analysis of seismic ambient noise, such as those by Schim-
 106 mel et al. (2021), are sensitive to the local crustal structure and therefore can be used
 107 as an independent data set to constrain the inversion. In fact, the predicted T_a for the
 108 crustal models in Table S1 differs from the measurements (Table S3) whereas, as shown
 109 in Fig. 2a, the corresponding predicted ellipticity curves fail to match the measured ε .
 110 These mismatches further encourage a joint inversion of the three data sets.

111 In this work, we analyze the ellipticity and phase shift of the Rayleigh waves re-
 112 lated to the large seismic events S1000a, S1094b and especially S1222a, to obtain fur-
 113 ther constraints on the structure of the Martian crust. In particular, we perform a joint
 114 inversion of Rayleigh wave ellipticity, P-to-s RFs and T_a , aiming to obtain a consolidated
 115 local crustal model.

116 2 Data processing and methods

117 2.1 Rayleigh wave ellipticity

118 To retrieve the characteristic ε at the InSight site, we performed a time-domain anal-
 119 ysis in a similar way as has been performed on Earth (e.g., Ferreira & Woodhouse, 2007;
 120 Berbellini et al., 2016; Ferreira et al., 2020), following the steps shown in Fig. 1c-f. First,

121 the time windows where the Rayleigh waves arrive are roughly estimated by manual inspection, following the identification by Kim, Banerdt, et al. (2022) for S1000a and S1094b
 122 and the time estimates by the Marsquake Service (MQS) for S1222a (Kawamura et al.,
 123 2022), and assuming they correspond to the fundamental mode (${}_0R_1$). These arrival time
 124 windows are shown in Fig. 1c.
 125

126 The raw data (InSight Mars SEIS Data Service, 2019) have been deglitched using
 127 the UCLA method (Scholz et al., 2020), restituted to ground motion (displacement) and
 128 transferred into the ZNE coordinates system. The data are then rotated into a ZRT co-
 129 ordinates system by using the back azimuth (BAZ) of each event. As S1000a and S1094b
 130 have been identified as impacts with source locations confirmed by orbital images (Po-
 131 siolova et al., 2022), their BAZs are well constrained at around 36° and 51° , respectively.
 132 For S1222a, even though there are estimates for the location of the source (e.g., Kawa-
 133 mura et al., 2022), there is no ground-truth location as for S1000a and S1094b.

134 We applied an independent methodology to estimate the BAZ of S1222a, based on
 135 the fact that the Rayleigh waves propagate in the ZR plane. Thus, we look for the max-
 136 imization of the ground-motion amplitudes on the Z and R components while minimiz-
 137 ing the ground motion on the T component. Further details on this methodology are pro-
 138 vided in section S2 (Supp. Material). For S1222a, our estimated BAZs are in the range
 139 114° - 137° , with a median value of 129° , which is used in this work. This BAZ is in the
 140 range obtained from the analysis of multi-orbit Rayleigh waves (Panning et al., 2022)
 141 and differs by $\sim 30^\circ$ from the BAZ derived from body waves by the MQS (Kawamura
 142 et al., 2022). After rotating into the ZRT coordinates system, the Z and R components
 143 are bandpass-filtered, with a 30% bandwidth, around different central periods T_0 (Fig.
 144 1d). In this way, every T_0 will be associated with one ellipticity value.

145 Subsequently, we obtained the optimal phase shift (ϕ) by maximizing the cross-correlation
 146 between the Z and R components. Examples of resulting ϕ and optimally shifted wave-
 147 forms are shown in Fig. 1d. Once optimally shifted, we computed the cross-correlation
 148 factor $\nu(t)$ between the shifted Z and R components (Fig. 1e). The normalized envelope
 149 $\eta(t)$, also shown in Fig. 1e and given by the multiplication of the Z and R envelopes, is
 150 calculated in order to find the specific time window where the energy is maximized. In
 151 order to discard low-amplitude signals with high cross-correlation (or vice versa), a char-
 152 acteristic function $\chi(t) = \nu(t)\eta(t)$ is calculated (Fig. 1e). The final ellipticity for the
 153 central period T_0 , computed as the R/Z quotient, is the average of the ellipticities in the
 154 time window where the characteristic function χ is larger than 0.8, as shown in Fig. 1f.

155 These processing steps are performed on the ${}_0R_1$ of events S1000a, S1094b and S1222a.
 156 In order to combine these measurements, even though radial and azimuthal anisotropy
 157 may be present, we assume that these would affect ellipticity to a lesser extent than the
 158 underlying crustal structure. To get a single characteristic ellipticity curve at the InSight
 159 landing site, only pairs of event- T_0 with actual Rayleigh wave signal are considered, i.e.,
 160 when the signal-to-noise ratio (SNR) is larger than a given threshold in both Z and R
 161 components. We use $\text{SNR}_Z \geq 10$ and $\text{SNR}_R \geq 5$ (given the higher noise levels on the
 162 horizontals). The final Rayleigh ellipticity curve for the InSight site is then calculated
 163 as the smoothed median of the valid ε values, obtained by means of a Savitzky-Golay
 164 filter (Savitzky & Golay, 1964).

165 2.2 Joint inversion of ellipticity, P-to-s RFs and T_a

166 As the trade-off between the layer velocities and the depth of the discontinuities
 167 is a well-known feature of the ellipticity inversion, the incorporation of independent mea-
 168 surements such as RFs and T_a can help to further constrain the possible models (Chong
 169 et al., 2016). We thus retrieved the 1D crustal structure underneath the InSight land-
 170 ing site by jointly inverting the ε measurements (this work) with the mean P-to-s RF
 171 by Joshi et al. (2023), obtained by stacking radial-component low-frequency RFs for eight

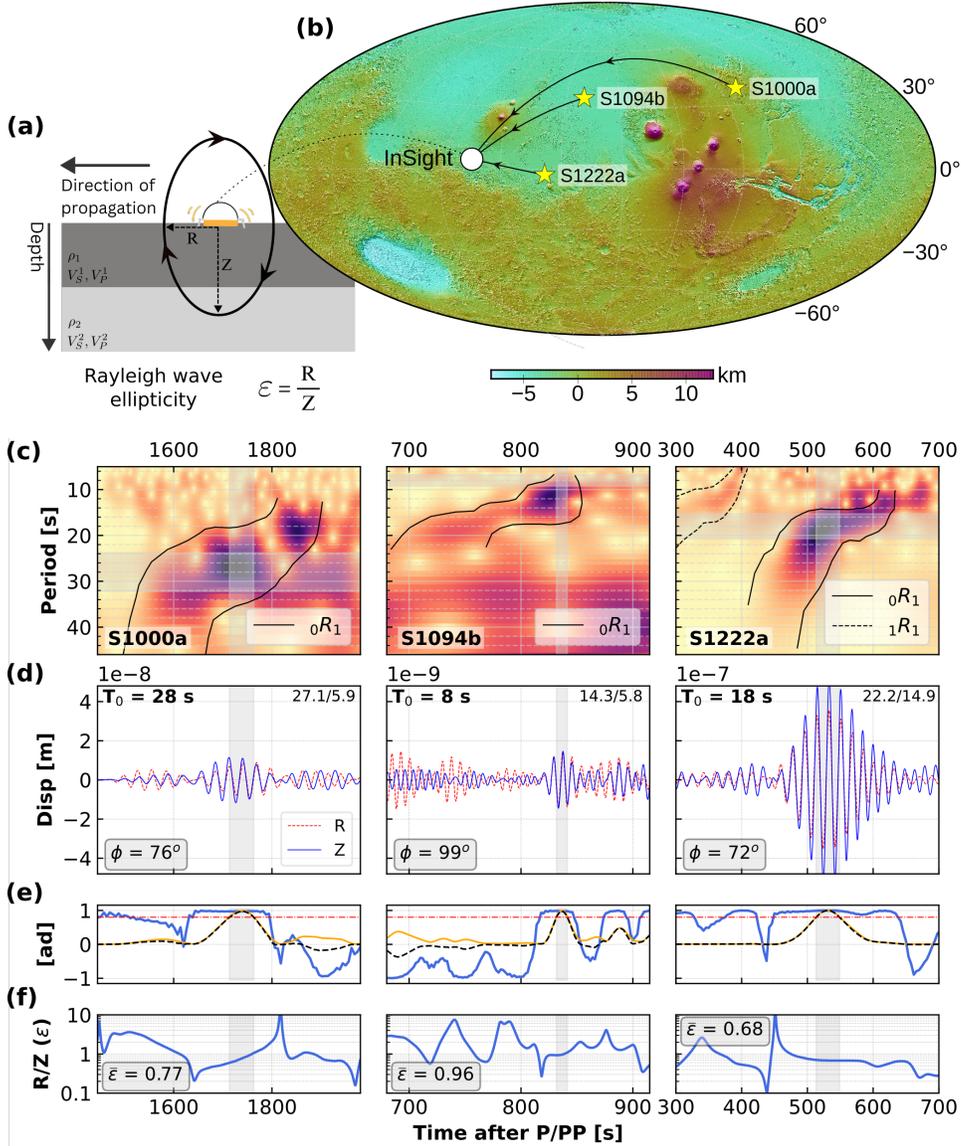


Figure 1. Rayleigh wave ellipticity from large marsquakes. (a) Schematic representation of Rayleigh wave ellipticity using the radial (R) and vertical (Z) ground motion. (b) Topographic map of Mars, from the Mars Orbiter Laser Altimeter (Smith et al., 2001), showing the epicenter of events S1000a, S1094b, S1222a and the InSight lander. Below the map, from left to right, columns correspond to events S1000a, S1094b and S1222a, respectively. From top to bottom: (c) S-transform spectrogram of the Z component, where the fundamental mode and first overtone of the Rayleigh waves are highlighted; (d) shifted Z and R waveforms using the optimal phase shift ϕ . The Z and R components are bandpass-filtered with a 30% bandwidth around the central period T_0 , as indicated. The SNR_Z and SNR_R are given at the top right corner of each plot; (e) cross-correlation factor between shifted Z and R components (blue), normalized envelope $\eta(t)$ (yellow) and characteristic function $\chi(t)$ (dashed black). The red dot-dashed line denotes a threshold at 0.8; (f) ellipticity $\varepsilon(t)$. For each event and each central period, the average ellipticity ($\bar{\varepsilon}$) is computed over the time window where $\chi \geq 0.8$ (gray box).

172 high-quality marsquakes mainly originating in the Cerberus Fossae region (see Joshi et
 173 al. (2023) for further details). Additionally, we included the T_a obtained by Schimmel
 174 et al. (2021), which are understood as the two-way travel times of P waves bouncing back
 175 from the first (~ 6.15 s) and second (~ 10.6 s) strong discontinuity within the crust.

176 For every model, the RF forward calculation is performed via a Python implemen-
 177 tation of the code by Shibutani et al. (1996), using a fixed ray parameter $p^* = 6.9$ s/deg,
 178 suitable for events in the Cerberus Fossae region (Joshi et al., 2023). In the RF, four peaks
 179 at around 0, 2.4, 4.8 and 7.2 s are observed, which are related to the direct P phase and
 180 three converted phases P_{1s} , P_{2s} , P_{3s} , where the subscript indicates the interface where
 181 each phase is generated. The predicted ellipticity curves are obtained via the *gpell* tool-
 182 box included in the Geopsy package (Wathelet, 2005). Lastly, the T_a are computed as-
 183 suming two-way vertical travel paths of the P waves.

184 The inversion scheme utilized is the Neighbourhood Algorithm (NA), first intro-
 185 duced by Sambridge (1999) and widely used for multiple geophysical applications, in-
 186 cluding subsurface characterization (e.g., Hobiger et al., 2013; Ferreira et al., 2020; Car-
 187 rasco et al., 2022). We applied a modified version of the NA, where the parameters space
 188 can be constrained by physical conditions and prior information (Wathelet, 2008). The
 189 parameter space is configured as 3-uniform-layers over a half-space (3LOH), where the
 190 depths of the discontinuities are constrained by Joshi et al. (2023). We inverted for shear-
 191 wave velocity V_s , P-wave velocity V_p and bottom depth z of each layer. Density was linked
 192 to V_p by using the empirical relationship by Brocher (2005), as suggested by Lin et al.
 193 (2012) for handling the trade-off with V_p/V_s , which was allowed to vary between 1.45
 194 and 2.2. The upper mantle is modeled as a half-space with V_s between 4 and 5 km/s,
 195 following the results for the surrounding region using surface waves (Kim, Stähler, et al.,
 196 2022) and teleseismic phases (e.g., Drilleau et al., 2022; Durán et al., 2022). See Table
 197 S4 for further parameterization details.

198 The NA is based on a fully non-linear, self-adaptive Monte Carlo approach, which
 199 efficiently explores the parameters space in order to find the model \mathbf{m} that minimizes
 200 the global misfit function Ψ between the observed and theoretical data. The global mis-
 201 fit depends on the misfit of each data set and their weighting. In this case, we used 0.45,
 202 0.45 and 0.1 as the weights of ε , RFs and T_a , respectively. Further details and other com-
 203 binations of weights can be found in sections S8 and S9 in the Supp Material.

204 3 Results and discussion

205 3.1 Ellipticity and phase-shift observations

206 The resulting ε and ϕ values are shown in Fig. 2a and b, respectively. Besides the
 207 limited number of events analyzed, there is a lack of valid ellipticity measurements due
 208 to the low-amplitude energy on the radial component. The valid ε values define an el-
 209 lipticity curve between 15 and 35 s and two points at 8 and 9 s (Fig. 2a). For the pur-
 210 pose of this paper, we focus on the period range between 15 and 35 s, where the hori-
 211 zontal contribution of the tilt induced by Rayleigh waves can be neglected. Due to the
 212 lack of further observations, a common 15% uncertainty is utilized, which properly cov-
 213 ers our observations.

214 On Earth, the ellipticity curve of one specific site is obtained from the recordings
 215 of a large number of earthquakes. In our work, even though three events were initially
 216 investigated, the ellipticity curve is mainly ruled by S1222a and, therefore, the lack of
 217 observations can be disadvantageous for further analysis and interpretation. Ellipticity
 218 data can have a large spread (e.g. Berbellini et al., 2016; Attanayake et al., 2017), which
 219 might be associated with complex Rayleigh wave propagation effects (Sexton et al., 1977;
 220 Pedersen et al., 2015) or due to local-scale heterogeneities near the receiver, as suggested
 221 by Ferreira & Woodhouse (2007) after analyzing two close stations in California. Con-

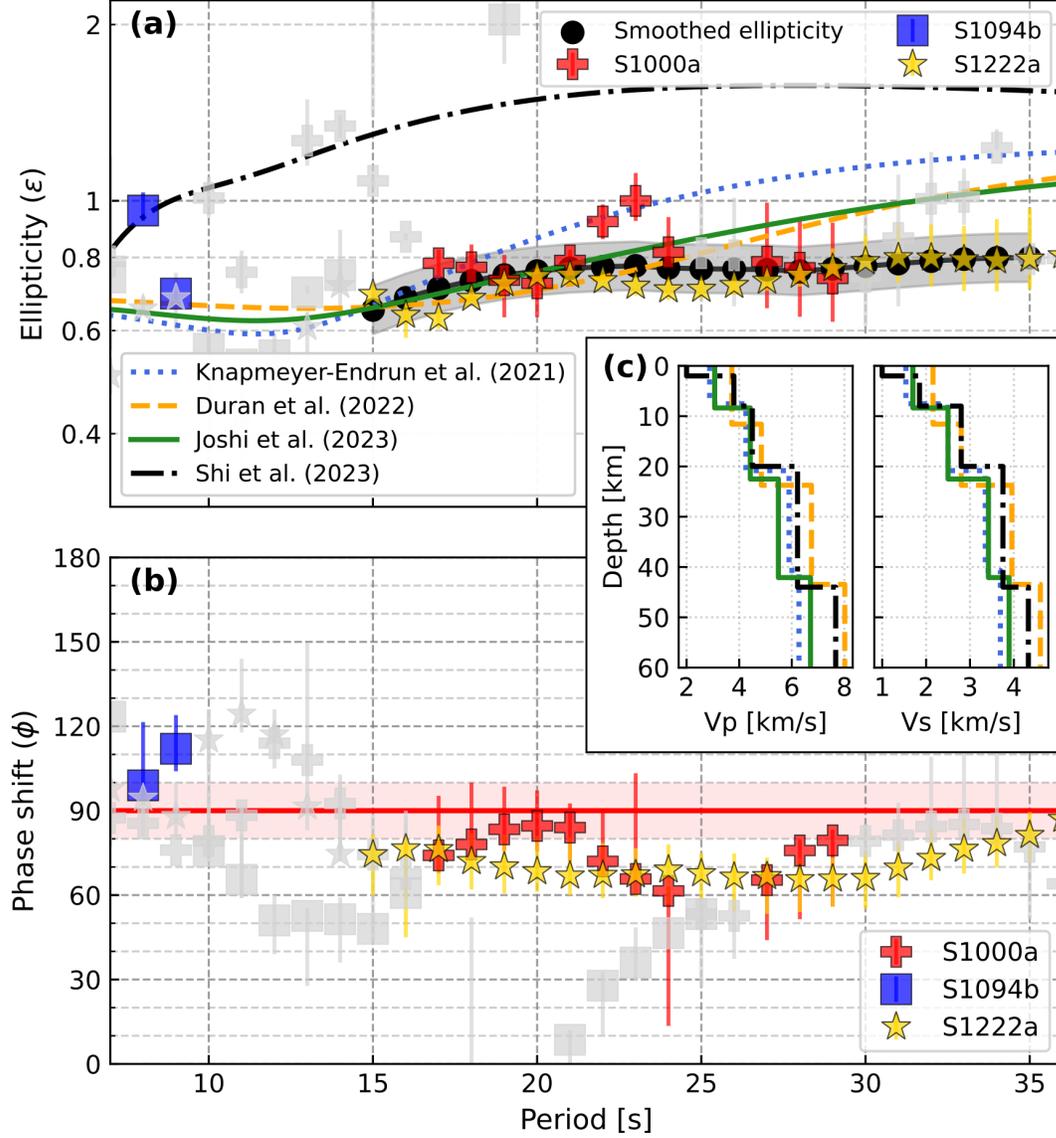


Figure 2. Rayleigh wave ellipticity and phase shift results. **(a)** Rayleigh wave ellipticity (ϵ) curve for events S1000a (red cross), S1094b (blue square) and S1222a (yellow star). Colored markers correspond to periods where the SNR criteria are fulfilled. Background light gray markers did not meet the minimum SNR threshold. The final ellipticity curve for the InSight site is denoted by the black circles, and a 15% uncertainty range has been imposed (gray area). Synthetic ellipticity curves are illustrated for characteristic crustal models from Knapmeyer-Endrun et al. (2021), Durán et al. (2022), Joshi et al. (2023) and Shi et al. (2023). **(b)** Phase shift ϕ between Z and R components of ${}_0R_1$ for the same events and with the same color code as in (a). The light red area corresponds to ϕ in the range $90 \pm 10^\circ$. **(c)** Vp and Vs models for the InSight crust used for synthetic calculation of ϵ in (a).

222 sidering the lower tectonic activity and smooth topography in the local surrounding of
 223 InSight, as compared to the terrestrial case, we assume that the ellipticity measurements
 224 from this event can be a good representation.

225 Besides, the ${}_0R_1$ of S1222a has a high SNR in both Z and R components, so the
 226 final smoothed ε is expected to be a good representation of the local structure. For a bet-
 227 ter assessment, the same ellipticity analysis was performed for one station on Earth (sta-
 228 tion BQ.DREG in Germany, see section S4 in Supp. Material). These measurements sug-
 229 gest that well-recorded events can show some variation with respect to the character-
 230 istic smoothed ellipticity curve of the site, but the broad shape of the curve remains sim-
 231 ilar. In particular, some events show different slopes and trends at specific period ranges,
 232 with respect to what is observed for the median ellipticity (Fig. S3 and S4), but the gen-
 233 eral trend is the same. Thus, we conclude that small oscillations on the ellipticity curve
 234 should not be over-interpreted.

235 The resulting Z-R phase shifts, mainly ruled by S1222a, suggest ϕ values between
 236 60° and 70° (Fig. 2b), which are persistently lower than the theoretical phase shift for
 237 Rayleigh waves under isotropic and homogeneous conditions (90°). On Earth, large de-
 238 viations up to $\pm 30^\circ$ have been found and have been related to small-scale heterogeneities
 239 (e.g., Ferreira & Woodhouse, 2007). From our observations on Earth, low ϕ ($60-75^\circ$)
 240 were measured for some events at specific period ranges (see Fig. S5). These observa-
 241 tions suggest that anomalous ϕ might be related to the Rayleigh-wave path, or even the
 242 source, rather than the characteristics of the receiver site. Tanuma & Man (2008) also
 243 proposed that variations of ϕ with respect to the isotropic theoretical value depend on
 244 the perturbation of the initial stress conditions of the medium. Although crustal anisotropy
 245 was reported along the path of S1222a (Beghein et al., 2022; Kim, Stähler, et al., 2022),
 246 its influence on the phase shift is unclear and should not be discarded. Thus, the low ϕ
 247 measured at the InSight landing site might be due to a combination of path or source
 248 effects with the initial stress conditions at this site, but not directly related to the me-
 249 chanical properties of the crust at the InSight landing site. In any case, low ϕ values have
 250 previously been observed on Earth to occur in some cases. Further analysis of this phe-
 251 nomenon is beyond the scope of this work.

252 **3.2 Local crust structure at the InSight landing site**

253 In order to derive a consolidated crustal structure for the InSight landing site, the
 254 joint inversion detailed in section 2.2 was performed. It is worth noting that the inver-
 255 sion of ellipticity data alone does not properly constrain the layering of the crustal struc-
 256 ture, as different settings between one and four layers over a half-space can equally well
 257 explain the ε observations (Fig. S8).

258 The initial joint inversion results suggest that these local measurements can be well
 259 explained by a three-layer model, with an overall good fitting of the three data sets (see
 260 Fig. S11). However, the amplitude of the P-phase (first peak on the RFs), which is not
 261 normalized as it provides information on the near-surface velocities (e.g., Ammon, 1991),
 262 is clearly overestimated. As shown by synthetic forward modeling (Fig. S12), lower am-
 263 plitudes for this arrival can be achieved by having lower velocities in the first layer. Be-
 264 cause of this, we performed a second inversion where we allowed a shallow low-velocity
 265 layer (4LOH, see Table S5 for details on the parameter space). These new resulting mod-
 266 els are preferred as they explain the three independent data sets reasonably well (Fig.
 267 3a-d), and especially the amplitude of the direct P-phase arrival is well retrieved.

268 Our inversion results show that a 4-layer crust structure (Fig. 3e-g), including a
 269 shallow low-velocity layer (sLVL) close to the surface, can explain the local-scale mea-
 270 surements well. The sLVL presents velocities between 1.6 and 1.8 km/s and extends down
 271 to 2-3 km depth. Its origin is unclear, but previous works on the shallowest structure
 272 of Mars have proposed the presence of sedimentary and heavily fractured basalt layers

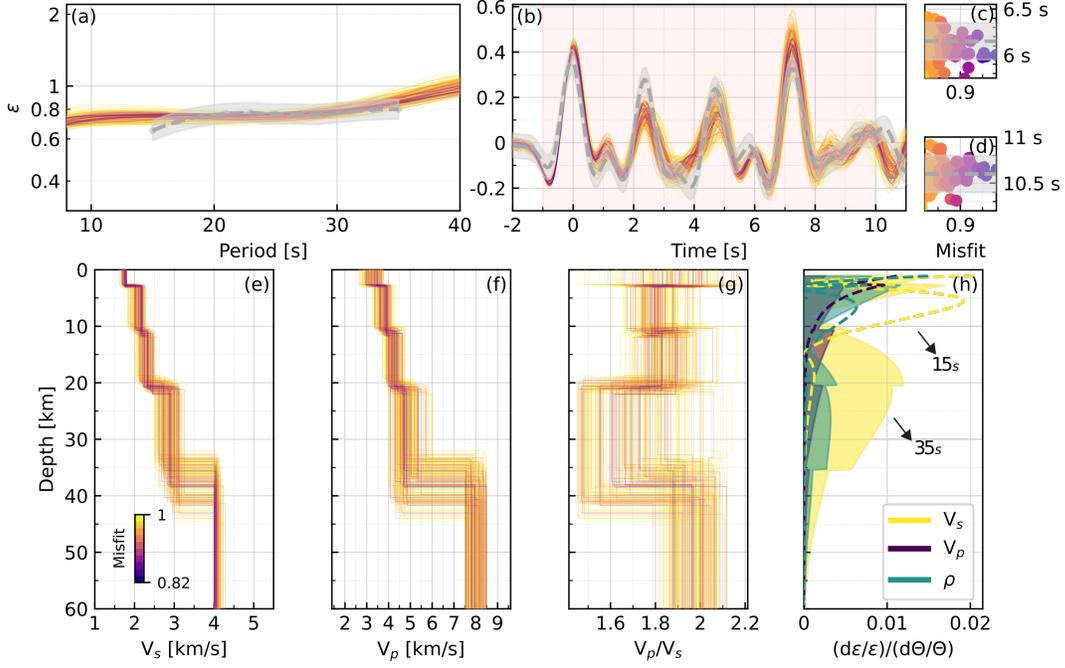


Figure 3. Joint-inversion results and resulting Martian crustal models using a 4LOH parameter space. The measurements and the synthetic modeling of 200 best-fitting crustal models are shown for (a) Rayleigh wave ellipticity, (b) P-to-s RFs from Joshi et al. (2023), and (c) first and (d) second P-wave reflection times from Schimmel et al. (2021). The light gray dashed line and area correspond to the measured data and their uncertainty, respectively, whereas the synthetic data are colored by their misfit. The corresponding crustal models are shown in terms of (e) V_s , (f) V_p and (g) V_p/V_s , as a function of depth. (h) Depth absolute sensitivity kernels of ellipticity data for the best-fitting model, calculated at 15 s (dashed line) and 35 s (filled area) for V_s , V_p and ρ . Light red box in (b) shows the time window where the data was fitted.

273 in the first hundreds of meters (Hobiger et al., 2021; Carrasco et al., 2022), where $V_s <$
 274 1 km/s were derived. Furthermore, Pan et al. (2020) suggested the presence of a sed-
 275 imentary layer extending from 200-300 m up to kilometers in depth, whereas Warner et
 276 al. (2022) showed evidence for heavily fractured basalt units in the surrounding area. Thus,
 277 we interpret this shallow discontinuity to be the bottom of this large-scale unit.

278 The sLVL results from combining the different data sets through the joint inver-
 279 sion, rather than from a specific feature of the RF. Indeed, Shi et al. (2023) showed that
 280 the ~ 1 s signal in the high-frequency RFs, related to the sLVL, is absent in the low-frequency
 281 RFs. Despite observing a similar 1 s signal in the low-frequency RF used in this work,
 282 the RF-only inversion with 3LOH shows that it is not necessarily related to an sLVL (Fig.
 283 S9). Instead, given the mantle velocity, the ellipticity data allows a range of average crustal
 284 V_s to fit the data, which leads to too high velocities near the surface for the RF, so a
 285 top layer with lower velocities is required.

286 In general, below the sLVL, our models have a similar structure and the discon-
 287 tinuities at ~ 10 and ~ 20 km depth are in good agreement with previously reported mod-
 288 els for the local crust (e.g., Knapmeyer-Endrun et al., 2021; Joshi et al., 2023). The dis-
 289 continuity around 37 km depth is interpreted as the crust-mantle boundary, where the
 290 sharpest velocity contrast is observed. Even though the ellipticity data has low sensi-
 291 tivity at this depth, it can properly constrain the shallow V_s structure (Fig. 3h), which
 292 directly affects the deeper structure as they are strongly tied through the RFs. Thus,
 293 as the shallow part is well constrained, the joint inversion can provide reliable depths
 294 for the Moho.

295 Overall, crustal V_p and V_s are lower than 5 and 3 km/s, respectively (Fig. 3e,f).
 296 V_p/V_s is around 1.8 in the first 20 km of the crust and, even though large variation is
 297 observed, it seems to decrease in the bottom layer of the crust (Fig. 3g). In particular,
 298 V_s are lower than in other regions on Mars, such as along the path between Amazonis
 299 Planitia and the InSight site in the lowlands (Kim, Banerdt, et al., 2022), or between
 300 S1222a and the lander, along the dichotomy region (J. Li, Beghein, Lognonné, et al., 2022).
 301 Although an increase in density is derived from the inversion, the inverted data sets are
 302 poorly sensitive to this parameter, as shown for the ellipticity (Fig. 3h), so further in-
 303 terpretation is avoided. At a local scale, our proposed V_s are also slightly lower than those
 304 previously reported for the InSight site (Knapmeyer-Endrun et al., 2021; Joshi et al., 2023).
 305 Therefore, considering that porosity plays an important role in decreasing the velocity
 306 of the seismic waves (e.g., Heap, 2019), these lower velocities might be related to an even
 307 larger porosity or stronger alteration of the whole crust at the InSight site than previ-
 308 ously thought.

309 4 Conclusions

310 For the first time, the ellipticity of the fundamental mode of direct Rayleigh waves
 311 was measured on Mars for periods between 15 and 35 s, especially thanks to S1222a. The
 312 obtained ellipticity curve is interpreted as a good representation of the characteristic el-
 313 lipticity at the InSight landing site. Anomalously low phase shifts were measured, which
 314 seem to be unrelated to the local crust structure. In order to consolidate the local crustal
 315 model, we performed a joint inversion of local-scale observations (ellipticity, P-to-s RFs
 316 and P-wave reflection times). Our results suggest that a four-layer crustal model, with
 317 a shallow low-velocity layer of about 2-3 km thickness at the InSight landing site explains
 318 the observations well. Other discontinuities around 10, 20 and 37 km depth correlate well
 319 with previously reported models. We propose lower P- and S-wave velocities for the cor-
 320 responding crustal layers, which might be related to a higher porosity or alteration of
 321 the Martian crust at the InSight landing site.

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Open Research

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The event information of the Martian seismic events can be found in the InSight seismic event catalog version 13 (InSight Marsquake Service, 2023). The waveform data and station metadata are available from IPGP's MSDC as well as from IRIS MDC (InSight Mars SEIS Data Service, 2019). Seismic data for station DREG (BQ network, 2016) are publicly available via EIDA (<http://eida.gfz-potsdam.de/webdc3/>). The GCMT catalog is available here <https://www.globalcmt.org/CMTsearch.html>, whereas the events analyzed for the Earth case and the codes utilized for the processing of the events can be found under <https://github.com/scarrascom/Rellipy>. Observed data used for the joint inversion, a Python implementation and preferred crustal models shown in Fig. 3 can be found at [10.5281/zenodo.8051337](https://zenodo.org/record/8051337) (Carrasco, 2023). The NA was implemented via the *dinverext* plugin, which can be obtained from <http://www.geopsy.org/>.

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