

Fluid-induced anthropogenic and natural earthquake swarms driven by aseismic slip

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Key Points:

- Scaling laws show that injection-induced and natural earthquake swarms have the same driving mechanism.
- We show that aseismic slip is always present in swarms, although its contribution differs from one swarm to another.
- We propose a model based on fluid-induced aseismic slip propagation that explains swarms behavior.

Abstract

Anthropogenic fluid injections at depth induce seismicity which is generally organized as swarms, clustered in time and space, with moderate magnitudes. While some similarities between swarms have already been observed, whether they are driven by the same mechanism is still an open question. Pore fluid pressure or aseismic processes are often proposed to explain observations, while recent models suggest that seismicity is triggered by fluid-induced aseismic slip. Using 22 natural and anthropogenic swarms, we observe that duration, migration velocity and total moment scale similarly for all swarms. This confirms a common driving process for natural and induced swarms and highlights the ubiquity of aseismic slip. We propose a method to estimate the seismic-to-total moment ratio, which is then compared to a theoretical estimation that depends on the migration velocity, the effective stress drop and the slip velocity. Our findings lead to a generic explanation of swarms driving process.

Plain Language Summary

Swarms are a particular type of seismic sequence, during which many earthquakes occur but with no mainshock distinguishable from the other events. They can be induced by anthropic injections at depth, like during geothermal exploitation. Natural swarms are also observed in a large variety of geological contexts. Natural and injection-induced swarms share a lot of similarities, like the migration of seismicity. But little is still known about their physics. Here, we explain the observed similarities by the fact that both types of swarms correspond to earthquakes triggered by the propagation of an aseismic slip transient, induced by fluid circulation. This allows to reconcile observations made over different length- and timescales, and provides a generic explanation of the processes occurring at depth.

1. Introduction

1.1 Natural and injection-induced swarms exhibit many similarities

Fluid pressure changes at depth can induce seismicity, as shown by the increase in seismicity near fluid injection sites during geothermal activities (e.g., in Basel, Switzerland; Diechmann and Giardini, 2009), wastewater storage in Oklahoma (Hincks et al., 2018), or during fault activation experiments in France (Guglielmi et al., 2015). On the other hand, earthquake swarms of natural origin (i.e., sequences of clustered earthquakes with moderate magnitudes, generally

below $M_w = 5$, without a mainshock-aftershock pattern) occur in various geological and tectonic contexts, such as mountain ranges (Jenatton et al., 2007), rift zones (De Barros et al., 2020) or along transform faults (Roland and McGuire, 2009). Earthquakes during those swarms are located over one or several fault planes (Lohman et McGuire, 2007; Baisch et al., 2009; Hong et al., 2020; Fischer et Hainzl, 2021). Migration of seismicity is the most characteristic behavior of both injection-induced and natural earthquake swarms (Goebel et Brodsky, 2018; Passarelli et al., 2018). Proposed physical explanations for swarm migration include fluid pressure diffusion (Shapiro et al., 1997), aseismic slip (Roland et McGuire, 2009), or a combination of both (De Barros et al., 2021), as well as cascading events (Fischer et Hainzl, 2021).

To gain deeper understanding into swarm processes and evaluate how generic they are, we compare injection-induced swarms with natural ones. Given the similarities identified between the two types of swarms, we aim at evaluating if a common mechanical process may drive swarms in different geological contexts and origins.

1.2 Understanding the role of aseismic slip in swarms

Aseismic moment release is thought to occur for injection-induced sequences, as revealed by moment-volume scaling relations (McGarr et Barbour, 2018; De Barros et al., 2019), by geodesy in the vicinity of a fluid injection site in the Brawley Basin (Wei et al., 2015), by measurements during field experiments (Guglielmi et al., 2015) or indirectly by studying repeating earthquakes during the Soultz-Sous-Forêt sequences (Bourouis and Bernard, 2007; Lengliné et al., 2014). The relatively weak values of seismic moment released compared to the spatial extent of seismicity also indicate that aseismic slip occurs over the whole seismicity area (Fischer and Hainzl, 2017). Aseismic slip has also been observed in association with natural swarms, using geodesy and slip

inversions (Lohman and McGuire, 2007; Gualandi et al., 2017), or by studying dual velocity migrations and repeating earthquakes during the 2015 swarm in the Gulf of Corinth (De Barros et al., 2020). Nevertheless, geodetic observations of aseismic slip associated with swarms remain rare and difficult to achieve, given the depth, long duration and low deformation of such sequences.

Numerical modeling showed that the increase of the critical earthquake nucleation size with fluid pressure first leads to aseismic slip, which may outpace the diffusing pressure front (Bhattacharya and Viesca, 2019; Laroche et al., 2021) and which triggers seismicity near its edges where shear stresses increase (Cappa et al., 2019; Wynants-Morel et al., 2020).

As illustrated on Figure 1, and based on the previously mentioned observations, while aseismic slip is directly induced by fluid pressure, earthquake swarms seem to be triggered by the shear stress perturbation resulting from aseismic slip propagation over brittle asperities, rather than by fluid overpressure. In this case, seismicity migration would be related to the aseismic slip propagation, and not to the diffusion of fluids (Bhattacharya et al., 2019; De Barros et al., 2021).

This is analogous to the observed co-location of seismic and aseismic slip areas during large slow slip events (SSEs) in subduction zones, as in Cascadia (Bartlow et al., 2011). The seismic events (tremors or earthquakes) are triggered by the stress transfer from the SSEs, even though such SSEs are not necessarily driven by pore fluid pressure perturbation.

In this work, we aim at exploring the similarities between injection-induced sequences and natural swarms in a general way, in order to infer if both types of seismicity can be explained by a common process, and in which extent aseismic slip is driving them. Using a dataset of 22 seismic sequences, we first investigate scaling relations between moment, migration and duration and we compare them to slow-slip events observations. As aseismic slip seems to be a common

feature among swarms, we then introduce a method to estimate the total and aseismic deformation. Finally, we propose a mechanical framework that relates the seismic to total ratio to seismic observables.

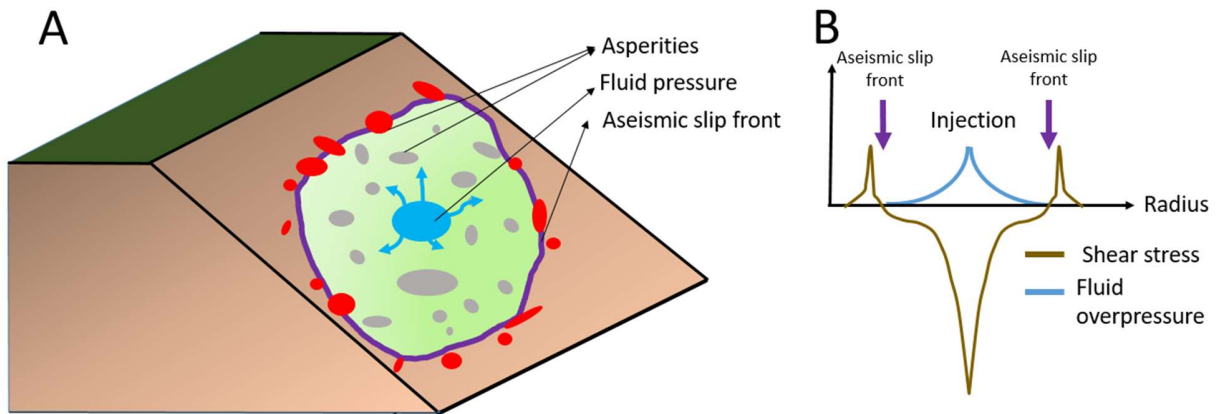


Figure 1. Common model for natural and induced swarms. (A) Conceptual view. Over a fault plane, fluid overpressure (blue arrows), either from anthropogenic or natural origin, induces aseismic slip (light green area). As the aseismic slip zone expands, it triggers seismicity near its edges (red patches) or within it (grey patches), through shear stress perturbation (brown curve, right). (B) Shear stress and fluid overpressure versus radial distance to the injection. The fluid overpressure induces an aseismic slip, with a shear stress drop within the slipping area and a stress concentration at its tip.

2. Materials and Methods

2.1 Data

We focus on a global dataset of 22 earthquake swarms, from either injection-induced or natural origin. For natural sequences, geological contexts are diverse: for instance, the 2003-2004 Ubaye swarm (Jenatton et al., 2007) occurs in a near-zero strain-rate area in the French Alps, while the 2015 Corinth swarm (De Barros et al., 2020) takes place in a very fast extensional (~ 15 mm/year) rift zone in Greece. We here focus on natural swarms in which fluid processes have been previously discussed, and we do not consider swarms taking place near volcanoes or in subduction zones. Most of the injection-induced swarms we consider originated from geothermal exploitation. However, they span a wide range of characteristics, including the injected fluid volume and the injection depth (see Supp.).

2.2 Migration velocity

Migration velocity of the 22 swarms is computed by fitting the seismicity front. Migration period is defined as the time during which the swarm is expanding. The spatial origin is chosen as the median of the coordinates of the 10 first events, and the origin time is defined as the time of the first event. We define the seismicity front as the 90th percentile of event distances in a sliding window containing 50 events. Seismicity front has been modelled by either diffusive law, linear fit or more complex relationships (Goebel et Brodsky, 2018; De Barros et al., 2021). The shape of the migration is here not investigated, as we only focus on an average migration velocity, in order to make comparisons among the different swarms. We fit a linear model over the seismicity front during the migration period of each sequence. This procedure yields migration durations and average migration velocities for each sequence (Figure 2A,B).

2.3 Seismicity area and effective stress drop computation

By analogy with the moment-size relationship for circular ruptures, the effective stress drop of a swarm is defined as (Fischer et Hainzl, 2017):

$$\Delta\sigma_e = \frac{7M_{0,seismic}}{16R^3} \quad (\text{Equation 1})$$

where R is the radius of the seismicity area and $M_{0,seismic}$ the seismic moment released during the swarm. The effective stress drop value is an indicator of the relative importance of aseismic moment release (Fischer and Hainzl, 2017): a low effective stress drop suggests distant seismic asperities embedded in a fault slipping aseismically, while values close to earthquake static stress drops suggest that seismic asperities cover most of the slipping area.

Following Fischer and Hainzl (2017), seismicity area is computed by fitting a 2D plane over the 3D distribution of hypocenters, after removing the outsiders biasing the plane fitting.

Hypocenters are then projected over the plane, and a ConvexHull algorithm delineates and returns the seismicity area S . We then compute a characteristic size, defined as $R = \sqrt{S/\pi}$, and with the cumulative seismic moment value, we compute the effective stress drop.

2.4 Total moment estimation

Aseismic slip quantification is difficult for injection-induced sequences as deformations associated with those episodes are small (0.5mm for the Guglielmi et al., 2015 field experiment), over long durations, leading to small strain rates hard to observe. This issue stays the same for natural swarms.

We here propose a simple way to estimate, roughly, the aseismic slip. Studies of slow slip transients have shown that the slip released by repeating earthquake sequences equals the surrounding aseismic slip (Matsuzawa et al., 2004; Uchida, 2019). As the front of the swarm seismicity is assumed to be directly triggered by the shear stress perturbation induced by aseismic

slip (Figure 1), we here make an analogy with slow slip transients. We suppose that the slip released seismically over discrete asperities equals the surrounding aseismic slip and neglect the afterslip (see Supp.). Assuming the asperity associated with the largest earthquake in the swarm only ruptures once, its slip gives an order of magnitude of the slip over the whole area. For each sequence, we isolate the largest event of moment $M_{0,max}$ and assuming a circular rupture with a static stress drop $\Delta\sigma_{max}$ of 10MPa (unless a more precise value is provided in the literature, see Supp.), we compute the slip D_{max} over this asperity (Madariaga, 1976) as:

$$D_{max} = M_{0,max}^{1/3} * \frac{(16\Delta\sigma_{max})^{2/3}}{G\pi^{2/3}} \text{ (Equation 2)}$$

Given that seismic moment is released over brittle asperities and aseismic slip is released in between them, we estimate the total moment (seismic + aseismic) over the seismicity area as

$$M_{0,total} = G * D_{max} * S \text{ (Equation 3)}$$

where G is the rock shear modulus (taken here as 30 GPa) and S the previously computed seismicity area. While the effective stress drop qualitatively indicates the importance of aseismic slip during a swarm, the rough quantification approach here allows us to better constrain aseismic moment release for each sequence.

2.5 Seismic to total moment ratio

By considering seismic and aseismic slip into one single slip event over a circular area of radius R, we have (Madariaga, 1976):

$$M_{0,total} = \frac{16}{7} \Delta\sigma_{total} R^3 \text{ (Equation 4)}$$

Where $\Delta\sigma_{total}$ is the total stress drop over the studied area.

The rupture velocity of slow slip events is related to its stress drop and to its maximum slip velocity V_{max} (Ampuero and Rubin, 2008; Rubin, 2008; Passelègue et al., 2020) with:

$$V_{rupt} = \frac{G * V_{max}}{n * \Delta\sigma_{total}} \text{ (Equation 5)}$$

To establish this, we assume that the stress drop of the slip event, $\Delta\sigma_{total}$, is proportional (factor $n > 1$) to the associated strength drop. This can be observed in several numerical simulations of slow slip, where $n \sim 10$ (Hawthorne and Rubin, 2013; Lambert et al., 2021).

In our case, we make the hypothesis that the seismicity is triggered by the fluid-induced aseismic slip. Therefore, the seismicity front follows the aseismic front (Wynants-Morel et al., 2020). The migration velocity of the swarms is then the rupture velocity of the aseismic slip ($V_{rupt} = V_{migr}$).

Combining Equation 4 and Equation 5 we then have:

$$M_{0,total} = \frac{16}{7} * \frac{G * V_{max}}{n * V_{migr}} * R^3 \text{ (Equation 6)}$$

This leads us to the following expression for the ratio r of seismic to total moment:

$$r = \frac{M_{0,seismic}}{M_{0,total}} = \frac{7 * M_{0,seismic}}{16 * R^3} \frac{n * V_{migr}}{G * V_{max}} \text{ (Equation 7)}$$

This equation can be written in a more compact form using the effective stress drop (see Equation 1). We then get:

$$r = \frac{n * \Delta\sigma_e * V_{migr}}{G * V_{max}} \text{ (Equation 8)}$$

This relation links the ratio of the cumulative seismic to total moment to the product of the migration velocity and the effective stress drop of the swarm.

3. Results

3.1 Aseismic slip drives the swarm's dynamics

The estimated velocities for the 22 swarms studied here range between a few meters per day, like for the Cahuilla swarm (Ross et al., 2020), to more than 1 km/day in the case of the Rittershoffen sequence (Lengliné et al., 2017). Figure 2C shows the migration velocity as a function of duration, for induced and natural swarms. For sake of comparison, we add the migration velocity of SSEs recorded on subduction zones (Gao et al., 2012). For these events, velocities correspond to the propagation of an aseismic slip, which is characterized either with geodesy (Schmidt and Gao, 2010) or with tremor migration (Bartlow et al., 2011; Ito et al., 2007).

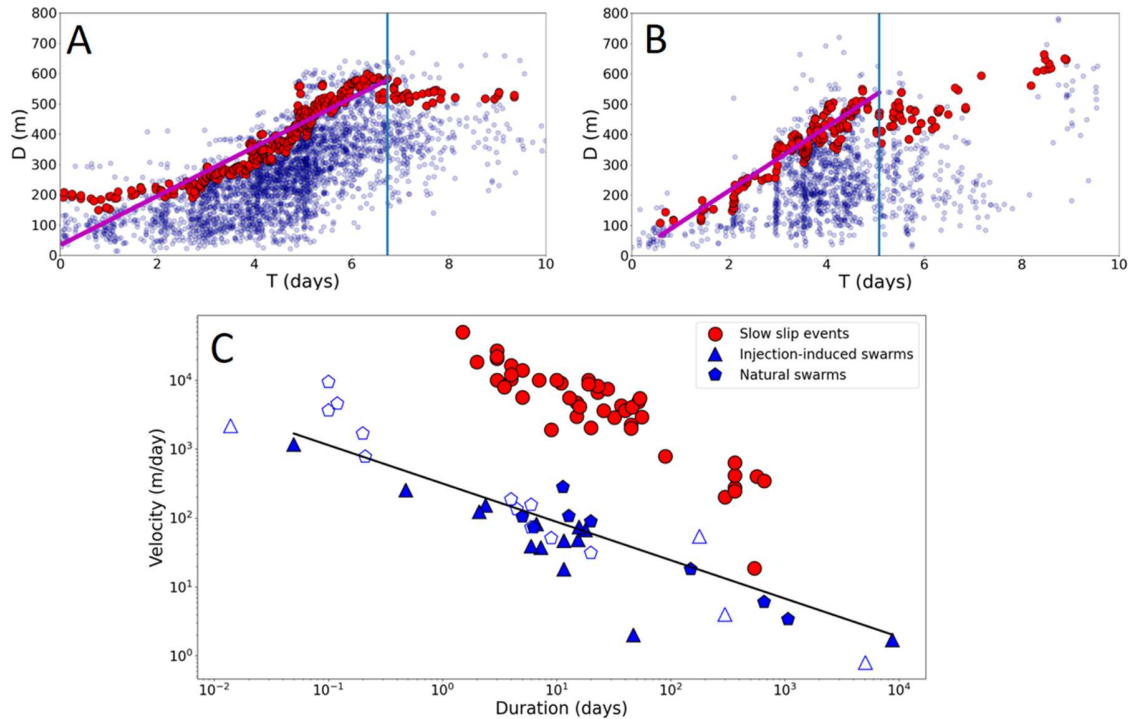


Figure 2. Scaling of propagation velocity with duration for swarms and slow slip events (SSEs). (A) Space-time distribution of seismicity in the Basel sequence (Herrmann et al., 2019). Blue dots represent individual event hypocenters while red circles represent the computed seismicity front. Linear fitting (magenta line) over the seismicity migration period, delimited by the vertical blue line, yields a propagation velocity of 81 m/day. (B) Same but for the Corinth 2015 swarm (De Barros et al., 2020). The seismicity front migrates at a velocity of 105 m/day. (C) Scaling of

velocity with duration. Red dots represent SSE data from (Gao et al., 2012). Filled triangles and pentagons represent injection-induced and natural swarms, respectively, for which we determined migration velocity and duration based on seismicity catalogs. Empty symbols represent migration velocities and durations directly taken from the literature (see Supp.). Black line represents the best-fitting line between our computed velocities and durations ($R^2 = 0.76$).

Two main observations can be made. First, injection-induced and natural swarms follow the same scaling $V \propto T^{-\gamma}$ of velocity with duration, with $\gamma = 0.6$ and $\gamma = 0.7$ when considering each subset individually. In addition to the other similarities discussed beforehand, the continuous scaling of velocity with duration for all swarms is direct evidence that both types of sequences obey the same physics for all velocity ranges (from a few meters per day in Ubaye or Cahuilla to 1160 m/day for Rittershoffen). As anthropogenic seismicity is induced (though indirectly) by fluid injection (Bentz et al., 2020), this confirms that natural swarms are also a consequence of fluid pressure perturbations.

Second, the velocity-duration scaling is the same for swarms and for the SSEs reported by Gao et al. (2012), despite higher velocities for the latter, typically around 1 to 10 km/day. This confirms that the migration of swarms globally behaves as the propagation of aseismic slip, hence the assumption made of $V_{\text{rupt}} = V_{\text{migr}}$. The observed scaling for swarms, $V \propto T^{-\gamma}$ with $\gamma = 0.55$, is compatible with fluid diffusion. However, a similar scaling is obtained for SSEs, which exhibit individual linear migrations (Houston et al., 2011) and are not driven by fluid diffusion. Other mechanisms have been proposed to explain such a scaling for SSEs, like a uniform stress drop or a uniform slip (Ide et al., 2007). These mechanisms might also be valid for swarms, explaining then the observed continuum of characteristics (Figure 2C). Therefore, a general scaling compatible with diffusion does not imply that individual swarm are directly driven by fluid diffusion.

The effective stress drop $\Delta\sigma_e$ is found to range between 1 kPa and 1 MPa (Figure 3). Those values are lower than typical values of static stress drop for earthquakes, which usually range between 1 and 100 MPa (Cocco et al., 2016) and are more similar to the stress drop values of SSEs (Brodsky and Mori, 2007). Thus, they indicate an aseismic component in the swarm processes. For instance, $\Delta\sigma_e = 1\text{kPa}$ for the Soultz-sous-Forêt stimulation, indicates an important aseismic moment release, while $\Delta\sigma_e = 1\text{MPa}$ for the Basel injection means that aseismic slip is relatively less important in this case. $\Delta\sigma_e$ ranges in a similar way for natural and injection-induced sequences (Figure 3), indicating once again that mechanisms of seismic and aseismic moment release are controlled by the same processes for both types of sequences.

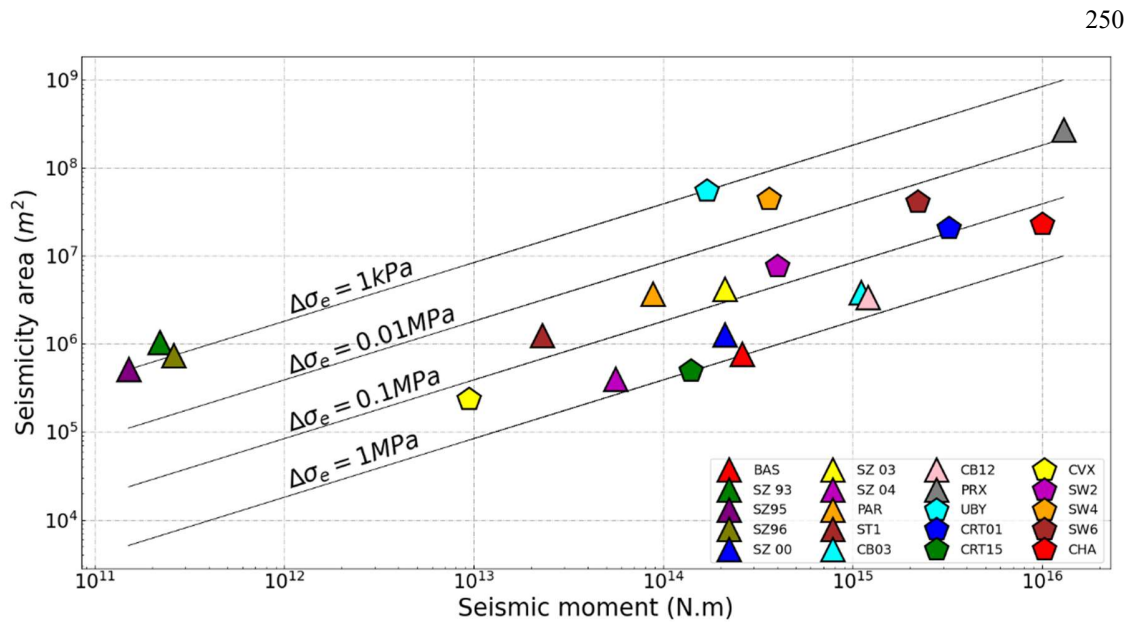


Figure 3. Seismicity area (m^2) as a function of the cumulative seismic moment released during 20 of the swarms studied here. Triangles correspond to injection-induced sequences while pentagons refer to natural swarms. Black lines represent different values of the effective stress drop $\Delta\sigma_e$.

Based on migration velocity and effective stress drop analysis, both natural and injection-induced swarms seem to share the same driving processes, in which aseismic slip seems ubiquitous, like depicted on Figure 1. The seismicity front is depicting the aseismic slip rupture propagation and the seismicity area corresponds to the aseismic slip area, in a similar way as tremors locations in SSEs zones delineate slip migration and area (Bartlow et al., 2011). However, the aseismic contribution might be different from one swarm to another.

3.2 Aseismic contribution differs among swarms

Once the total moment $M_{0,total}$ for each swarm is computed, we compare it to the seismic moment released by using the seismic to total moment ratio r (Equation 7).

A value of r close to 1 indicates that moment release is mainly seismic, while a lower value shows that moment release is significantly aseismic. We compute $M_{0,total}$ and r for the swarms studied here. As shown in Figure 4A, r ranges from 0.001 to almost 1. For the Basel injection-induced sequence, $r = 0.97$, suggesting that aseismic deformation is low in this case, while for the Ubaye natural swarm, $r = 0.005$, indicating an important aseismic moment release.

For the Soultz 1993 sequence, despite an injected fluid volume of the same order of magnitude as in the Basel injection (Diechmann et Giardini, 2009), the cumulative seismic moment is 3 orders of magnitude lower than the Basel one. This can be explained here by an important aseismic moment release ($r \sim 0.001$) taking place during the Soultz sequence. Therefore, the strong difference of seismic moment release for similar volumes can simply reflect the amount of induced aseismic deformation (McGarr and Barbour, 2018; De Barros et al., 2019).

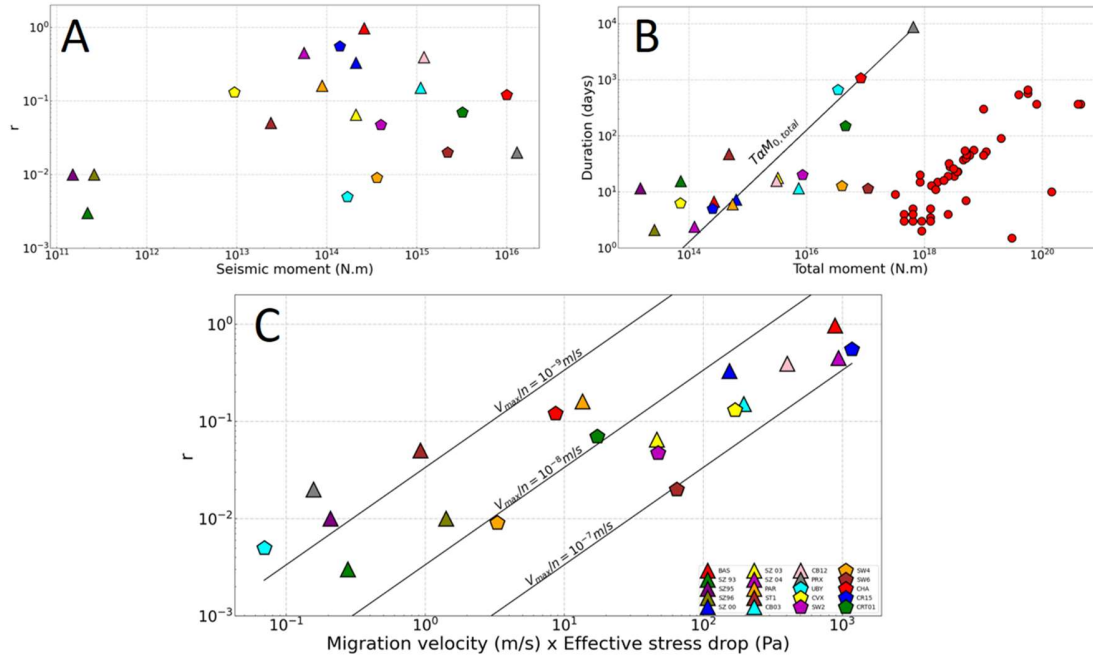


Figure 4. (A) Seismic to total moment ratio, as a function of the seismic moment released during each swarm, for the sequences studied here. (B) Duration as a function of the estimated total moment. Black line represents the 1:1 scaling. Red dots correspond to the SSE data from Gao et al., 2012. (C) Seismic to total moment ratio for the swarms studied here, as a function of the product of the migration velocity and the effective stress drop. The black lines correspond to different values of $\frac{V_{max}}{n}$, assuming $G = 30\text{GPa}$ (see Equation 8).

Interestingly, one can also note that the scaling of duration with estimated total moment (Figure 4B) seems to be 1:1, similarly to the scaling between event duration and aseismic moment observed for SSEs (Ide et al., 2007; Peng and Gomberg, 2010), while seismic moment vs duration does not exhibit such a scaling (Passarelli et al., 2018). Indeed, with our total moment estimate, we are able to measure the “hidden” aseismic slip release. As hypothesized (Peng and Gomberg, 2010), the apparent branching off of the swarms in the moment duration can be corrected when considering the aseismic deformation.

Using Equation 8, we can relate the seismic to total moment ratio to two observables, effective stress drop and migration velocity (see Figure 4C). Following Equation 8, we estimate V_{max}/n being between 10^{-10} and 10^{-7} m/s, which is consistent with expected orders of magnitudes (Roland and McGuire, 2009; Glowacka et al., 2001) if we consider a value of $n \sim 10$ (Hawthorne and Rubin, 2013; Lambert et al., 2021). Variability in V_{max} explains why the observed scaling between r and $\Delta\sigma_e * V_{migr}$ is not as linear as expected. As the general trend shows a scaling different than the isovalues of V_{max}/n , it means that V_{max} also depends, through fault and stress properties, on the seismic-to-total seismic ratio. The slip velocity, together with the migration velocity and the effective stress drop, are the crucial parameters to characterize the seismic and aseismic moment partitioning in swarms. Among other properties, these three parameters depend on the stress state and on the proximity of the fault to failure (Hainzl and Fischer, 2002; Fischer and Hainzl, 2017; Passelègue et al., 2020; Wynants-Morel et al., 2020; De Barros et al., 2021). These relationships therefore deserve to be investigated in order to anticipate the swarm evolution, especially given that similarities are found between swarms and foreshock sequences of some major earthquakes (Chen and Shearer, 2013).

4. Conclusions

In this work, we confirmed that injection-induced and natural swarms are governed by the same physics, as was previously shown for particular sequences (Fischer and Hainzl, 2017). By analyzing sequences covering a wide range of geological contexts, migration velocities, durations and injected fluid volumes, we showed a global unity in the swarm's dynamics. After confirming that fluid-induced aseismic slip explains observations made on swarms, like their migration or their spatial seismic moment release, we exploited the similarities between swarms and slow slip events to introduce a simple mechanical framework that relates the seismic and aseismic moment

partitioning to physical and observable parameters (Equation 8). This opens interesting perspectives to better understand swarms, their propagation, and improve their monitoring in order to anticipate potential large earthquakes. It also paves a way to studying natural and injection-induced swarms as the same phenomena.

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