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2 **First Detection of Extremely Enhanced Solar Wind Helium-3 Originated from the**
3 **Lunar**

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9 **Key Points:**

- 10 • The first observations of ${}^3\text{He}^{2+}$ density extremely enhanced to thousand times of average
11 solar wind by SWID/Chang'E-2 is reported.
- 12 • Neutrals ionization and surface charging caused by energetic electrons induce ${}^3\text{He}$ losses
13 from the lunar exosphere and may form ion voids.
- 14 • The loss of the exosphere and the formation of ionosphere ion void will be recurrent for
15 the frequent energetic electrons injection.
16

17 Abstract

18 The first observations of ${}^3\text{He}^{2+}$ density extremely enhanced to thousand times of average solar
19 wind by SWID/Chang'E-2 is reported in this paper. The two events occurred near the lunar and
20 the Lagrange L2 point respectively. The synchronized energetic electrons increasing can induce
21 the lunar surface charging to a large negative voltage and produce the electron-impact double
22 ionization of helium-3 in the exosphere. ${}^3\text{He}^{2+}$ will converge to the lunar surface by the negative
23 potentials, and fully pick up via non-resonance stochastic heating process by Alfvén wave of the
24 solar wind. It indicates ${}^3\text{He}^{2+}$ in the events can originate from the lunar and transport to L2.
25 Besides SEP (solar energetic events), the increasing of energetic electrons occurs frequently in
26 quiet time, so the loss of the lunar exosphere will be recurrent, and the ionosphere ion void may
27 be formed.

28 Plain Language Summary

29 Chang'E-2 is a lunar orbiter of launched in 2010. Solar wind ions are measured by SWIDs on
30 board Chang'e-2. Helium-3 is one stable isotope of helium, which is expensive in earth because
31 of its sparsity. The Solar wind Helium-3 from the sun inject and store in the lunar regolith. So
32 there is abundant helium-3 in the lunar environment. It's the first time we found the abnormal
33 increase to thousands times of helium-3 ions in the solar wind when the Chang'E-2 is near the
34 lunar and at the downstream of the lunar. Due to the scarcity of helium-3 on earth, it can be
35 considered that the sun or the moon is its source. In the past, the observations of helium-3
36 enhancement in the solar wind originated from the sun have been reported. The enhancement can
37 only reach 100 times. Our observations and researches show that the abnormal increasing of
38 helium-3 originated from the moon rather than the sun, and can be increased to thousands times.
39 These ions can move downstream far from the lunar. It will cause the ion void of the lunar
40 ionosphere and the escape of lunar material.

41 1 Introduction

42 The photoionization, surface adsorption, and thermal escape are the primary loss
43 mechanisms of the atomic species from the lunar exosphere (Stern, 1999; Killen and Ip,
44 1999; Wurz et al. 2007). These loss processes are a relatively long and slow process. The rapid
45 loss process and mechanism of the exosphere have not been reported. The moon has a tenuous
46 exosphere and no intrinsic magnetic field. Helium is one of the main components in the lunar
47 exosphere with surface densities of $2\sim 7\times 10^4/\text{cm}^3$ (Killen and Ip, 1999; Hodges, 1973, Stern et al,
48 2012, Tirtha et al, 2017). Besides photoionization is regarded as the dominant loss process to
49 helium, the double ionization of helium by the electrons impact can occur once the energy of the
50 injected electrons comes up to 90.2eV (Shah, et al. 1988). Lunar surfaces exposed to the ambient
51 plasma, charged energetic particles, and UV radiation will become charged and develop a sheath
52 with an electric field normal to the surface. Energy electron injections can induce the lunar
53 surface charging to large negative potentials during SEP or BEE (Halekas, et al. 2009; Wang, et
54 al. 2012, 2016). Newborn ions can be fully pickup by the solar wind Alfvén waves via the non-
55 resonance stochastic heating process (Yoon and Wu, 1991; Wang et al. 2009, 2011, Sun et al.
56 2014; Liu et al. 2014). In the non-resonance stochastic heating process, ions are picked up by
57 Alfvén wave in several gyro-periods and are heated in the process. The acceleration of low
58 charge-to-mass will be more efficient. The ion velocity distribution becomes gradually isotropic
59 when it catches up with Alfvén speed by stochastic heating. This means that ions in lunar space
60 outside the sheath could be fully pickup by the solar wind in the presence of Alfvén waves. It has

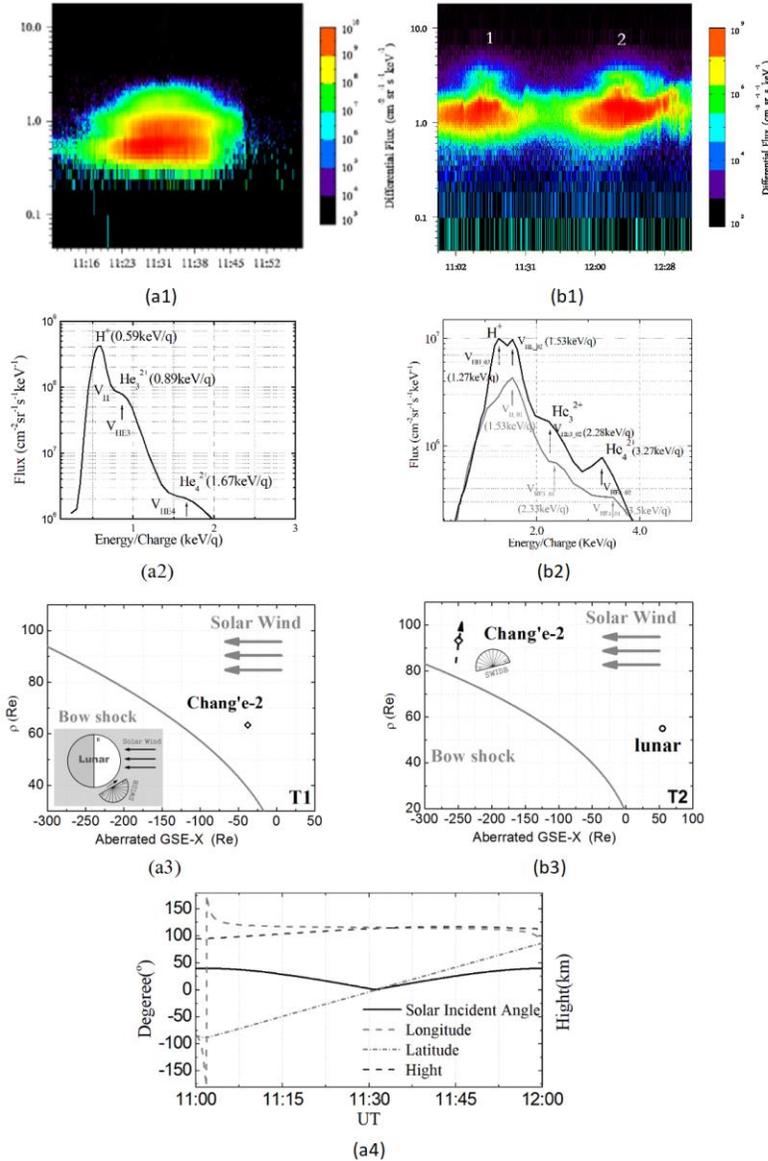
61 been previously observed that the ions picked up by the solar wind from the moon can form a
62 partial or incomplete shell, and the phenomenon of fully pickup ions has not been observed.

63 Launched on Oct.1, 2010, Chang'E-2, as the second lunar orbiter of China, successfully
64 entered a 100km altitude circular lunar orbit after a cislunar transfer trajectory and lunar capture
65 phase. Then during a half-year flight in lunar orbit, Chang'E-2 completed the transfer to the Sun-
66 Earth Lagrange L2 point and entered into the L2 point Lissajous orbit in August 2011 (Huang et
67 al., 2012; Liu et al., 2014). On Chang'E-2, SWIDs and HPD are the backups of the instruments
68 on Chang'E-1 (Wang et al. 2012; Kong et al.2011). SWIDs are designed to observe the
69 0.05~20keV/q solar wind ions, the energy sweep step is 48 with the resolution of 8.5%
70 (FWHM). HPD is designed to observe the energetic charged particles respectively. SWID
71 utilizes a typical top-hat electrostatic analyzer (ESA) with a fan-shaped field-of-view (FOV) of
72 about $180^{\circ} \times 6.7^{\circ}$. The ESA is separated into 12 apertures defined as C1 to C12. The HPD is
73 outfitted with a sensor head composed of three silicon semiconductors. Before the flight, SWIDs
74 were calibrated by the ion beam facility at IRAP (former CESR of Toulouse in France, Institut
75 de Recherche en Astrophysique et Planétologie). HPD was calibrated on a tandem accelerator at
76 CIAE (China Institute of Atomic Energy) and cyclotron at IMP (Institute of Modern Physics,
77 Chinese Academy of Sciences).

78 On Oct.17, 2010, in the lunar orbit and on Sep.27, 2011, at the L2 point, solar wind ${}^3\text{He}^{2+}$
79 extreme enhancement and the 0.1~2MeV electrons increasing were almost synchronized
80 observed by Chang'E-2. In this paper, we analyze the phenomenon and find that the most
81 feasible origins of the extremely enhanced ${}^3\text{He}^{2+}$ lie in the charged particles inducing surface
82 charging and the ionization of the neutrals. Ions can concentrate on the surface by the sheath
83 electric field and fully pickup by the Alfvén wave of the solar wind. This may lead to the rapid
84 loss of helium in the lunar exosphere and the formation of ion void.

85 **2 Observations**

86 Distinct ${}^3\text{He}^{2+}$ fluxes peaks of solar wind ions observed by C12 of SWIDB on Oct.17, 2010,
87 UT 11:00-12:00 (T1) and Sep.27, 2011, UT 10:45-12:40 (T2) are shown in Fig.1. And the other
88 phenomenon as 0.1 ~ 2MeV electrons increase and solar wind ${}^4\text{He}^{2+}$ acceleration were also
89 observed in Fig.2-3. During T1, Chang'E-2 was in a polar orbit with an altitude of 90-120 km
90 and was located at (-50.9RE, 37.6RE, 5.4RE)GSE on the lunar dayside and near side. In T2,
91 Chang'E-2 moved to Lagrange-2 point, and the spacecraft traveled to (-243.2RE, 108.0RE,
92 7.0RE)GSE, downstream of the lunar. For the limited power source, the observation time
93 interval of SWIDs on Sep.27, 2011, just lasted for 2 hours and 37 minutes. The solar incident
94 angle between the symmetric axis of the measurement channel C12 and the Sun-Moon
95 connection line is about 23° , as shown in Fig.1 (b3). Based on the model of the Earth's distant
96 bow shock (Bennett et al. 1997), Chang'E-2 was in the solar wind on both T1 and T2. The
97 schematic diagram in Fig. 1(a3), (b3) demonstrates the plasma regimes and the observation
98 geometry of SWIDB in aberrated GSE coordinate. The longitude, latitude, and altitude of
99 Chang'E-2 in the selenocentric solar ecliptic (SSE) coordinate system of T1 are shown in Fig.1
100 (a4).



101
 102 **Figure 1.** Observations of solar wind fluxes peaks spectrum of H⁺, ³He²⁺, ⁴He²⁺ at UT 2010-10-17 11:00-12:00
 103 (T1, a1-a4) and UT 2011-9-27 10:55-12:40 (T2, b1-b3) by C12 of SWIDB. (a1, b1) the time-energy spectrogram of
 104 solar wind ions in T1 and T2. The vertical coordinate is Energy/Charge and the horizontal is universal time. (a2, b2)
 105 the energy-flux profiles of Solar wind ion in T1 and T2. The subscript 01 and 02 in (b1) identify the sequential
 106 number for the periodical precession of Chang'e-2. (a3, b3) the observation geometry of SWIDB and the location of
 107 the spacecraft during T1 and T2, (a4) the solar incident angle of C12, longitude, latitude, and altitude of the
 108 Chang'e-2 in selenocentric solar ecliptic (SSE) coordinate system in T1.

109 **2.1 Extreme Enhancement of ³He⁺⁺ in the solar wind (EEH)**

110 The obvious differential flux peaks of H⁺, ³He²⁺, and ⁴He²⁺ are shown in Fig.1 (a2) and Fig.1
 111 (b2) and are signified by arrows. There are one proton beam and two alpha beams of ³He²⁺ and
 112 ⁴He²⁺ in Fig.1 (a2) and the first period of Fig.1 (b2). But two proton beams and two alpha beams
 113 of ³He²⁺ and ⁴He²⁺ exist in the second period in Fig.1 (b2). The bulk speed of H⁺, ³He²⁺, and
 114 ⁴He²⁺ beams are defined as V_H, V_{HE3}, and V_{HE4} in Fig. 1(a2). In Fig. 1(b2), V_{H_01}, V_{HE3_01}, and
 115 V_{HE4_01} flux peaks in the gray line of the first period indicate the bulk speed of H⁺, ³He²⁺, and

116 ${}^4\text{He}^{2+}$ beams. The H^+ beams and two alpha beams of ${}^3\text{He}^{2+}$ and ${}^4\text{He}^{2+}$ are referred to as V_{HL_02} ,
 117 V_{HH_02} , V_{HE3_02} , and V_{HE4_02} in the second period with the black line.

118 The penultimate particle peak is verified as ${}^3\text{He}^{2+}$ because of the $M/Q=1.5$ amu/e in Fig1.
 119 The 1-8 orbits H^+ and ${}^4\text{He}^{2+}$ observation result from 0:00-16:00 UT on 2010-10-17 around T1
 120 (the 6th orbit) was shown in Fig 2. In Fig2(c), the peak of the last particle is verified as ${}^4\text{He}^{2+}$
 121 because the $M/Q=2$ amu/e on 1-3th orbits. ${}^4\text{He}^{2+}$ velocity is divided by $(M_{\text{He4}}/Q_{\text{He4}})^{1/2} = (2)^{1/2}$ and
 122 ${}^3\text{He}^{2+}$ velocity is divided by $(M_{\text{He3}}/Q_{\text{He3}})^{1/2} = (1.5)^{1/2}$. In Fig2. (b), the ${}^4\text{He}^{2+}$ acceleration and
 123 moderated to the proton P1 speed is observed. In T1, the ratio of Energy/Charge between the first
 124 two peaks of H^+ , ${}^3\text{He}^{2+}$ is factor 1:1.5 with the V_{H} and V_{HE3} of $\sim 336\text{km/s}$. The solar wind bulk
 125 speed by ACE/SWEPAM is $\sim 331\text{km/s}$ as V_{H} in T1 through a time shift method of about 1h
 126 delay. Moreover, the V_{HE4} for ${}^4\text{He}^{2+}$ of $\sim 400\text{km/s}$ approximates 64 km/s higher than the V_{H} and
 127 V_{HE3} . In the two periods of T2, the FOV of C12 is derived from the solar wind center of about 23
 128 $^\circ$. For the first period, the ratio of Energy/Charge between the proton peak and ${}^3\text{He}^{2+}$ peak is
 129 factor 1.5:2 with the V_{H_01} and V_{HE3_01} of $\sim 541\text{km/s}$, while the V_{HE4_01} for ${}^4\text{He}^{2+}$ of $\sim 570\text{km/s}$
 130 approximates 29km/s higher than the V_{HE3_01} and V_{H_01} . During this period, the average solar
 131 wind bulk speed is $\sim 600\text{km/s}$ by ACE/SWEPAM through a time shift method of about 1.5h
 132 delay. In the second period, the Energy/Charge of the middle two peaks of high-speed H^+ and
 133 ${}^3\text{He}^{2+}$ is factor 1:1.5 with the V_{HH_02} , and V_{HE3_02} of $\sim 540\text{km/s}$, which exceeds the low-speed
 134 proton beam's velocity V_{HL_01} of $\sim 493\text{km/s}$ by about 57 km/s. The V_{HE4_02} for ${}^4\text{He}^{2+}$ of
 135 $\sim 560\text{km/s}$ approximates 20 km/s higher than the V_{HE3_02} and V_{HH_02} . For this period, the average
 136 solar wind bulk speed is $\sim 567\text{km/s}$ by ACE/SWEPAM.

137 We calculate the ratio of phase space density for ${}^4\text{He}^{2+}/\text{H}$, ${}^3\text{He}^{2+}/\text{H}$, and ${}^3\text{He}^{2+}/{}^4\text{He}^{2+}$ during
 138 these two events. The ${}^4\text{He}^{2+}/\text{H}$ ratio is ~ 0.01 . It is similar to the average solar wind value. The
 139 ${}^3\text{He}^{2+}/\text{H}$ ratio is ~ 0.03 and ${}^3\text{He}^{2+}/{}^4\text{He}^{2+}$ ratio is ~ 3 . ${}^3\text{He}^{2+}/{}^4\text{He}^{2+}$ ratio equating to about 6000
 140 times the average solar wind value of $\sim 4.9 \times 10^{-4}$ (Ansgar et. al 2008; Robert et al.,2002; P.
 141 Bochsler et al.,1990; George et al. 1998). In T1, the minimum solar incident angle is about 0° ,
 142 and the density of ${}^3\text{He}^{2+}$ is $\sim 0.01\text{cm}^{-3}$. Thus, we call these performances as Extreme
 143 Enhancement of ${}^3\text{He}^{2+}$ events (EEH).

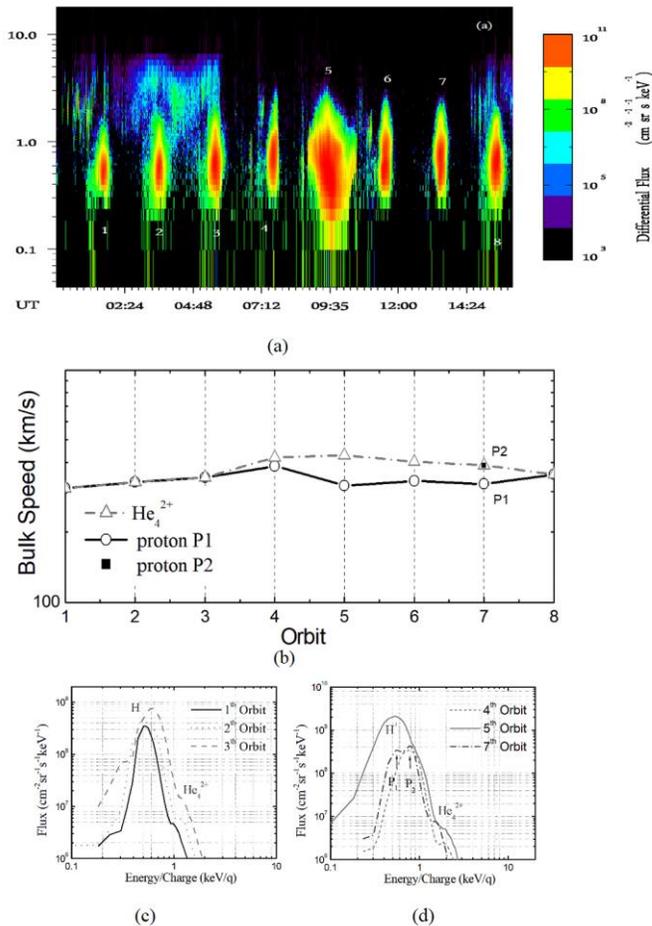
144 In T1 and T2, the thermal speed/solar wind speed is ~ 0.15 and ~ 0.14 separately. Despite the
 145 different solar wind conditions for the two EEH, the core distribution of the thermal is broad.
 146 And no obvious ${}^4\text{He}^+$ observed in these events.

147 2.2 The turbulence of the space environment

148 The weekly highlights of Space Weather Prediction Center (SWPC) reported that a
 149 transient-like feature was observed at ACE on October 17 between 0111 UTC and 1150 UTC as
 150 B_z went through an extended negative period. The peak negative values around -7nT when a
 151 possible source of the slow CME was observed on October 14. And a group of super-thermal
 152 ions from UT01:20 to 06:10 were observed simultaneously by Chang'E-2 at 1-3th orbit in Fig.2
 153 (a). The bulk speed variations of the protons and ${}^4\text{He}^{2+}$ are unveiled in Fig.2 (b). ${}^4\text{He}^{2+}$
 154 acceleration is commencing on the 4th orbit. The difference of bulk velocity between the proton
 155 beam and ${}^4\text{He}^{2+}$ beam was 34-110km/s in 4-6 orbit. In Fig.2 (a), the acceleration of the ions
 156 occurred at 10:10 and lasted until the FOV of C12 was far from the solar wind center. The ions'
 157 central energy increased at least 0.4keV in 20 minutes. The EEH of T1 that occurred in the 6th
 158 orbit is mentioned in section 2.1. The energy-fluxes profiles of the protons and ${}^4\text{He}^{2+}$ in 4, 5, 7th
 159 orbit can be seen from Fig.2(c). In the 5th orbit, the energy-fluxes profiles broadened in the polar
 160 direction, the total flux increased and the temperature does not change significantly. The double

161 proton beams for P1 and P2 arose in the 7th orbit. The bulk speed of P2 in sync with ${}^4\text{He}^{2+}$ and
 162 the velocity difference between the two proton beams is $\sim 66\text{km/s}$. In the 8th orbit, the solar wind
 163 proton beam reverted to a single peak with the coincidence of bulk speed of ${}^4\text{He}^{2+}$ and proton.

164 A SEP began on Sep. 23, 2011, was reported by SWPC and observed by HPD. In Fig3.(b),
 165 the flux peaks of $>10\text{MeV}$ protons reached to $40 (\text{cm}^2 \cdot \text{sr} \cdot \text{s} \cdot \text{MeV})^{-1}$. For T2, an ICME without the
 166 magnetic cloud, starting from 20:00 on Sep.26 and ending at 15:00 on Sep.28, was publicized in
 167 the ICME tables list by Richardson and Cane ([http://www.srl.caltech.edu/ACE/](http://www.srl.caltech.edu/ACE/ASC/DATA/level3/icmetable2.htm)
 168 [ASC/DATA/level3/icmetable2.htm](http://www.srl.caltech.edu/ACE/ASC/DATA/level3/icmetable2.htm)). During T2, the obvious acceleration of ions was observed
 169 at 12:22~12:32 in Fig.1 (b1) and Fig.3 (d). The ions sustained acceleration until the FOV of the
 170 C12 far from the solar wind center, and the peak energy of the ions increased $\sim 0.46\text{keV}$.



171 **Figure 2.** The observation result of C12 measurement channel of SWID-B from 0:00-16:00 UT on 2010-10-17.
 172 1-8 is the sequential number for Chang'E-2 orbiting the lunar and indicates the spacecraft orbits the lunar for 8
 173 times. (a) Energy/Charge-Time spectrogram of C12. (b) The bulk speed of the solar wind ${}^4\text{He}^{2+}$ and protons in each
 174 of the orbits. The curve with circles, squares, and triangles are the bulk speed of protons and ${}^4\text{He}^{2+}$. P1 and P2 show
 175 the velocity of double proton beams in the 7th orbit. (c) Energy/Charge-Fluxes profiles of Solar wind ions in 1, 2, 3rd
 176 orbit. (c) Energy/Charge-Fluxes profiles of Solar wind ions in 4, 5, 7th orbit.

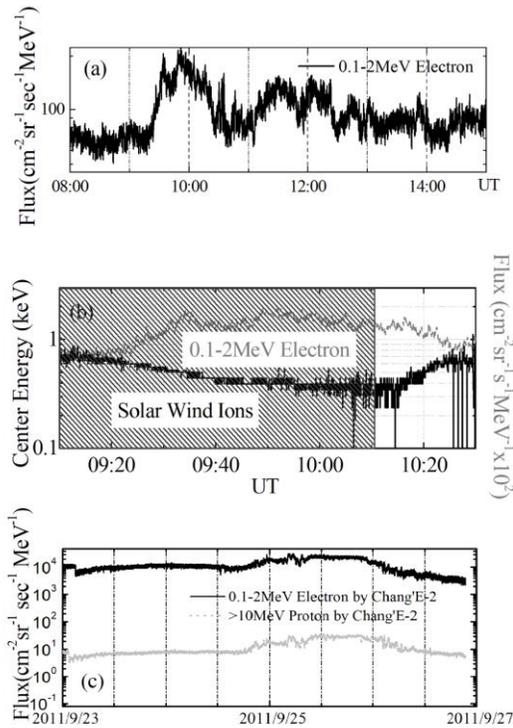


Figure 3. 0.1-2.0MeV keV energetic electrons increase and the acceleration of solar wind ions observed by HPD/Chang'E-2. (a) The rising of the energetic electrons manifests a BEE event that occurred on 2010-10-17 9:17 and lasted for about 5 hours. (b) In the 5th orbit on 2010-10-17, the 0.1-2.0MeV electron fluxes peaked at 9:50 and subsequently increased at least 0.4keV of the ions' central energy. (c) The fluxes of energetic electrons increased greatly during SEP on 2011-9-23 16:00 ~2011-9-27 16:00. (d) The peak energy of the ions increased ~0.46keV on 2011-9-27 12:26~12:36.

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2.3 Increasing of Energetic electrons and lunar surface charging

The rise of 0.1-2.0MeV energetic electrons was observed by HPD/Chang'E-2 during the two EEH in Fig3. In T1, the bursting of the energetic electrons indicated the occurrence of a BEE event (Wang et al. 2012, 2011). The energetic electrons during BEE are directional obviously. However, in T2, the fluxes of energetic electrons surged is due to the solar energetic events (SEP). The general magnitude of the ubiquitous background 0.1~2.0MeV electrons fluxes observed by HPD was $\sim 10^1 (\text{cm}^2 \cdot \text{sr} \cdot \text{s} \cdot \text{MeV})^{-1}$. On Oct.17, 2010, 9:17 UT, 0.1~2.0MeV energetic electrons injected impulsively in Fig.3 (a) and sustained high flux for about 5 hours with the flux peaks reaching the magnitude of $\sim 10^2 (\text{cm}^2 \cdot \text{sr} \cdot \text{s} \cdot \text{MeV})^{-1}$. For T2, the 0.1~2.0MeV energetic electrons injected sustained high flux with the peak fluxes reaching the magnitude of $\sim 10^4 (\text{cm}^2 \cdot \text{sr} \cdot \text{s} \cdot \text{MeV})^{-1}$.

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The increasing of energetic electrons suggested the EEHs were subsequent with the lunar and spacecraft surface charging. Wang et al. (2012, 2016) demonstrated that BEE can induce circumlunar spacecraft and the lunar surface charging to negative kilovolts when the temporal integral of the energetic electrons fluxes exceeds 10^{11}cm^{-2} . And the charging process is sustained for the same interval as the BEE. For T1, we calculated the surface potentials of the lunar utilizing the simulation method by Wang et al. (2012, 2016). The calculation result shows that the temporal integral of the electrons fluxes can reach up to $>10^{11} \text{cm}^{-2}$ and the lunar surface charging voltage can reach up to more than -0.75kV . The result is consistent with the magnitude

205 of at least 0.4keV increase of solar wind ions energy as mentioned in section 2.2. For the lunar
 206 surface potentials is much more negative than kT_e/e , an ion matrix or transient sheath forms
 207 around the surface and expands to the solar wind. The total sheath thickness (S_C) can be obtained
 208 from the quasistatic Child–Langmuir law

$$209 \quad S_C = \frac{\sqrt{2}}{3} \left(\frac{2eU_0}{kT_e} \right)^{3/4} \lambda_D.$$

210 In the above expressions, λ_D is the Debye length, T_e is the electron temperature, U_0 is
 211 negative surface bias. According to the above formula, sheath thickness S_C is at least 30 meters.

212 For T2, SEP occurred with the fluxes of 0.1~2MeV electrons reaching to $>10^3$ (cm²
 213 ·sr·s·MeV)⁻¹ and sustained for more than 4 days from Sep.23, 2011, to Sep.27, 2011. Halekas et
 214 al. (2009) proposed that the spacecraft and lunar surface potentials can reach negative kilovolts
 215 values during SEP. This indicates that in T2, the lunar surface charging potentials can reach up to
 216 a large negative voltage. It was confirmed by the 0.46keV increment of ion energy in section 2.2.
 217 So the lunar surface potentials amplitude of variations can reach up to -0.46kV.

218 **3 Origin, ionization, and pickup of ${}^3\text{He}^{2+}$ in EEHs**

219 **3.1 The origin of ${}^3\text{He}^{2+}$ enhancements**

220 The abnormal increase of helium-3 in the two EEHs can originate from the sun or the moon.
 221 In these events, ${}^3\text{He}^{2+}/{}^4\text{He}^{2+}$ ratios elevated by a factor of 6000 exceedingly surpass those in early
 222 reports of keV ${}^3\text{He}^{2+}$ enhancement events (EHs) in the solar wind. In EHs, the factors are 4~100
 223 than those in the average solar wind (Ho, 1998; Gloeckler, 1999; Wittenberg, 2000). The EHs
 224 are associated with CME passage and ${}^3\text{He}^{2+}$ stems from the solar. Most of the EHs were within
 225 or trailing the magnetic clouds. An overabundant ${}^4\text{He}^+$ was observed in these EHs. It implied that
 226 because of the expansion and cooling of the CME transit from the Sun, the temperature in EHs is
 227 much lower than ($<10^5$ K) the typical million-degree freeze observed in the normal solar wind.
 228 However, an ICME without magnetic cloud was reported during T2 and none of the ICME
 229 occurred during T1. Unlike EHs, none of the noticeable ${}^4\text{He}^+$ was observed in EEHs. The
 230 temperature is of $\sim 2 \times 10^5$ K in T1 and $\sim 3 \times 10^5$ K in T2 confirmed that the core distributions of the
 231 thermal plasma are not as narrow as in CME. Extra high fluxes of ${}^3\text{He}^{2+}$ ions, the absence of
 232 ${}^4\text{He}^+$, and high kinetic temperature suggest the different origins of solar wind ${}^3\text{He}^{2+}$
 233 enhancements between the EHs and EEHs.

234 In addition to the origin of the sun, there are two mechanisms of the lunar origins. The first
 235 is from the helium ionization of the lunar exosphere. Since the majority of the lunar exosphere
 236 Helium is of the solar wind origin, it is known to scale with the solar wind alpha particle flux
 237 (Hurley et al., 2016). The diurnal variation of Helium is at around $\sim 2 \times 10^3/\text{cm}^2$ to $\sim 4 \times 10^4/\text{cm}^2$
 238 (Hodges, 1975). CHACE/MIP /Chandrayaan-1 evidenced that the upper limit of neutral helium
 239 density can reach $8 \times 10^2/\text{cm}^2$ in the sunlit when the Helium abundance in the Moon had hit one of
 240 its lowest values (Tirtha et al., 2017). Because of the observations of Helium rich solar wind in
 241 two EEHs, we assume that the concentration of helium in the exosphere can reach the magnitude
 242 of $10^2 \sim 10^4/\text{cm}^2$. The double ionization of helium-3 can occur once the energy of the injected
 243 electrons $>90.2\text{eV}$. The cross section is $\sim 10^{-20} \text{cm}^2$ for the helium double ionized by $>1\text{keV}$
 244 electrons (Shah et al. 1998). The trend of surface charge to negative potentials is consistent with
 245 that of energetic electron increasing. Combined with the average ${}^3\text{He}/{}^4\text{He}$ ratio, the total amount
 246 of double ionized ${}^3\text{He}^{2+}$ in the sheath with a thickness of 30 meters and an area of 1cm^2 is
 247 around 0.0002~0.02. The produced ions will converge to the lunar surface in the negative sheath

248 electric field. Since the lunar regolith is an electrical insulator, deposited charged particles can
 249 remain in the lunar surface regolith for prolonged intervals. When the lunar surface potential is
 250 reversed, the ions deposited on the surface will be modulated by the solar wind electromagnetic
 251 field and pick up by the solar wind.

252 There is another possible loss mechanism by energetic particles penetrate the lunar regolith
 253 cause deep dielectric charging and breakdown(Jordan et al. 2014, 2019). Campins and Krider's
 254 experiments suggest that dielectric breakdown can melt and boil off material. The breakdown
 255 can occur when the difference in fluxes between protons and electrons is at least $\sim 10^{10} \text{ cm}^{-2}$. It
 256 can be derived that the helium-3 released from the area of 1 cm^2 is about several to hundreds of
 257 particles in the two EEHs. According to the collision cross section, the ionized helium-3 is much
 258 lower than the observations in the EEHs. Therefore, the deep dielectric charging and breakdown
 259 in lunar regolith is not enough to comprehend the abnormal increase of helium-3 in the solar
 260 wind. These results indicate that the primary source of abnormal helium 3 in this paper should be
 261 the ionization of the exosphere.

262 3.2 ${}^3\text{He}^{2+}$ ions fully pickup by the solar wind

263 The double proton beams and the acceleration of the solar wind helium ions mentioned in
 264 sections 2.1-2.2 suggested the existence of Alfvén wave during T1 and T2. For T1, The IMF $|B|$
 265 value for the orbit 4-7th in Fig.2 is 5-10nT from ACE by the time-shift method and the solar wind
 266 density is 3.0 cm^{-3} by Chang'E-2. Accordingly, the local Alfvén velocity was calculated as
 267 $63\sim 126 \text{ km/s}$. It is almost consistent with the variation of velocity difference between the ${}^4\text{He}^{2+}$
 268 and proton P1 in Fig.2 (b). For the orbit 2-3th, the difference between the center speed of the
 269 super-thermal ions and the bulk speed of the proton in Fig.2 (a) was $\sim 120 \text{ km/s}$. In the 7th orbit,
 270 the high- speed proton beam P2 appeared with the same speed as ${}^4\text{He}^{2+}$. Then in the 8th orbit, the
 271 single proton beam recovered with the same velocity as ${}^4\text{He}^{2+}$. It suggests that the Alfvén wave
 272 has existed in the solar wind for more than 3 hours before T1. Although the resonance conditions
 273 were not satisfied because the initial bulk speed of ${}^3\text{He}^{2+}$ is of ~ 0 , ${}^3\text{He}^{2+}$ ions can be picked up by
 274 the solar wind via non-resonance interaction and sustaining stochastic heating due to the
 275 existence of the Alfvén wave. It is generally known that solar wind Alfvén waves originate from
 276 the sun. Solar wind proton and helium-4 were accelerated by Alfvén wave except for helium-3
 277 ions, which suggests that helium-3 is not originating from the solar.

278 For T2, the spacecraft was located in the solar wind downstream distant to the lunar.
 279 Although the FOV has deviated from the solar wind center, the ${}^3\text{He}^{2+}$ and ${}^4\text{He}^{2+}$ beams of ~ 57
 280 km/s and $\sim 25 \text{ km/s}$ higher than the low-speed proton beam's velocity were observed in the two
 281 periods respectively. It is indicating the existence of Alfvén wave and consistent with the non-
 282 resonance stochastic heating process, where the velocity of the solar wind pickup ions will
 283 increase by an Alfvén speed after enough time.

284

285 4 Discussion and conclusions

286 Two extreme enhancements of ${}^3\text{He}^{2+}$ by thousands of times average solar wind value and
 287 the $0.1\sim 2 \text{ MeV}$ electrons bursting are observed by SWIDs and HPD when Chang'E-2 orbits
 288 around the moon and at L2 point in the solar wind. The two EEHs occurred when the spacecraft
 289 was in the solar wind around the lunar and the lunar downstream respectively. The difference
 290 from EHs indicates the helium-3 in EEHs is not originated from the sun. Both the lunar

291 exosphere and lunar regolith are rich in helium, the lunar should be the potential candidate to be
292 the provenance. The dielectric breakdown caused by energetic particles penetrates the lunar
293 regolith can release outgases to the exosphere. Double ionization of helium-3 will occur by the
294 $>90.2\text{eV}$ energetic electrons injection. Energetic electrons increasing can induce the lunar
295 surface charging to negative kilovolts in the sheath of >30 meters. The ions in the sheath will
296 converge to the surface by the electric field. When the lunar surface potential is reversed, the
297 ${}^3\text{He}^{2+}$ ions distributed on the lunar surface will be pickup by the turbulent solar wind
298 alternatively. The speed difference of the solar wind proton and helium beams in both EEHs
299 identified the existence of Alfvén wave. Helium-3 deriving from the lunar can be fully picked up
300 by the solar wind in the non-resonance stochastic heating process and traveling downstream to
301 L2 point. In the two events, the calculation shows that helium-3 produced by dielectric
302 breakdown is much lower than in EEHs. The maximum surface density for ${}^3\text{He}^{2+}$ produced in the
303 exosphere can reach 0.02 cm^{-2} . So the primary lunar source of solar wind abnormal helium-3 in
304 the two EEHs should be the lunar exosphere.

305 In this paper, we find only two EEHs, which is much lower than the probability of BEE.
306 Firstly, the SWIDs were unable to distinguish M/Q in the isotropic plasma; secondly, the
307 different processes of ions with diverse mass numbers or charges, such as the acceleration by the
308 electric field, and the non-resonance stochastic heating process, leading to their fractionation.
309 And for the lack of magnetic field data, it's difficult to deduce the exact local positions of the
310 Helium-3 deriving from the lunar or to confirm the concrete pickup process. The results indicate
311 the occurrence of ions void in the lunar exosphere. Since the frequent occurrence of BEEs can
312 lead the surface charging to large negative potentials and the neutrals ionizing, the phenomenon
313 of ion void in the exosphere will be recurrent. It will cause the instability of the lunar ionosphere
314 and the loss of lunar material. Moreover, because the lunar surface potentials can be charged to
315 large negative potentials and last a long time during SEP, the speculation can be made that the
316 neutrals ionizing and loss in the lunar exosphere could arise frequently in the high-level solar
317 activity period.

318 Future more, the detection of the electromagnetic field, neutrals, and charged particles on
319 the following Chang'E series could offer us the chance to evidence the effects of energetic
320 electron injection in the lunar space environment and verify the lunar Ionospheric disturbances.

321

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332 **References**

- 333 Deng, A., & Stauffer, D. R. (2006), On improving 4-km mesoscale model simulations. *Journal*
 334 *of Applied Meteorology and Climatology*, 45(3), 361–381. doi:10.1175/JAM2341.1
- 335 Ansgar Grimberg, Heinrich Baur, Fritz Buhler, Peter Bochsler, & Rainer Wieler. (2008), Solar
 336 wind helium, neon and argon isotopic and elemental composition: Data from the metallic glass
 337 flown on NASA's Genesis mission. *Geochimica et Cosmochimica Acta.*,72, 626-645
- 338 Stern S A., (1999),The lunar atmosphere: history, status, current problems, and context. *Reviews*
 339 *of Geophysics*, 37:453-491
- 340 Bennett L., Kivelson M. G., Khurana K. K., Frank L.A., & Paterson W.R., (1997), A model of
 341 the Earth's distant bow shock. *Journal of Geophysical Research*, 102 (A12), 26927-26941
- 342 Bochsler P., Geiss J., & Maeder A., (1990), The abundance of ³He in the solar wind a constraint
 343 for models of solar evolution, *Solar physics*. 128,203-215
- 344 Campins, H., & E. P. Krider. (1989), Surface discharges on natural dielectrics in the solar
 345 system, *Science*, 245, 622–624, doi:10.1126/ science.245.4918.622
- 346 George Gloeckler, & Johannes Geiss, (1998), Measurement of the abundance of helium-3 in the
 347 sun and in the local interstellar cloud with SWICS on ULYSSES. *Space Science Reviews*,
 348 84:275-284
- 349 Gloeckler G., Fisk L. A., Hefti S., Schwadron N. A., Ipavich, T.F. M. Geiss J., et al., (1999),
 350 Unusual composition of the solar wind in the 2-3 May 1998 CME, observed with SWICS on
 351 ACE. *Geophys. Res. Lett.* 26(2), 157-160
- 352 Halekas, J. S., Delory G. T., Lin R. P., Stubbs T. J., & Farrell W. M. (2009), Lunar surface
 353 charging during solar energetic particle events: Measurement and prediction. *J. Geo. Res. Let.*,
 354 114, A05110
- 355 Hurley, D.M., Cook, J.C., Benna, M., Halekas, J.S., Feldman, P.D., Retherford, K.D., et al.,
 356 (2016), Understanding temporal and spatial variability of the lunar helium atmosphere using
 357 simulta- neous observations from LRO, LADEE, and ARTEMIS. *Icarus*, 273, 45–52
- 358 Ho, C. W. (1998), Helium-3 enhancements and unusual ion charge state composition in coronal
 359 mass ejections. *University of Maryland College Park*.
- 360 Hodges R.R., & JR, J., (1973), Helium and Hydrogen in the Lunar Atmosphere, *Geo. Res.*,
 361 78(34), 8055-8064
- 362 Hodges, R.R., (1975), Formation of the lunar atmosphere. *Moon*, 14, 139–157
- 363 Huang H., (2012), chang'e-2 satellite Lagrange L2 point Mission. *63th International*
 364 *astronautical Congress*, Naples, Italy
- 365 Jordan A. P., Stubbs T. J., Wilson J. K., Schwadron N. A., Spence H. E., Joyce C. J., (2014),
 366 Deep dielectric charging of regolith within the Moon's permanently shadowed regions, *J.*
 367 *Geophys. Res. Planets*, 119, 1806–1821
- 368 Jordan A. P., Stubbs T. J., Shusterman M.L., Izenberg N.R., Wilson J. K., Hayne P.O. et al,
 369 (2019), How dielectric breakdown may contribute to the global weathering of regolith on the
 370 moon, *Icarus*, 319, 785–794

- 371 Killen R., & Ip W.H., (1999), The surface-bounded atmospheres of Mercury and the Moon, *Rev.*
372 *Geophys.* 37(3), 361-406 .
- 373 Kong L. G., Wang S. J., Wang X. Y., Zhang A.B., Zhu G.W., Yu D.J. et al., (2011), In-flight
374 Performance and Preliminary Observational Results of Solar Wind Ion Detectors (SWIDs) on
375 Chang'E-1. *Planetary and space Science*, 62(1),23-30.
- 376 Liu H. F., Wang S. Q., Li K. H., & Tang C. J., (2014), Nonresonant and stochastic heating of
377 ions by low-frequency wave. *Physics Letters, A*. 378, 3614–3616
- 378 Liu L., Liu Y., Cao J. F., Hu S.J. Tang G.S. & Xie J. F., (2014), Chang'e-2 lunar escape
379 maneuvers to the sun-Earth L2 libration point mission. *Acta Astronautica.*, 93, 390-399
- 380 Robert F.Wimmer-Schweingruber. (2002),The composition of the solar wind. *Adv Space Res.*,
381 30(1),23-32
- 382 Shah M. B., Elliott D. S., McCallion P., & Gibody H. B., (1988), Single and double ionization of
383 helium by electron impact, *J. Phys. B: At. Mol. Opt. Phys.*, 21, 2751-2761
- 384 Stern, S.A., Retherford, K.D., Tsang, C.C.C., Feldman, P.D., Pryor, W., & Gladstone, G.R.,
385 (2012), Lunar atmospheric helium detections by the LAMP UV spectrograph on the Lunar
386 Reconnaissance Orbiter. *Geophys. Res. Lett.*, 39, L12202. doi:10.1029/2012GL051797.
- 387 Sun J. C., Gao X. L., Lu Q. M., & Wang S., (2014), The Efficiency of Ion Stochastic Heating by
388 a Monochromatic Obliquely Propagating Low-Frequency Alfvén Wave. *Plasma Science and*
389 *Technology*, 16(10)
- 390 Tirtha P. D., Thampi S. V., Dhanya M.B., Anil B., Ahmed S.M., & Sridharan R., (2017), Upper
391 limit of helium-4 in the sunlit lunar exosphere during magnetotail passage under low solar wind
392 condition: Result from CHACE aboard MIP in Chandrayaan-1, *Icaru* , 297, 189–194
- 393 Wang B., & Wang C. B., (2009), Heating rate of ions via nonresonant interaction with turbulent
394 Alfvén waves with ionization and recombination. *Physics of plasmas*, 16, 082902
- 395 Wang B., Wang C. B., Yoon P. H., & Wu C. S., (2011), Stochastic heating and acceleration of
396 minor ions by Alfvén waves. *Geophys. Res. Lett.*, 38, L10103, doi: 10.1029/2011GL047729
- 397 Wang X. Y., Zhang A. B., Zhang X.G., Reme H., Kong L.G., Zhang S.Y. et al., (2012), Bursts of
398 Energetic Electron Induced Large Surface Charging Observed by Chang'E-1. *Planetary and*
399 *space science*, 71(1), 1-8
- 400 Wang X. Y., Zhang A. B., Jing T., Reme H., Kong L. G., Zhang S. Y., et al., (2016),
401 Synchronization of Energetic Electrons Bursting and the Lunar Orbiter Surface Charging to
402 negative Kilovolts. *Chinese J. Geophys.*, 59(10), doi: 106038/cjg20161001, 2016
- 403 Wittenberg E., Hjections L., George C.H., Douglas C.Hamilton, George Gloeckler, & Peter
404 Bochsle. (2000), Enhanced Solar Wind 3He²⁺Associated with Coronal Mass, *Geophys. Res.*
405 *Lett.*, 27 (3), 309-312
- 406 Wurz P., Rohner U., Whitby J.A., Kolb C., Lammer H., Dobnikar P. et al., (2007), The lunar
407 exosphere: The sputtering contribution, *Icarus*, 191,486-496
- 408 Yoon P. H., & Wu C. S., (1991), Ion Pickup by The Solar Wind Via Wave-particle Interactions.
409 *American Geophysical Union*, 61, 241-258