

# Supporting Information for “Impact of remineralization profile shape on the air-sea carbon balance”

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## Introduction

This Supplementary Information contains supporting materials and methods, figures for individual model simulations presented in aggregated form in the main text, as well as tables containing additional globally-aggregated values and anomalies.

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### Alternative remineralization profile shapes.

Here we outline the derivation and assumptions behind six different remineralization profiles. Assuming timescales of sinking ( $\leq 1$  month) are shorter than transport by ocean circulation ( $\sim 1$  year), biological material can be approximated as instantaneously redistributed and remineralized in the vertical.  $f(z)$  is the fraction of the flux of particulate organic matter from a productive layer near the surface (Buesseler et al., 2020) sinking through the depth horizon  $z$  [m].

The most basic curve is a “simple exponential”, assuming constant first-order remineralization kinetics and velocity (Banse, 1990; Dutkiewicz et al., 2005; Marsay et al., 2015; Gloege et al., 2017):

$$f_e(z) = C_e e^{-\frac{z}{\ell_e}}, \quad (1)$$

where  $C_e$  is a scaling coefficient and  $\ell_e$  [m] is a characteristic lengthscale—the ratio of remineralization timescale and sinking speed.

Including an additional flux of refractory material  $c$  to increase the net sinking flux produces the “ballast” model (Armstrong et al., 2001; Gloege et al., 2017):

$$f_b(z) = C_b e^{-\frac{z}{\ell_b}} + c, \quad (2)$$

while explicitly considering the transformation of labile particles with a characteristic lengthscale,  $\ell_{d1}$ , and more refractory material with a characteristic length scale,  $\ell_{d2}$  results in a “double exponential” profile (Lutz et al., 2002):

$$f_d(z) = C_{d1} e^{-\frac{z}{\ell_{d1}}} + C_{d2} e^{-\frac{z}{\ell_{d2}}}. \quad (3)$$

Relaxing the assumption of constant remineralization timescale, and considering a decrease in the rate of remineralization as labile material is preferentially transformed and refractory material is left behind (as in Eq. 1), leads to a “stretched exponential”:

$$f_s(z) = C_s e^{-z^{(1-s)}}, \quad (4)$$

where  $s$  is a scale factor between 0–1, for example, if marine particles are degrade similarly to marine sediments,  $s \approx 0.95$  (Middelburg, 1989; Cael & Bisson, 2018). A three-parameter stretched exponential with  $z$  normalized by  $z_o$  is used in many applications. However, fitting  $z_o$  and  $s$  simultaneously is ill-conditioned, i.e. parameter values are not identifiable, so we have used the simpler two-parameter stretched exponential function, which still provides fits well within global particle flux uncertainty (see Fig. 1).

Second-order degradation kinetics leads to a rational form (Suess, 1980):

$$f_r(z) = \frac{C_r}{z + a} \quad (5)$$

where  $C_r$  [m] is determined by remineralization and sinking while  $a$  [m] is determined by remineralization, sinking, and the initial flux (Cael & Bisson, 2018).

One can model sinking particles as heterogeneous media containing organic carbon, ballast minerals, and heterotrophic bacteria where remineralization slows with time (Rothman & Forney, 2007). This translates to an upper incomplete “gamma” function curve,  $\Gamma(a, x)$ , of zeroth order (Cael & Bisson, 2018):

$$f_g(z) = C_g \Gamma\left(0, \frac{z}{\ell_g}\right), \quad (6)$$

where  $\ell_g$  [m] relates to sinking speed and intraparticle bacterial concentration (Rothman & Forney, 2007; Cael & Bisson, 2018). The upper incomplete gamma function can take any value for order  $a$ , and is a component of other remineralization profile studies (e.g. Aumont et al 2017, Kriest and Oschlies 2008). For positive integer values of  $a$ , when  $a = 1$  then  $C_g\Gamma(1, z/\ell_g) = C_g e^{-z/\ell_g}$  (which is the simple exponential profile), and if  $a = 2$  then there is a recursion relation where  $C_g\Gamma(2, z/\ell_g) = C_g\Gamma(1, z/\ell_g) + C_g z/\ell_g e^{-z/\ell_g}$ , and so on. We repeated our statistical fit to the reference power law curve for upper incomplete gamma functions with orders of  $a$  between 0–3 (Figure S1). The best fits to the reference power-law ( $b=0.84$ ) are in fact when  $a = 0$  regardless of minimizing absolute or relative errors (AFIT or RFIT), and also in the EFD case. The simple exponential case, equivalent to  $a = 1$ , is a poor fit to the power-law for all three cases, and the dissimilarity increases where  $a > 1$ . Indeed, if  $a = 2$  or  $3$  then attenuation increases with depth in the upper ocean, which is uncharacteristic of sinking particle observations and the Martin Curve, and the existence of a particle class whos flux profile would depend on  $C_g z/\ell_g e^{-z/\ell_g}$  is unlikely. Thus, we do not run simulations where  $a > 1$ , which do not correspond to a plausible mechanism for sinking particle remineralization.

A more general three-parameter upper incomplete gamma function parameterization  $C_g\Gamma(a_g, z/\ell_g)$  fits the Martin curve very well with  $a_g \approx -0.8$  (Figure S1), and would correspond to a constant-sinking reactivity continuum model (Aumont et al., 2017) with a power-law reactivity distribution. However, reactivity continuum models do not describe reactivity using a power law, and instead use lighter-tailed distributions such as the gamma (Boudreau & Ruddick, 1991), beta (Vähätalo et al., 2010), or log-normal distribution

(Forney & Rothman, 2012) that do not permit closed-form solutions. Again, we do not consider this parameterization as a plausible alternative to the Martin Curve for this reason.

Finally, we note that the “power-law” (Eq. 1) assumes either an increase in sinking speed with depth or a decrease in remineralization rate (Kriest & Oschlies, 2011; Cael & Bisson, 2018).

Coefficients derived by non-linear statistical fit of profiles given by Eqn. S1–S6 to the reference power-law curve are given in Table S1.

### Supporting simulation results:

Fig. S2 shows steady-state zonal averages for the power-law simulations where  $b=0.84\pm 0.14$ . Particulate organic carbon fluxes in the ocean interior increase when the power-law exponent decreases to 0.70 (Fig. S2, left column) and decrease when the power-law exponent is increased to 0.98 (Fig. S2, right column). The negative feedback between nutrient availability and biological production of particles (e.g. global experiments in Kriest & Oschlies, 2011) can be seen in the surface export flux anomalies (Fig. S2d, f). More efficient export and lower rates of upper ocean remineralization leads to a decrease in recycled nutrient availability, and therefore less overall biological productivity, and lower shallow particulate organic carbon flux when  $b$  is reduced (also see globally-integrated community production,  $\Delta B_C$ , and integrated export fluxes through the deepest annual mixed layer,  $\Delta E_{mld}$ , values in Fig. S3a and Table S2). However, reduced shallow export across the deepest mixed layer depth is compensated by lower upper ocean flux attenuation, due to reduced exponent value, resulting instead in enhanced export flux anomalies

integrated at 1 km depth. Nutrient availability increases when the remineralization profile is more attenuating in the upper ocean, driving enhanced shallow particulate production and export, but this flux is quickly depleted by the same remineralization profile attenuation, resulting in lower interior ocean export fluxes to the deep ocean.

Biological carbon concentration ( $C_{bio}$ ) integrates these export fluxes, so that when interior ocean export increases, the deep ocean biological carbon store increases (Fig. S3b), particularly in the Southern Ocean (Fig. S2g), and the deep North Pacific (Fig. S2j), but decreases in the North Atlantic (Fig. S2g) and vice versa (Fig. S2f, i, l) (Kwon & Primeau, 2006; Kriest & Oschlies, 2011; Kriest et al., 2012; Romanou et al., 2014).

The deep ocean store of biological carbon is directly linked to air-sea carbon partitioning, thus greater  $\Delta C_{bio}$  indicates uptake of atmospheric carbon by the ocean, and  $p\text{CO}_2$  declines. Conversely, when  $\Delta C_{bio}$  decreases, that carbon outgasses from the ocean causing atmospheric  $p\text{CO}_2$  to increase (Table S2).

Fig. S4 shows zonally-averaged anomalies with respect to the reference power-law of export fluxes for the three sets of parameter values and the six different functional forms of remineralization profile. Anomalies largely have an inverse surface-deep ocean contrast, which is captured by the differences in fluxes through the deepest annual mixed layer depth ( $\Delta E_{mld}$ , Table S3) for the surface ocean changes in particulate export, and fluxes through the 1 km depth horizon ( $\Delta E_{1km}$ , Table S3) for the deep ocean changes in particulate export. However, the ballast functional form (Eq. S2) has a more complex distribution of particulate flux anomalies in surface, intermediate, and deep waters associated with the

additional refractory flux  $c$ , which becomes important when the exponentially decaying labile portion of the sinking flux becomes attenuated to low levels.

Fig. S5–S7 shows Atlantic and Pacific Ocean zonally-averaged anomalies with respect to the reference power-law of the concentration of biological carbon ( $\Delta C_{bio}$ , Table S3), which integrates the vertical flux and remineralization of particulate organic matter. Again, anomalies largely have an inverse surface-deep ocean contrast, with similar sign anomalies in the older upwelling waters of the Southern Ocean and deep North Pacific, in contrast to the youngest waters in the North Atlantic (Kwon & Primeau, 2006; Kriest & Oschlies, 2011; Kriest et al., 2012; Romanou et al., 2014). The deep ocean store of  $C_{bio}$  is inversely proportional relationship with atmospheric  $\text{CO}_2$  content (Table S3).

Fig. S8 illustrates the crucial effect of biological pump nonlinearity in achieving significant atmospheric  $\text{CO}_2$  drawdown. In our NOPOM simulation (Table S2) there are no sinking particle fluxes, and the oceanic biological carbon store is reduced by over 90% compared to the reference power-law profile, which leads to a  $\sim 160\mu\text{atm}$  increase in atmospheric  $\text{CO}_2$  (compare to the  $200\mu\text{atm}$  difference between biotic and abiotic oceans Volk and Hoffert (1985)). This is roughly twice the  $\sim 80\mu\text{atm}$  rise in AFIT and EFD experiments with simple exponential remineralization profiles, even though these simulations have a reasonably strong particulate organic matter fluxes through 1 km (only reduced by  $\sim 20\%$  compared to the reference power-law, see Table S3) and a substantial internal store of biological carbon (also only reduced by  $\sim 20\%$  compared to the reference power-law). Thus, only about half of the biological pump’s influence on atmosphere-ocean carbon partitioning ( $\sim 80\mu\text{atm}$ ) can be attributed to export of particulate organic matter and

biological carbon storage. The remainder must be attributed to differences in the shape of the remineralization profile, and the mechanisms responsible. Simple exponential profiles have a first-order, constant, rate of remineralization, i.e. fixed attenuation of the sinking particle flux with depth, whereas attenuation for the power-law curve decreases significantly with depth. For AFIT and EFD exponential profiles, this difference in attenuation results in the majority of the sinking particle flux from the surface ocean being remineralized in the upper 2 km (export fluxes through this horizon are 0.204 and 0.140 PgC  $y^{-1}$  compared to 0.802 PgC  $y^{-1}$  for the reference power law curve), and therefore there is very little abyssal biological carbon storage to act as a significant reservoir of atmospheric CO<sub>2</sub>. In other words, significantly decreasing upper ocean particle attenuation or increasing remineralization lengthscale with depth appears to be equally important for air-sea carbon partitioning as export and storage of biological carbon in the ocean interior (even when it occurs below 1000 m).

Finally, Fig. S9 shows a comparison of model phosphate fields with the World Ocean Atlas climatology (v2013, Garcia et al., 2014) for the 21 simulations with different remineralization profiles. Despite reorganization of carbon and nutrient concentrations, comparison of model output to climatological data does not significantly change, underscoring that although these profile choices result in non-negligible differences in ocean biogeochemical distributions (Kriest et al., 2012; Aumont et al., 2017) and atmospheric CO<sub>2</sub> levels (Kwon et al., 2009), the differences are small enough that all the curves still quantitatively reproduce the observations to a similar degree.

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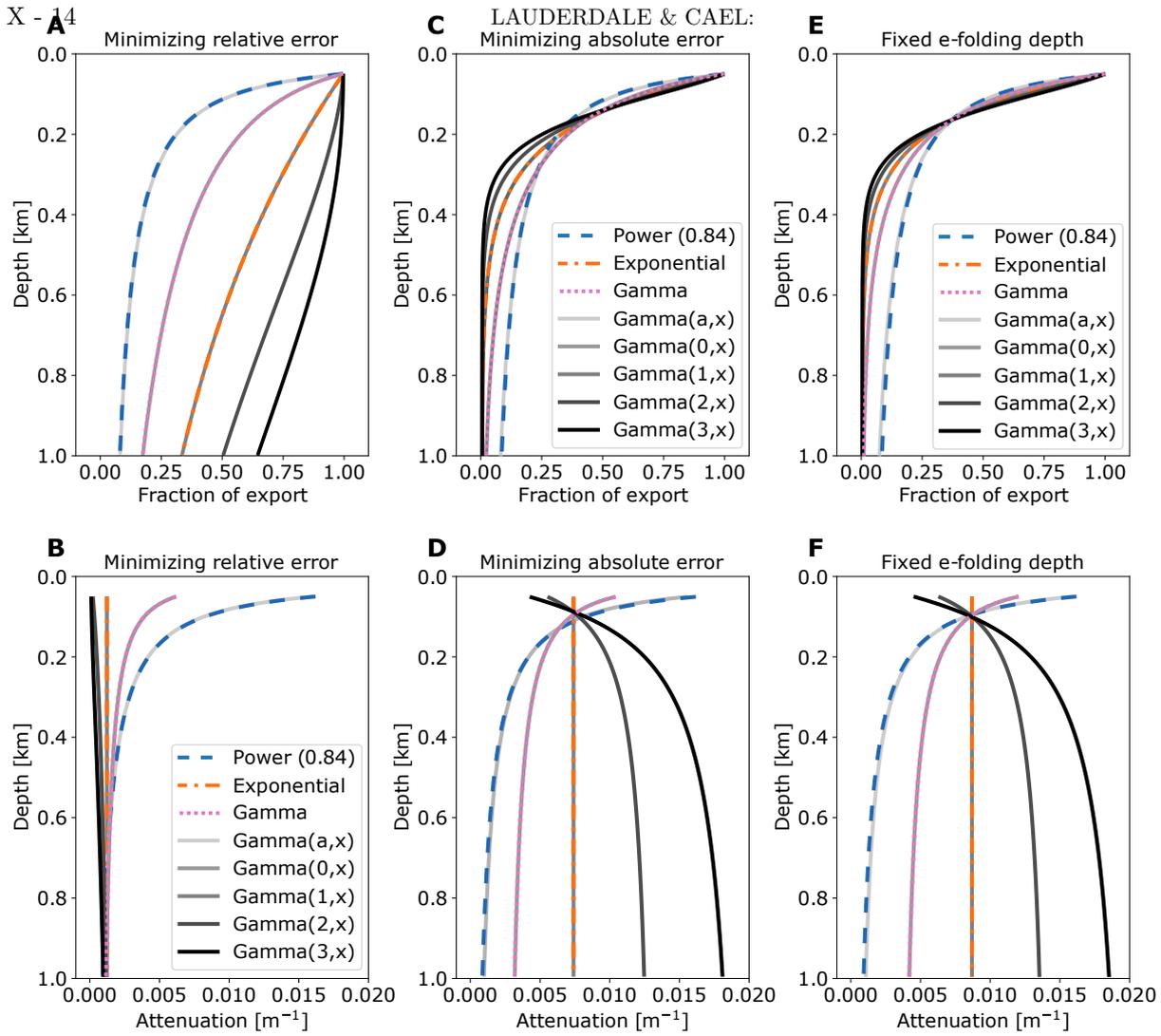
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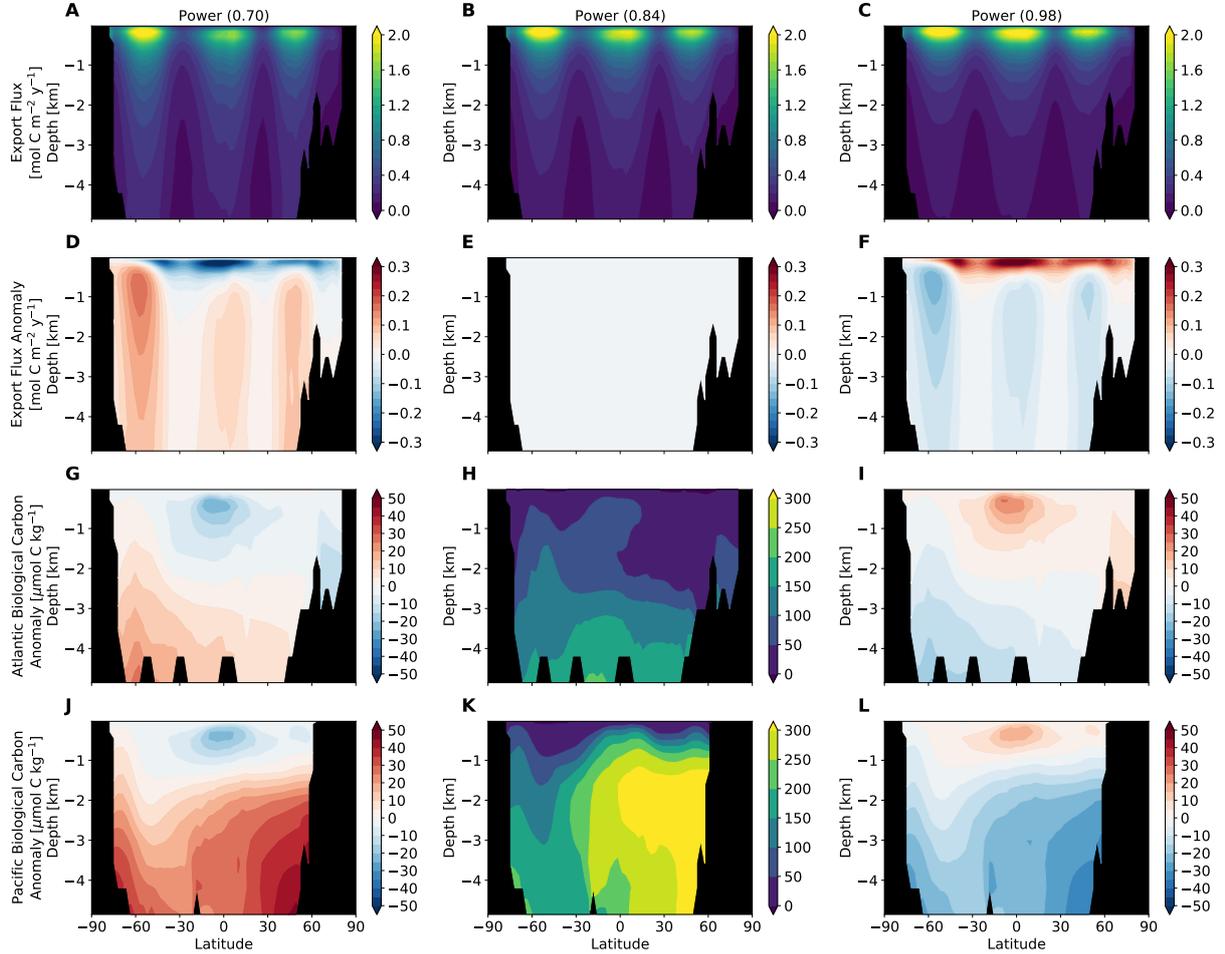
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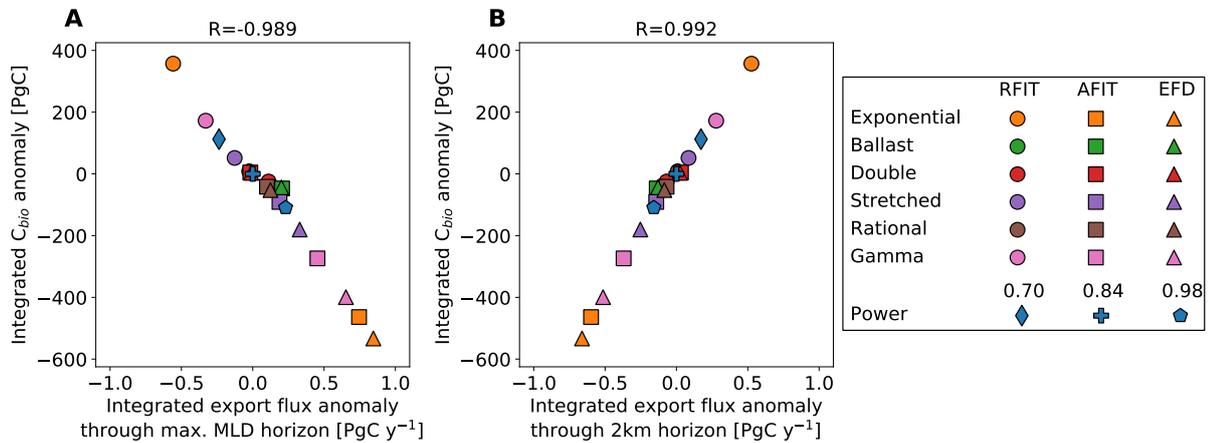
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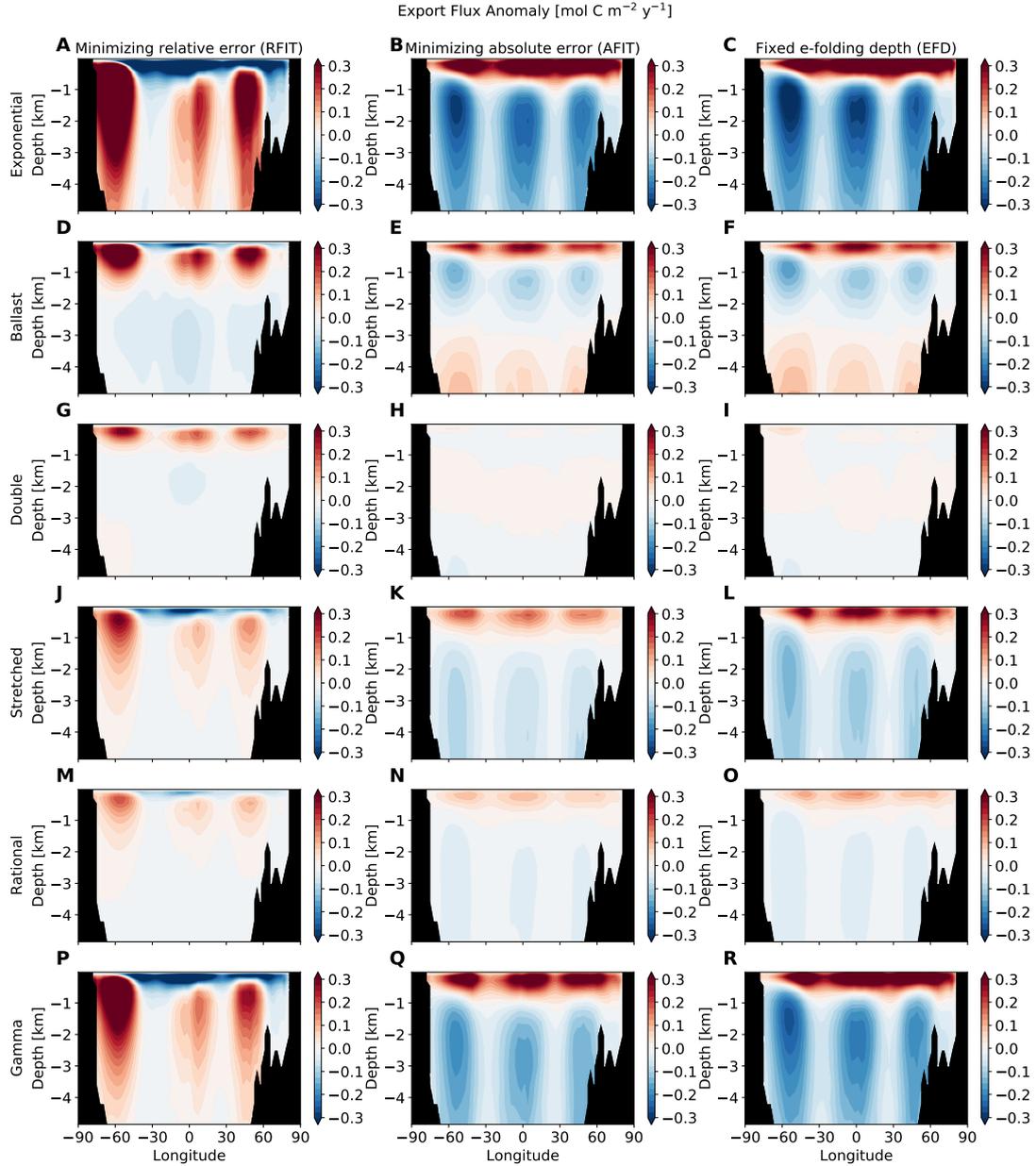
**Figure S1.** Fraction of sinking particulate organic matter exported from the 50 m surface layer remaining at each depth for different evaluations of the upper incomplete gamma function,  $\Gamma(a, x)$ , with varying values of the order  $a$  (increasing greyscale tones) and coefficients (e.g. Eqn. S6) statistically fit to the reference power-law curve ( $b=0.84$ ) by (a) statistically minimizing the relative error (“RFIT”), or (c) the absolute error (“AFIT”), and (d) matching the e-folding depth scale of 164 m (“EFD”). “Gamma( $a, x$ )” profiles are cases where the order is determined as part of a three-parameter non-linear fit giving  $a \approx -0.8$ . Profiles for the reference power-law ( $b=0.84$ ), simple exponential, and gamma function (Eqn. S6) curves from the main text (Fig. 2) are shown in colored, dashed, lines. Plots b, d, and f show the attenuation rate of the export flux with depth  $\left[\frac{1}{f} \frac{\partial f}{\partial z}, m^{-1}\right]$ .



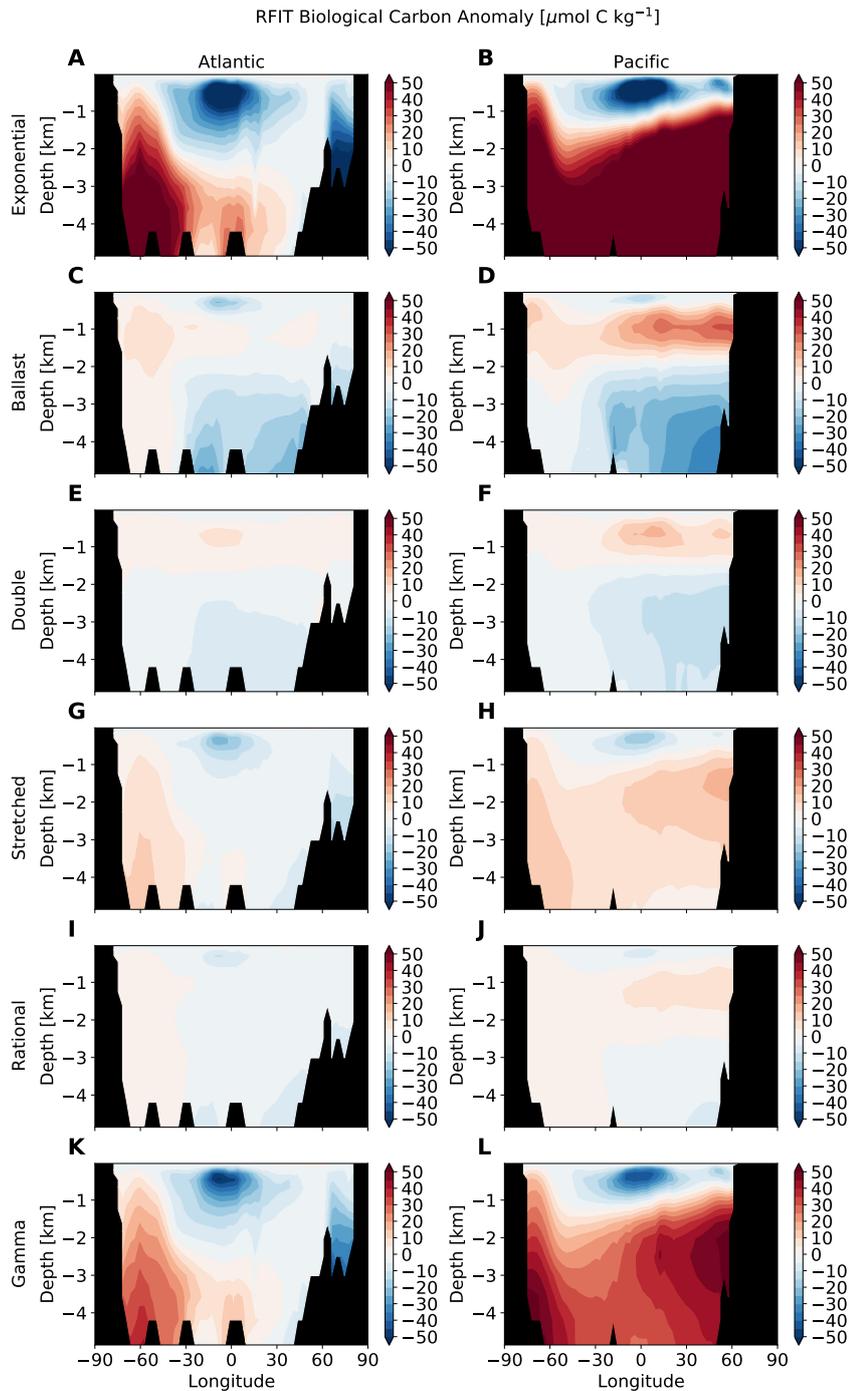
**Figure S2.** Zonal-average properties for power-law simulations where  $b = 0.84 \pm 0.14$  (a–c) particulate organic carbon export fluxes ( $\text{mol C m}^{-2} \text{y}^{-1}$ ), (d–f) anomalies of sinking particle export flux compared to the reference power-law simulations (i.e. the middle column,  $\text{mol C m}^{-2} \text{y}^{-1}$ ), and biological carbon concentration ( $C_{bio}$ ,  $\mu\text{mol C kg}^{-1}$  in the (g–i) Atlantic Ocean, and (j–l) Pacific Ocean. Reference  $C_{bio}$  concentrations are shown in the middle column, with anomalies for the decreased and increased values of the power-law exponent,  $b$ , in the left and right columns.



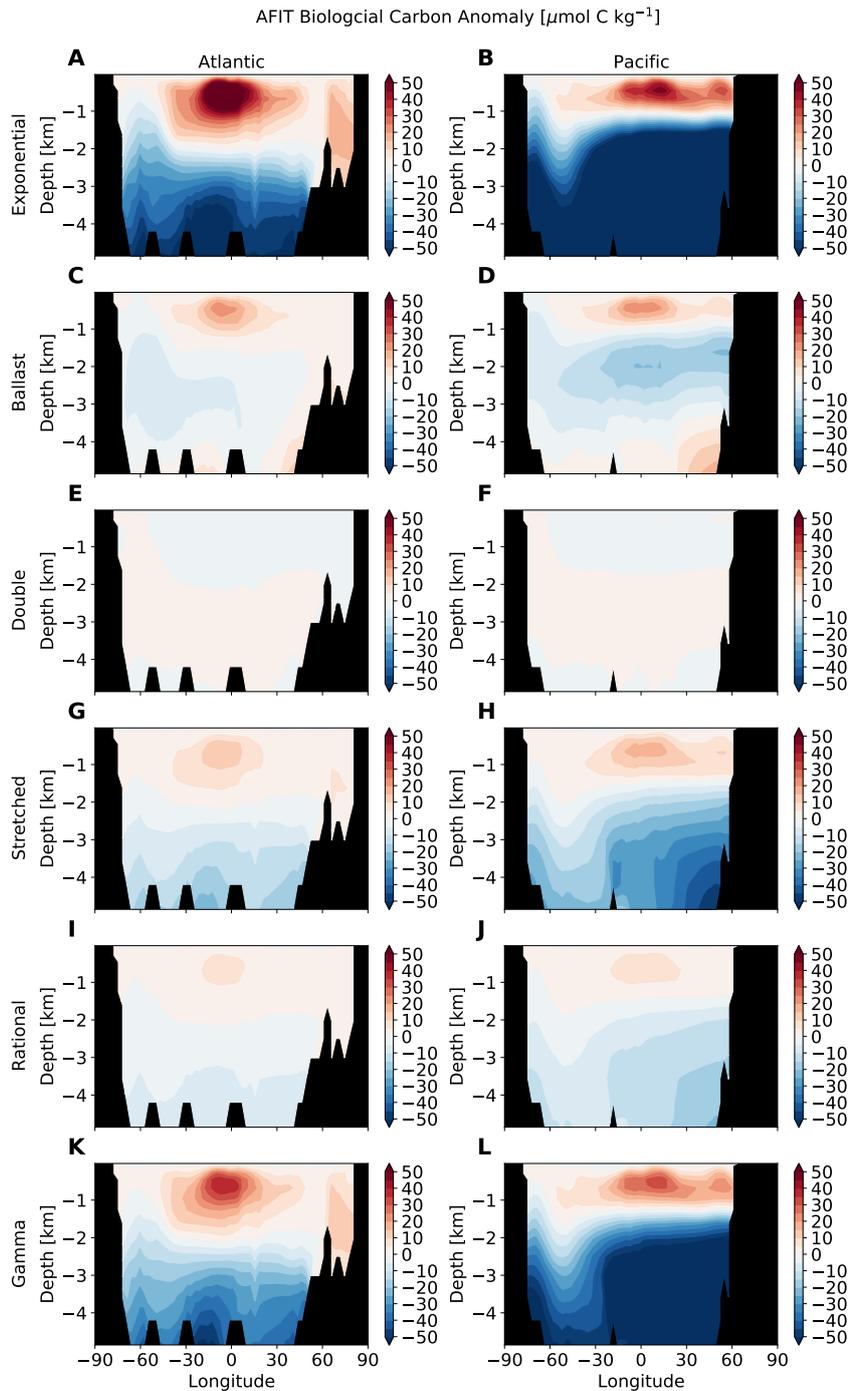
**Figure S3.** Change in the integrated particle export flux rate [ $\text{PgC y}^{-1}$ ] passing through (a) the horizon of deepest annual mixed layer depth, and (b) the 2 km depth horizon, against integrated biological carbon reservoir anomaly [ $\text{PgC}$ ], both with respect to the power-law curve where  $b=0.84$ . Three power-law simulations ( $b = 0.84 \pm 0.14$ ) are indicated by the blue symbols (diamond, cross, and pentagon), circle, square, and triangle symbols indicate that profile coefficients (Eq. S1–S6) were derived by minimizing the relative fit error (“RFIT”), minimizing the absolute fit error (“AFIT”), and fixing the e-folding depth of remineralization (“EFD”) to the reference power-law curve. Values are shown in Table S3.



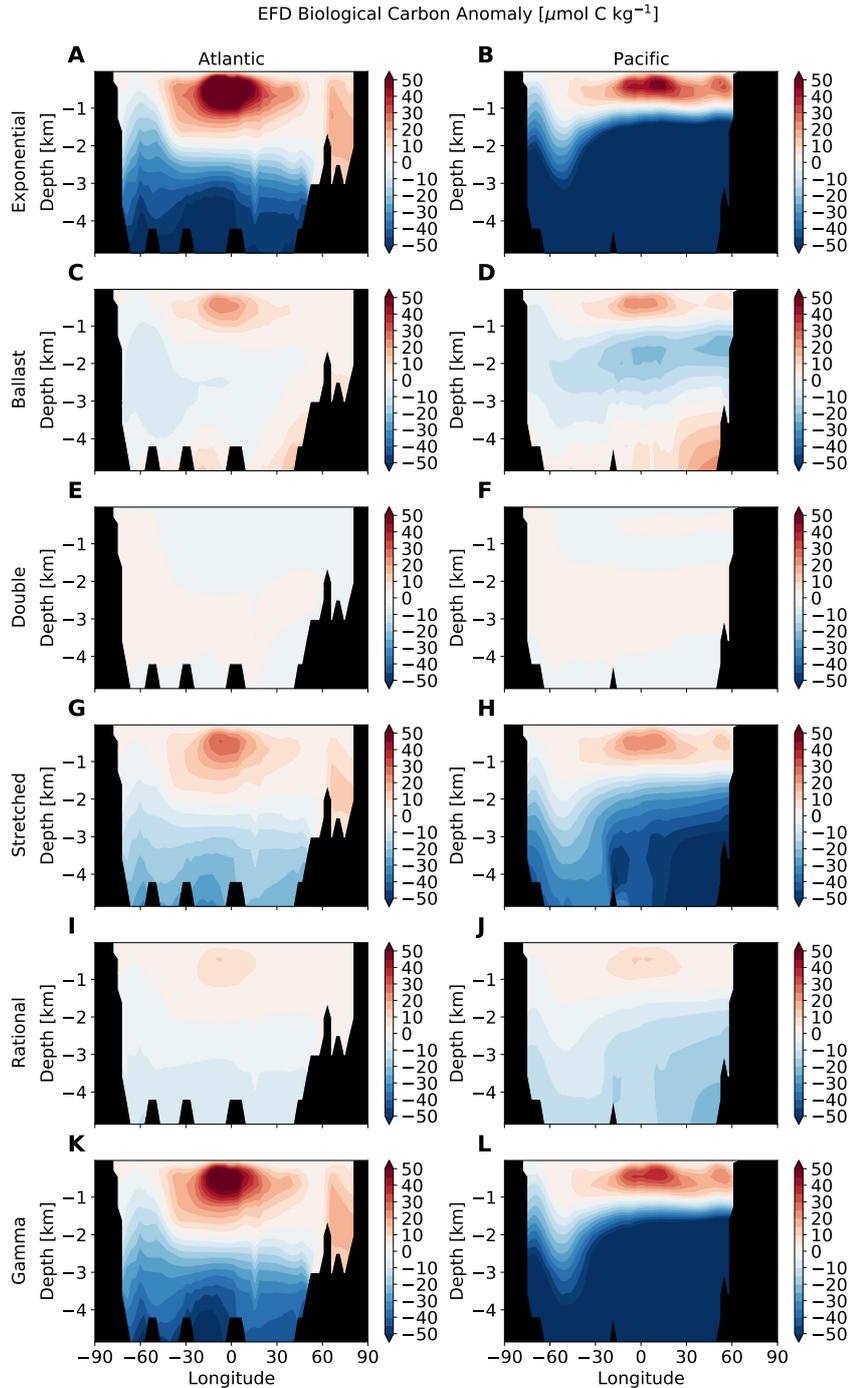
**Figure S4.** Zonally-averaged export flux anomaly with respect to the reference power-law curve where  $b=0.84$ , for the parameter sets where the relative error of the fit is minimized (RFIT, left column), where the absolute error of the fit is minimized (AFIT, middle column), and where the e-folding depth of remineralization is matched to the 164 m of the control curve (EFD, right column), where (a–c) is the simple exponential profile, (d–f) is the ballast profile, (g–i) is the double exponential profile, (j–l) is the stretched exponential profile, (m–o) is the rational profile, and (p–r) is the gamma profile.



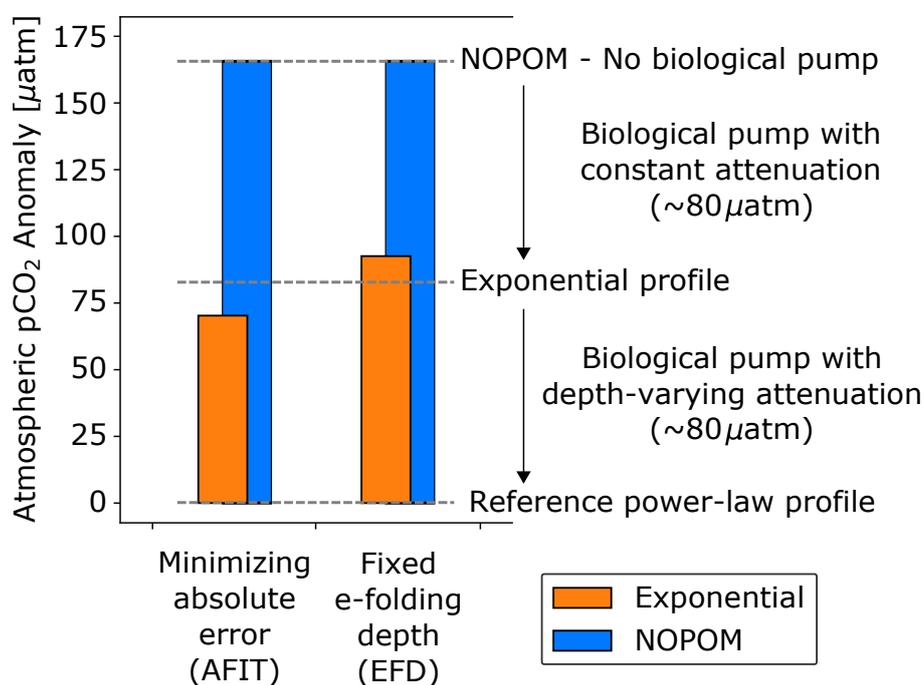
**Figure S5.** Zonally-averaged biological carbon ( $C_{bio}$ ,  $\mu\text{mol C kg}^{-1}$ ) anomalies with respect to the reference power-law curve for coefficients minimizing the relative error of the fit (RFIT) in the Atlantic Ocean (left column) and Pacific Ocean (right column) using (a–b) exponential, (c–d) ballast, (e–f) double exponential, (g–h) stretched exponential, (i–j) rational, and (k–l) gamma function remineralization profiles.



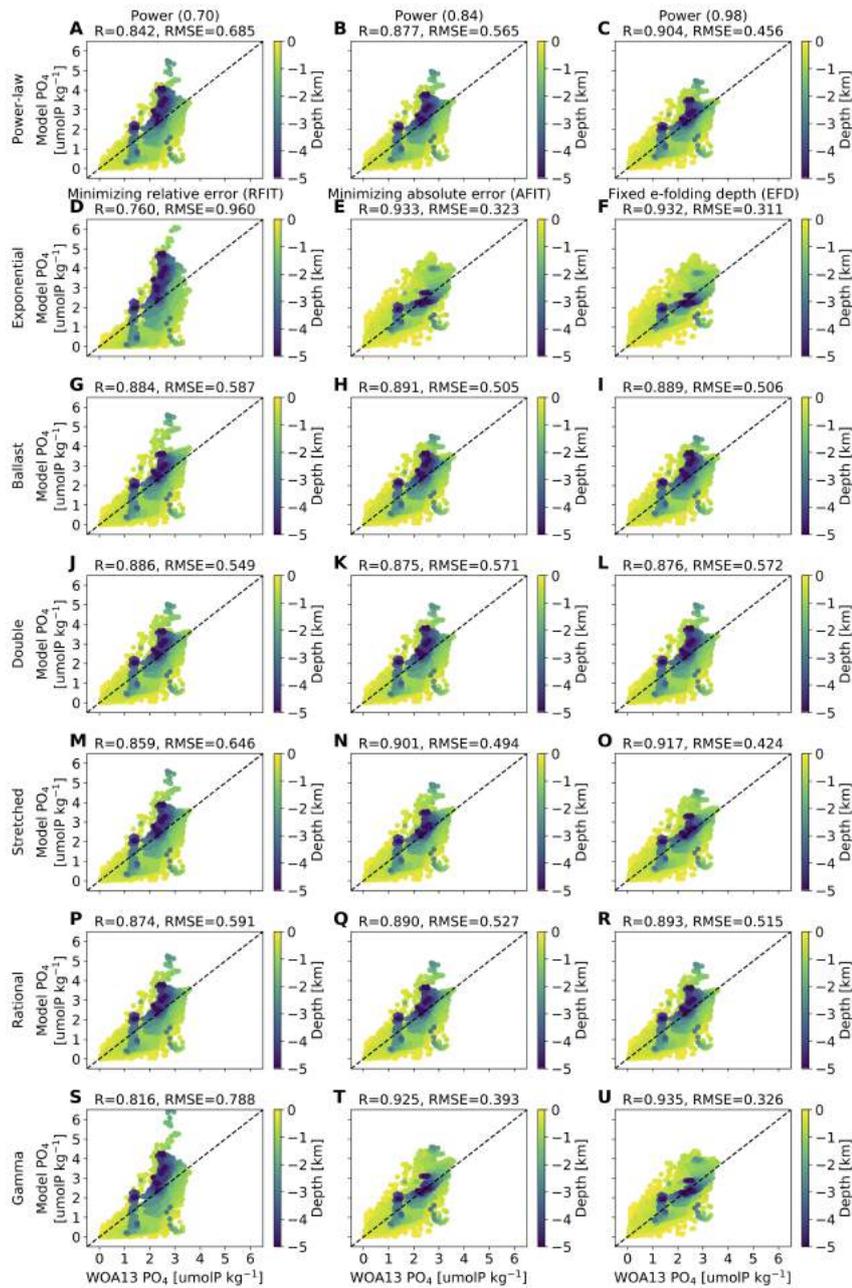
**Figure S6.** Zonally-averaged biological carbon ( $C_{bio}$ ,  $\mu\text{mol C kg}^{-1}$ ) anomalies with respect to the reference power-law curve for coefficients minimizing the absolute error of the fit (AFIT) in the Atlantic Ocean (left column) and Pacific Ocean (right column) using (a–b) exponential, (c–d) ballast, (e–f) double exponential, (g–h) stretched exponential, (i–j) rational, and (k–l) gamma function remineralization profiles.



**Figure S7.** Zonally-averaged biological carbon ( $C_{bio}$ ,  $\mu\text{mol C kg}^{-1}$ ) anomalies with respect to the reference power-law curve for coefficients matching the 164 m e-folding depth (EFD) in the Atlantic Ocean (left column) and Pacific Ocean (right column) using (a–b) exponential, (c–d) ballast, (e–f) double exponential, (g–h) stretched exponential, (i–j) rational, and (k–l) gamma function remineralization profiles.



**Figure S8.** Illustration of the importance of non-linearity in the biological pump remineralization profile. The simple exponential profile, with constant attenuation with depth, leads to significant outgassing of ocean carbon and increases in atmospheric pCO<sub>2</sub> of  $\sim 80 \mu\text{atm}$  for AFIT and EFD simulations, despite being statistically fit to the reference power-law. This is roughly half the  $\sim 165 \mu\text{atm}$  increase that results from removing the biological pump altogether (NOPOM), highlighting the importance for the air-sea carbon balance, not only of the existence of a biological pump that maintains interior ocean biological carbon stores, but also its non-linearity.



**Figure S9.** Model phosphate concentrations [ $\text{PO}_4$ ,  $\mu\text{mol kg}^{-1}$ ] for each remineralization profile against observations from the World Ocean Atlas 2013 (WOA13, Garcia et al., 2014) colored by model depth. The black dashed line indicates a 1:1 relationship. Model-data likeness is evaluated by the Pearson correlation coefficient,  $R$ , and the root-mean-squared concentration error,  $RMSE$ . Although remineralization profile choices result in non-negligible differences in macronutrient distributions, the differences are small enough that all the curves still quantitatively reproduce the observations to a similar degree.

**Table S1.** Parameter values and fit statistics for remineralization functions (Eq. S1–S6). Each function was matched to the reference power-law (Eq. 1) with exponent  $b=0.84$  by statistically minimizing the relative (“RFIT”) or absolute (“AFIT”) misfit of the curves, or by matching e-folding remineralization depth scale (“EFD”), Fig. 2b–d. Note different units of coefficients. Goodness of fit is evaluated by  $\mathcal{S}$ , the Standard Error of Regression (smaller numbers indicate better fit).

Shape	Parameter	Units	RFIT	AFIT	EFD
Exponential	$C_e$		1.059	1.451	1.548
	$\ell_e$	m	871.5	134.2	114.5
	$\mathcal{S}$		1.107	0.0701	0.0700
Ballast	$C_b$		1.200	1.487	1.530
	$\ell_b$	m	226.8	108.6	101.9
	$c$		0.03111	0.04159	0.04139
	$\mathcal{S}$		0.3838	0.0453	0.0440
Double Exponential	$C_{d1}$		1.326	1.583	1.522
	$\ell_{d1}$	m	124.3	70.38	75.09
	$C_{d2}$		0.08668	0.1466	0.1492
	$\ell_{d2}$	m	2521	1144	1170
Stretched Exponential	$\mathcal{S}$		0.1559	0.0175	0.0175
	$C_s$		10.88	13.91	15.81
	$s$		0.7776	0.7526	0.7404
Rational	$\mathcal{S}$		0.2499	0.0260	0.0314
	$C_r$	m	88.75	69.87	66.61
	$a$	m	38.75	19.87	16.61
Gamma Function	$\mathcal{S}$		0.1174	0.0112	0.0119
	$C_g$		0.3214	0.6003	0.7267
	$\ell_g$	m	1950	419.6	300.6
	$\mathcal{S}$		0.6272	0.0499	0.0543

**Table S2.** Supplementary quantities for power-law remineralization simulations with exponents of  $b=0.70$ ,  $0.84$ , and  $0.98$ , as well as the “NOPOM” simulation where there are no particulate organic matter export fluxes. Reference rates/concentrations are presented for the control power-law curve where  $b=0.84$ , while values presented for simulations where  $b=0.70$ ,  $0.98$ , and NOPOM are anomalies with respect to the control.  $\Delta B_C$  is the change in globally-integrated rate of net community production,  $\Delta E_{mld}$ ,  $\Delta E_{1km}$ , and  $\Delta E_{2km}$  are the change in areally-integrated particulate organic carbon export flux through the deepest mixed layer depth, 1 km, and 2 km horizons, respectively,  $\Delta C_{bio}$  is the globally-integrated change in biological carbon (evaluated as dissolved inorganic carbon minus the preformed carbon concentration), and  $\Delta pCO_2^{atm}$  is the change in atmospheric CO<sub>2</sub> partial pressure.

		Exponent ( $b$ )			
		0.70	0.84	0.98	NOPOM
$C_p$	$m^b$	1.000	1.000	1.000	
$\Delta B_C$	PgC $y^{-1}$	-5.231	29.570	5.175	39.65
$\Delta E_{mld}$	PgC $y^{-1}$	-0.236	2.349	0.230	-2.349
$\Delta E_{1km}$	PgC $y^{-1}$	0.141	1.749	-0.173	-1.749
$\Delta E_{2km}$	PgC $y^{-1}$	0.172	0.802	-0.159	-0.802
$\Delta C_{bio}$	PgC	112.32	2363.4	-109.30	-2187
$\Delta pCO_2^{atm}$	$\mu atm$	-21.59	269.33	24.77	165.4

**Table S3.** Supplementary anomalies for alternative remineralization profile simulations. Reference power-law values ( $b=0.84$ ) are given in Table S2.  $\Delta B_C$  is the change in globally-integrated net community production rate,  $\Delta E_{mld}$ ,  $\Delta E_{1km}$ , and  $\Delta E_{2km}$  are the change in areally-integrated particulate organic carbon export flux through the deepest mixed layer depth, 1 km, and 2 km horizons, respectively,  $\Delta C_{bio}$  is the globally-integrated change in biological carbon (evaluated as dissolved inorganic carbon minus the preformed carbon concentration), and  $\Delta pCO_2^{atm}$  is the change in atmospheric CO<sub>2</sub> partial pressure.

Shape	Parameter	Units	RFIT	AFIT	EFD
Exponential	$\Delta B_C$	PgC y <sup>-1</sup>	-17.39	7.866	11.01
	$\Delta E_{mld}$	PgC y <sup>-1</sup>	-0.558	0.745	0.845
	$\Delta E_{1km}$	PgC y <sup>-1</sup>	0.6860	-0.2890	-0.511
	$\Delta E_{2km}$	PgC y <sup>-1</sup>	0.5250	-0.5980	-0.662
	$\Delta C_{bio}$	PgC	357.1	-463.9	-533.3
	$\Delta pCO_2^{atm}$	$\mu atm$	-62.94	70.28	92.59
Ballast	$\Delta B_C$	PgC y <sup>-1</sup>	-5.218	2.549	3.108
	$\Delta E_{mld}$	PgC y <sup>-1</sup>	0.124	0.205	0.201
	$\Delta E_{1km}$	PgC y <sup>-1</sup>	0.3380	-0.244	-0.292
	$\Delta E_{2km}$	PgC y <sup>-1</sup>	-0.097	-0.138	-0.130
	$\Delta C_{bio}$	PgC	-38.40	-47.16	-43.58
	$\Delta pCO_2^{atm}$	$\mu atm$	-12.39	10.99	12.55
Double Exponential	$\Delta B_C$	PgC y <sup>-1</sup>	-1.099	-0.2730	-0.553
	$\Delta E_{mld}$	PgC y <sup>-1</sup>	0.111	-0.017	-0.011
	$\Delta E_{1km}$	PgC y <sup>-1</sup>	0.019	-0.021	-0.010
	$\Delta E_{2km}$	PgC y <sup>-1</sup>	-0.069	0.031	0.025
	$\Delta C_{bio}$	PgC	-24.66	4.514	3.255
	$\Delta pCO_2^{atm}$	$\mu atm$	-1.761	-1.082	-1.821
Stretched Exponential	$\Delta B_C$	PgC y <sup>-1</sup>	-5.272	1.226	4.483
	$\Delta E_{mld}$	PgC y <sup>-1</sup>	-0.127	0.188	0.329
	$\Delta E_{1km}$	PgC y <sup>-1</sup>	0.234	0.056	-0.075
	$\Delta E_{2km}$	PgC y <sup>-1</sup>	0.085	-0.142	-0.253
	$\Delta C_{bio}$	PgC	51.61	-90.93	-180.2
	$\Delta pCO_2^{atm}$	$\mu atm$	-17.92	10.31	28.60
Rational	$\Delta B_C$	PgC y <sup>-1</sup>	-2.525	0.892	1.593
	$\Delta E_{mld}$	PgC y <sup>-1</sup>	-0.025	0.100	0.124
	$\Delta E_{1km}$	PgC y <sup>-1</sup>	0.120	0.000	-0.026
	$\Delta E_{2km}$	PgC y <sup>-1</sup>	0.009	-0.069	-0.085
	$\Delta C_{bio}$	PgC	8.479	-41.63	-52.47
	$\Delta pCO_2^{atm}$	$\mu atm$	-7.745	5.612	8.583
Gamma Function	$\Delta B_C$	PgC y <sup>-1</sup>	-11.42	4.023	8.622
	$\Delta E_{mld}$	PgC y <sup>-1</sup>	-0.330	0.454	0.653
	$\Delta E_{1km}$	PgC y <sup>-1</sup>	0.471	-0.009	-0.282
	$\Delta E_{2km}$	PgC y <sup>-1</sup>	0.279	-0.370	-0.514
	$\Delta C_{bio}$	PgC	172.29	-273.74	-399.44
	$\Delta pCO_2^{atm}$	$\mu atm$	-40.38	35.70	66.35