

1 **First studies of mesosphere and lower thermosphere**
2 **dynamics using a multistatic specular meteor radar**
3 **network over southern Patagonia**

4 **J. Federico Conte¹, Jorge L. Chau¹, J. Miguel Urco¹, Ralph Latteck¹, Juha**
5 **Vierinen⁴, and Jacobo O. Salvador^{2,3}**

6 ¹Leibniz Institute of Atmospheric Physics at the University of Rostock, Kühlungsborn, Germany

7 ²Universidad Nacional de la Patagonia Austral, Ro Gallegos, Argentina

8 ³CONICET, Argentina

9 ⁴Arctic University of Norway, Tromsø, Norway

10 **Key Points:**

- 11 • First observations of MLT dynamics over the region of strongest stratospheric wave
12 activity are analyzed.
13 • Estimates of mean horizontal winds and their gradients are possible, thanks to the
14 multistatic configuration.
15 • Mean momentum fluxes are estimated with vertical velocity estimates free of hor-
16 izontal divergence contamination.

Abstract

This paper presents for the first time results on winds, tides, gradients of horizontal winds, and momentum fluxes at mesosphere and lower thermosphere (MLT) altitudes over southern Patagonia, one of the most dynamically active regions in the world. For this purpose, measurements provided by SIMONe Argentina are investigated. SIMONe Argentina is a novel multistatic specular meteor radar system that implements a SIMONe (Spread Spectrum Interferometric Multistatic meteor radar Observing Network) approach, and that has been operating since the end of September 2019. Average counts of more than 30000 meteor detections per day result in tidal estimates with statistical uncertainties of less than 1 m/s. Thanks to the multistatic configuration, horizontal and vertical gradients of the horizontal winds are obtained, as well as vertical winds free from horizontal divergence contamination. The daily averages of these preliminary results on the gradients are consistent with the expected MLT behavior during the summer and a strong southern hemisphere polar vortex. Mean momentum fluxes are estimated after removing the effects of mean winds using a four-hour, four-kilometer window in time and altitude, respectively. Reasonable statistical uncertainties of the momentum fluxes are obtained after applying a 28-day averaging. Therefore, the momentum flux estimates presented in this paper represent monthly mean values of waves with periods of four hours or less, vertical wavelengths shorter than four kilometers, and horizontal scales less than 400 km.

1 Introduction

The mesosphere and lower thermosphere (MLT) is the atmospheric region that couples the lower and upper parts of the terrestrial atmosphere. For this reason, knowledge of its dynamics is of great importance in order to understand the behavior of the atmosphere as a whole. The coupling is accomplished mainly via propagation of three dominant types of waves: planetary waves (PWs), tides and gravity waves (GWs). PWs are waves with scales of thousands of kilometers and periods of up to ~ 30 days. They are mainly generated in the troposphere by land-sea discontinuities, or triggered in-situ by, e.g., baroclinic instabilities and filtered gravity waves (e.g., Rossby, 1939; McCormack et al., 2014; H.-L. Liu & Roble, 2002). Tides are also waves with horizontal scales of thousands of kilometers, but periods that are sub-harmonics of the solar and lunar days. Thermal tides are mainly a consequence of solar radiation absorption by water vapour in the troposphere and ozone in the stratosphere, while the lunar tide results from the gravitational pull of the Moon (e.g., Lindzen & Chapman, 1969; Forbes, 1984). GWs are small to medium scale waves with periods ranging from a few minutes to several hours. They can be triggered by a myriad of different sources, e.g., the orography, thunderstorms, shear instabilities, convection, etc. (e.g., Hines, 1988; Piani et al., 2000; Fritts & Alexander, 2003).

During the last decades, specular meteor radars (SMRs) have been extensively used to study winds and atmospheric waves in the MLT (e.g., Hocking, 2005; Clemesha et al., 2009; Hoffmann et al., 2010; A. Z. Liu et al., 2013; Laskar et al., 2016; Jia et al., 2018, and references therein). They have also been used to study GWs, which are known to play an important role in determining the wind and thermal structure of the MLT (e.g., Fritts, 1984). Particularly, some studies have focused on extracting information about GW-driven momentum fluxes from SMR measurements (e.g., Fritts et al., 2010; Placke et al., 2011a; Andrioli et al., 2015). However, understanding the results on momentum flux estimates based on SMR winds is not trivial, mainly because of the uncertainties associated with the estimation procedure (e.g., Fritts et al., 2012b). In fact, Vincent et al. (2010) showed that the accuracy in the momentum flux estimation is highly dependent on the number of meteor detections. Consequently, the usage of multistatic meteor radar systems represents one way to reduce the uncertainties of the momentum flux es-

68 timates (e.g., Spargo et al., 2019). Furthermore, by detecting more meteors and being
 69 able to observe them from different viewing points, multistatic SMR systems also allow
 70 for more reliable estimations of horizontal wind gradients.

71 The MLT over the southern part of Argentina and Chile is considered to be one
 72 of the most dynamically active regions in the globe. Satellite-based studies have revealed
 73 that GW-driven momentum fluxes increase considerably at both stratosphere and MLT
 74 altitudes over Patagonia (e.g., Trinh et al., 2018; Vadas et al., 2019a). Numerical model
 75 simulations have reported generation of secondary GWs with horizontal scales of up to
 76 2000 km at mesospheric altitudes over the southern Andes (Vadas & Becker, 2019). Nev-
 77 ertheless, wave coupling processes in the MLT region over the Patagonian sector are still
 78 not well understood, partly because the installation of ground-based instruments has not
 79 been possible, either due to logistics challenges or instrument requirements. In this work,
 80 we present preliminary results of a multistatic SMR network that allows, for the first time,
 81 measurements of MLT dynamics in the Patagonian region. Besides the local support,
 82 our success has been possible thanks to a novel approach that we call SIMONE (Spread
 83 Spectrum Interferometric Multistatic meteor radar Observing Network) (Chau et al., 2019).
 84 SIMONE makes use of modern radar practices like spread-spectrum, MIMO (Multiple-
 85 Input, Multiple-Output), and compressed sensing applied to atmospheric radars (Vierinen
 86 et al., 2016; Urco et al., 2018, 2019). This allows for much easier installation, operation
 87 and expansion of the network than previous equivalent systems.

88 The paper is organized as follows. Section 2 introduces the SIMONE Argentina sys-
 89 tem. Section 3 provides a detailed description of the different analyses performed to the
 90 data. Then, we present and discuss the main results and findings in Section 4. Finally,
 91 in section 5 the concluding remarks are presented.

92 **2 SIMONE Argentina**

93 SIMONE Argentina is a state-of-the-art network of multistatic specular meteor radars
 94 that was installed in September of 2019 in the southern province of Santa Cruz, Argentina.
 95 It is comprised of one single transmitting site with five linearly polarized Yagi antennas
 96 in a pentagon configuration, and five receiving sites with one dual-polarization Yagi an-
 97 tenna each. The receivers are placed between 30 and 270 km of distance from the trans-
 98 mitting site, which is located at 49.6° S, 71.4° W (see Figure 1 for details on the geo-
 99 graphical distribution of the sites). This type of network configuration is known as MISO
 100 (Multiple-Input, Single-Output), since only one antenna is used on reception (e.g., Chau
 101 et al., 2019).

102 SIMONE Argentina is the result of an effort led by the Leibniz Institute of Atmo-
 103 spheric Physics (Germany) in collaboration with the Universidad Nacional de la Patag-
 104 onia Austral (Argentina), and the Arctic University of Norway. A similar system has been
 105 installed in Peru (SIMONE Peru). SIMONE systems use coded spread spectrum on trans-
 106 mission (Vierinen et al., 2016). A phase coded signal based on pseudo-random sequences
 107 is generated and transmitted on each antenna independently. Transmission is done at
 108 a frequency of 32.55 MHz and with an average power of 400 W per antenna. All five trans-
 109 mitted codes are simultaneously decoded at each receiving site by means of compressed
 110 sensing (e.g., Urco et al., 2019). Hardware and software details of both systems, i.e., SI-
 111 MONE Peru and SIMONE Argentina, can be found in Chau et al. (2020).

112 SIMONE Argentina started operations by the end of September 2019 and has been
 113 running since then with almost no interruptions. Figure 2 shows a summary of the de-
 114 tection statistics for the first seven months of operations. The upper panel indicates the
 115 normalized percentage of meteor counts for each individual link. The bottom panel is
 116 used to present the average daily total counts for each month. Problems with the local
 117 power supply at the transmitting site resulted in fewer meteor detections during April

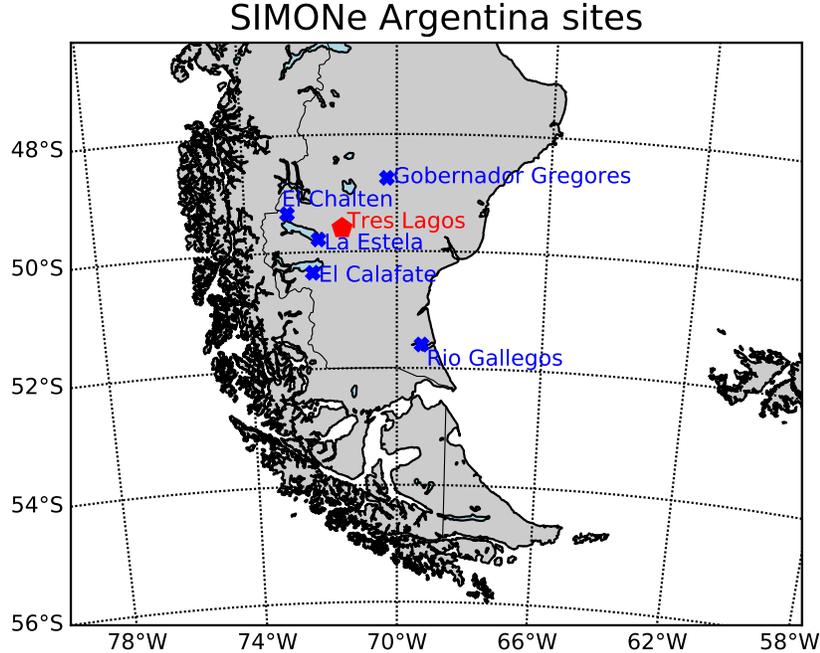


Figure 1. Map of SIMONe Argentina. The transmitter is indicated with the filled red pentagon, and the receivers are indicated with the blue crosses.

118 of 2020. Besides, the receiving site at Gobernador Gregores was out of operation dur-
 119 ing December of 2019 due to damage in the antenna cables. Nevertheless, for most of
 120 the time, the number of counts was much higher than in standard specular meteor radars.
 121 The links concentrating most of the meteor detections are Tres Lagos-El Calafate and
 122 Tres Lagos-La Estela. Starting in March 2020, the link Tres Lagos-Rio Gallegos exhibits
 123 a considerable increase in the counts, as a result of having rotated the transmitting an-
 124 tennas by 90 degrees. By month, January presents the largest counts, with an average
 125 of more than 50000 meteor detections per day.

126 3 Data analysis

127 Specular meteor radars (SMRs) are used to measure the Doppler shift of meteor
 128 trails due to their drifting with the mesospheric neutral winds (e.g., Jones et al., 1998).
 129 In order to extract the wind information from the measurements, one may implement
 130 an all-sky fit of the Doppler velocities measured during a certain period of time and within
 131 a given altitude interval (e.g., Hocking et al., 2001; Holdsworth et al., 2004). In other
 132 words, one must solve the following equation:

$$\mathbf{u} \cdot \mathbf{k} = 2\pi f + \zeta, \quad (1)$$

133 where $\mathbf{u} = (u, v, w)$ is the neutral wind vector, with u , v and w being its zonal (east-
 134 west), meridional (north-south) and vertical (up-down) components, respectively. $\mathbf{k} =$
 135 (k_u, k_v, k_w) is the Bragg wave vector (scattered minus incident) in the meteor-centered
 136 east-north-up coordinate system (perpendicular to the meteor trail); f is the Doppler
 137 shift; and ζ is the Doppler shift uncertainty. For this equation to be valid, one must make
 138 the assumption that the winds at each given height interval are uniform during the se-
 139 lected period of time (homogeneous method). The results using the homogeneous method
 140 have been obtained assuming $w = 0$.

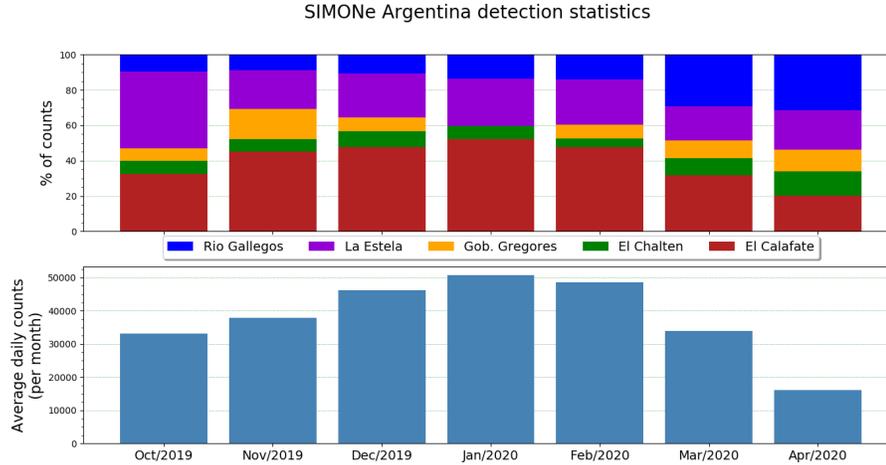


Figure 2. Upper panel: normalized percentage of meteor counts, color coded for each individual link. Bottom panel: monthly daily mean counts for all the links combined (see text for details).

141 Monostatic SMRs commonly allow for hourly horizontal wind estimations every 2-
 142 3 km in altitude (e.g., Jacobi et al., 1999; Hoffmann et al., 2010). The wind estimates
 143 are representative of mean values over an area of ~ 200 km in radius. These horizon-
 144 tal winds can be further processed in order to obtain information about large scale waves
 145 such as planetary waves and tides (e.g., Murphy et al., 2006; Chau et al., 2015; Conte
 146 et al., 2018).

147 One of the advantages of using multistatic SMR networks is that the amount of
 148 meteor detections is significantly increased (e.g., Stober & Chau, 2015). During most of
 149 the time since it started operations, SIMONE Argentina has been capable of detecting
 150 more than 30000 meteors per day (see Figure 2). With such amounts of meteor detec-
 151 tions, one can not only reliably estimate horizontal winds with higher time and verti-
 152 cal resolutions (~ 15 min and 1 km, respectively), but also investigate second order pa-
 153 rameters such as the squares of the perturbation components of the radial velocities, i.e.,
 154 the momentum flux components (e.g., Spargo et al., 2019). Another advantage multi-
 155 static SMR networks present is that they are capable of sampling the observed volume
 156 from different viewing angles. In this scenario, the increased number of meteor detec-
 157 tions can be further exploited in order to estimate first-order derivatives of the horizon-
 158 tal winds, and in this way include in the estimation a mean vertical wind free of hori-
 159 zontal divergence contamination (e.g., Chau et al., 2017).

160 Below, we describe the main procedures used in this work, i.e., tidal estimation,
 161 wind estimation using a gradient method, and mean momentum flux estimation. For all
 162 three procedures, a first wind estimation is carried out in order to remove outliers. This
 163 means solving Equation 1 in bins of the same size as that used later on in each given pro-
 164 cedure (e.g., 4 hour, 4 km in the gradient method); and then removing the radial veloci-
 165 ties the values of which have a corresponding residual of more than 25 m/s. This first
 166 wind estimation is carried out only in those bins containing a minimum of 10 meteor de-
 167 tectons.

168

3.1 Tidal estimation

169

170

171

172

173

174

175

Horizontal winds obtained from meteor radar measurements have been used for several decades to investigate tides and planetary waves in the MLT (e.g., Hocking & Thayer, 1997; Fritts et al., 2012a). Different mathematical techniques such as least squares or wavelet analysis can be applied to the wind data in order to extract the tidal information (e.g., Stening et al., 1997; Sandford et al., 2006; He et al., 2017). To avoid zero-padding or interpolating when encountering data gaps, a least squares approach was selected for this study.

176

177

178

179

180

181

182

183

After removing the outliers, Equation 1 was again fitted to the Doppler shift measurements. For this purpose, a weighted least squares (WLS) technique was implemented using bins of 1 hour and 2 km (in altitude), shifted by half an hour and 1 km, respectively. The inverse of the squared Doppler shift uncertainties (ζ in Eq. 1) were used as weights. The WLS was carried out only in those bins containing a minimum of 10 meteor detections. Then, under the assumption that the obtained hourly horizontal winds are the result of the superposition of a mean wind and different period oscillations, the following equation was fitted to the zonal (u) and meridional (v) wind components

$$[u + \psi_u, v + \psi_v] = [U_0, V_0] + \sum_{i=1}^4 [A_u, A_v]_i \cos\left(2\pi \frac{t - [\phi_u, \phi_v]_i}{T_i}\right). \quad (2)$$

184

185

186

187

188

189

190

191

Here, ψ_u and ψ_v are the Doppler shift uncertainties (error) propagated into the estimated winds; U_0 and V_0 are the mean zonal and meridional winds; A_{u_i} (A_{v_i}) and ϕ_{u_i} (ϕ_{v_i}) are the amplitude and phase, respectively, of the zonal (meridional) component of each considered wave; $T_1 = 2$ day; T_i for $i > 1$ is the period of each considered tide ($T_2 = 24$ h, $T_3 = 12$ h and $T_4 = 8$ h, for the diurnal, semidiurnal and terdiurnal solar tides, respectively); and t is the Universal Time (UT) in hours. The cosine of a sum was used to linearize Equation 2, which then was solved by applying the WLS method using a running window of 4 days shifted by 1 day.

192

3.2 Wind field gradient method

193

194

195

196

197

If one relaxes the assumption of homogeneity, the wind field inside the observed area may be estimated using the gradient method. This method consists in approximating the horizontal winds with their first-order Taylor expansion terms (e.g., Burnside et al., 1981; Browning & Wexler, 1968; Chau et al., 2017). This means introducing the following expression into Equation 1

$$\mathbf{u} = \mathbf{u}_0 + \mathbf{u}_x(x - x_0) + \mathbf{u}_y(y - y_0) + \mathbf{u}_z(z - z_0), \quad (3)$$

198

199

where $\mathbf{u}_0 = (u_0, v_0, w_0)$ represents the mean wind; $\mathbf{x}_0 = (x_0, y_0, z_0)$ is a reference point; and

$$\begin{aligned} \mathbf{u}_x &= \left(\frac{du}{dx}, \frac{dv}{dx}, \frac{dw}{dx} \right) \\ \mathbf{u}_y &= \left(\frac{du}{dy}, \frac{dv}{dy}, \frac{dw}{dy} \right) \\ \mathbf{u}_z &= \left(\frac{du}{dz}, \frac{dv}{dz}, \frac{dw}{dz} \right). \end{aligned}$$

200

201

202

203

204

205

206

207

The coordinates (x, y, z) are in km, and calculated taking into consideration the latitude, longitude and altitude of each meteor detection and the radius of the Earth at the reference point. \mathbf{x}_0 is determined using the latitude and longitude of the transmitting site, and the altitude of each height level considered in the WLS fit. The latter was implemented using bins of 4 hour and 4 km (in altitude), shifted by half an hour and 1 km, respectively. For this study, we have assumed that $(dw/dx, dw/dy, dw/dz) = \mathbf{0}$, which means solving for nine unknowns. In this method, a conservative condition of having a minimum of eighty one meteor detections per bin was selected.

3.3 Momentum flux estimates

The procedure followed in this study to estimate the GW momentum flux is based on the works by Thorsen et al. (1997) and Hocking (2005). It consists in applying a least square method to solve the following equation

$$(\mathbf{u}' \cdot \mathbf{k})^2 = (2\pi(f - \hat{f}))^2. \quad (4)$$

In this expression, \mathbf{k} and f are, respectively, the same Bragg wave vector and same Doppler shift as in Equation 1, but \mathbf{u}' represents the perturbed wind vector instead. \hat{f} is a so-called mean Doppler shift,

$$\hat{f} = \mathbf{u} \cdot \mathbf{k} / 2\pi, \quad (5)$$

where \mathbf{u} is the mean wind that results from solving Equation 1 or 3. To guarantee more reliable values of the six unknowns ($\langle u'u' \rangle$, $\langle v'v' \rangle$, $\langle u'w' \rangle$, $\langle v'w' \rangle$, $\langle u'v' \rangle$ and $\langle w'w' \rangle$), the fit was performed only in those bins containing forty meteor detections or more.

From Equation 5, it follows that different momentum flux estimates may be obtained depending on the mean wind that is used to determine \hat{f} . In this study, the wind estimates that result from the gradient method were used to calculate \hat{f} (see previous section). That is, Equation 3 (with $dw/dx = dw/dy = dw/dz = 0$) was introduced into Equation 1, and the latter then solved using bins of 4 hour and 4 km, shifted by 30 min in time and 1 km in altitude. Finally, Equation 4 was solved in bins of 1 hour and 2 km (in altitude), shifted by 30 min and 1 km, respectively.

4 Results and discussion

One of the main goals of this work is to provide, for the first time, information on the MLT dynamics over southern Patagonia obtained using SIMONe Argentina. For this reason, our results are discussed as they are presented.

4.1 Mean winds and tides

In Figure 3, we present the mean zonal (U_0) and meridional (V_0) winds, and the total amplitudes of the quasi two-day planetary wave (Q2DW) and the diurnal (D1), semidiurnal (S2) and terdiurnal (T3) solar tides. The vertical black dashed line indicates January 1st 2020. The term *total amplitude* refers to the magnitude of the vector sum of the corresponding zonal and meridional components of each fitted wave. The statistical uncertainties of the estimated parameters are shown in the right column panels. All quantities were obtained after applying the procedure detailed in Section 3.1. Data gaps are shown in white.

From inspection of Figure 3, two features stand out: the S2 tide is the dominant wave, with amplitudes in the order of 40-65 m/s, and the Q2DW exhibits strong enhancements after 4 January 2020. It is well known that the semidiurnal solar tide at middle latitudes dominates over all other tidal components (e.g., Andrews et al., 1987; Pancheva & Mukhtarov, 2011). Furthermore, many studies of tides in the northern hemisphere (NH) have reported that S2 decreases significantly around the onset of a sudden stratospheric warming (SSW) event, to later on recover and reach even larger amplitudes than those exhibited prior to the SSW (e.g., Chau et al., 2015; Siddiqui et al., 2018; Conte et al., 2019). In September of 2019, approximately 12 days before the 27th (first day of available data from SIMONe Argentina), a SSW event was registered in the southern hemisphere (e.g., Yamazaki et al., 2020). Interestingly enough, the largest amplitudes of S2 are seen between 27 September and 12 October, which might be an indicative of the recovery phase of S2 after the weakening associated with a SSW event. For the entire dataset analyzed in this study, S2 exhibits significant intraseasonal variability, which becomes evident in the many, although weaker, enhancements observed after \sim 31 October.

253 The Q2DW at middle latitudes has been reported to reach maximum amplitudes
 254 during the summer (e.g., Kumar et al., 2018). In our results, the Q2DW is active mostly
 255 in summer, in agreement with previous studies. Even more, it becomes the dominant wave
 256 by the end of January 2020, with amplitudes larger than those corresponding to S2. Offermann
 257 et al. (2011) showed that the Q2DW exhibits a triple peak structure in the NH during
 258 summer. Although it may not be obvious at first, after a more through inspection of Fig-
 259 ure 3, it can be noticed that the largest amplitudes of the Q2DW are distributed in three
 260 subsequent enhancements, around 7, 13 and 19 January. A fourth enhancement can be
 261 seen around 9 February, but the latter is significantly weaker than the previous three.

262 Both the diurnal and terdiurnal solar tides exhibit considerable intraseasonal vari-
 263 ability. In the case of D1, its activity becomes more evident mostly below ~ 90 km and
 264 during summer. Above ~ 92 km, and mainly during equinox times, T3 becomes more
 265 noticeable, with amplitudes similar to those corresponding to D1.

266 Compared to five-year average values at 54° S presented by Conte et al. (2017), the
 267 summer reversal of the mean zonal wind shown in Figure 3 is observed at altitudes ap-
 268 proximately two km lower. This is consistent with previous studies reporting a decrease
 269 with latitude of the height of the zonal wind summer reversal (e.g., Hoffmann et al., 2010;
 270 Wilhelm et al., 2019). Besides, it might seem that U_0 starts the transition into summer
 271 conditions relatively early, around 3 October. However, above ~ 92 km, U_0 experiences
 272 a late reversal to westward conditions around 24 October, to finally go into summer con-
 273 ditions (i.e., eastward above the mesopause) about five days later. V_0 blows mainly to-
 274 wards the equator, and only after 10 March poleward values start to dominate.

275 The statistical uncertainties of all the fitted parameters presented in Figure 3 are
 276 very small. Only above 103 km and below 77 km, values of ~ 2 -3 m/s are obtained (not
 277 shown). The low statistical uncertainties are a consequence of the large amount of me-
 278 teor detections provided by SIMONe Argentina. By solving Equation 1 in bins of one
 279 hour and two km, one guarantees wind estimates with very low uncertainties. The lat-
 280 ter, combined with the fact the WLS method used to solve Equation 2 is applied to a
 281 very well conditioned matrix, results in small statistical uncertainties.

282 4.2 Gradients and vertical wind

283 In Figure 4, we present the winds and gradients obtained after applying the pro-
 284 cedure described in Section 3.2. Given that we want to minimize the effects that tides
 285 may have on the gradients, only daily averages are presented. The first four panels on
 286 the left show the mean zonal wind and the mean zonal eastward, northward and upward
 287 derivatives, i.e., u_0 , u_x , u_y and u_z , respectively. The corresponding four panels on the
 288 right show the mean meridional wind and mean meridional first-order derivatives, i.e.,
 289 v_0 , v_x , v_y and v_z , respectively. The fifth panel is reserved to show the mean vertical wind
 290 (w_0) on the left, and its statistical uncertainty ($\sigma(w_0)$) on the right. The gradients are
 291 in m/s/km, while the remaining parameters are in m/s. Data gaps are shown in white.

292 Before starting the description of the main features observed in Figure 4, it is im-
 293 portant to stress here that the variability seen in the mean winds is representative of large
 294 scale structures, with periods greater than four hours and vertical wavelengths larger than
 295 four km. u_0 and v_0 do not differ much from the mean winds shown in Figure 3. They
 296 do show more variability, although this is expected since u_0 and v_0 are daily averages
 297 of 4-hour winds, while the u and v presented in the previous section are 4-day mean winds.
 298 Part of this variability may be attributed to the Q2DW, which becomes noticeable around
 299 5 January in u_0 , and around 25 January in the case of v_0 .

300 Some previous observational studies have investigated the horizontal gradients in
 301 the zonal and meridional winds (e.g., Conde & Smith, 1998; Meriwether et al., 2008; Chau
 302 et al., 2017). However, the present study is the first one to show results on both the hor-

SIMONe Argentina - Mean winds & tides

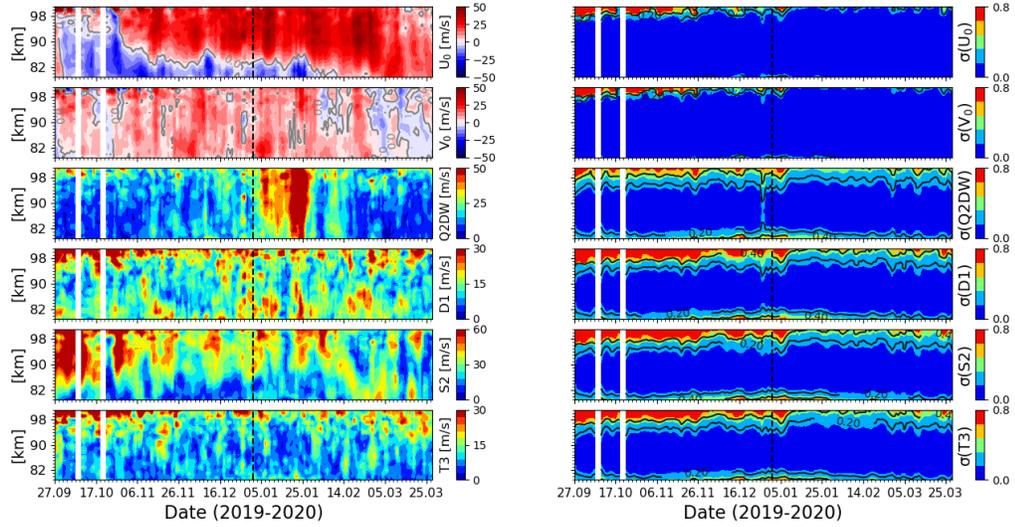


Figure 3. Mean zonal (U_0) and meridional (V_0) winds, the quasi two-day planetary wave (Q2DW), and the diurnal (D1), semidiurnal (S2) and terdiurnal (T3) solar tides during September 2019 - April 2020 over SIMONe Argentina. The corresponding statistical uncertainties (in m/s) are shown in the right column. All parameters were estimated using a four-day running window shifted by one day. The vertical black dashed line indicates January 1st 2020.

303 zontal and vertical gradients of MLT horizontal winds over southern Patagonia. From
 304 Figure 4, it can be seen that u_x and v_y have amplitudes that, when added together, may
 305 result in values of u_x+v_y in the order of ~ 0.4 m/s/km. Chau et al. (2017) showed that
 306 horizontal divergence values of ~ 0.1 m/s/km are large enough to introduce an appar-
 307 ent vertical wind of 1-2 m/s. Consequently, one can now understand the importance of
 308 estimating the vertical wind together with the gradients. By doing the latter, one reduces
 309 significantly the effects of biases introduced in w_0 by horizontal variability of u and v .
 310 On the other hand, u_x is expected to be zero (on average) during the summer. After com-
 311 puting the mean over longer intervals of time (e.g., 10 days) and hence averaging out long
 312 period waves typical of summer, such as the Q2DW, u_x becomes indeed very small, with
 313 values of ~ 0.01 - 0.03 m/s/km and mainly negative above ~ 82 km, and positive below
 314 (not shown here).

315 In the case of u_y and v_x , daily amplitude values are in the order of 0.3 m/s/km.
 316 The larger amplitudes of u_y (compared to u_x) may be an indicative of zonal wind lat-
 317 itudinal changes related to the SH polar vortex, part of which usually locates below the
 318 region seen by SIMONe Argentina (e.g., Figure 8 in Orte et al. (2019)). For altitudes
 319 lower than ~ 84 km, u_y (v_x) is mainly negative (positive). The upward component of
 320 the relative vorticity can be coarsely approximated by v_x-u_y (for a precise calculation,
 321 one needs to include the latitude information; see Equation A16 in Chau et al. (2017)).
 322 Notice that below ~ 84 km, v_x-u_y is mainly positive, which suggests an upward move-
 323 ment due to vortical effects. Besides, differences between the amplitudes of u_y , v_x and
 324 u_x , v_y , suggest that changes in the horizontal gradients due to GWs depend on the prop-
 325 agation direction of the these waves.

326 The vertical gradients of both the zonal and the meridional wind components, i.e.,
 327 u_z and v_z , show the largest amplitudes of all the first-order derivatives analyzed in this

work. Between 81 and 90 km of altitude, and during approximately the same period of time the zonal wind summer reversal is observed, u_z shows strong positive amplitudes (~ 15 m/s/km or more). During the summer, the zonal wind decreases with the altitude. This negative vertical gradient of the zonal wind allows eastward propagating GWs to easily reach mesospheric altitudes, where they break and deposit momentum and energy. This deposition of momentum creates an eastward drag that decelerates the zonal (westward) wind. Due to the Coriolis effect, the deceleration of the zonal wind introduces an equatorward meridional wind component, which in turn leads to an upward motion and a subsequent (adiabatic) cooling of the mesopause region (e.g., Smith, 2012). Finally, due to the thermal wind equation, $u_z \approx -\partial T/\partial y$ (where y is the latitude, and T is the temperature), the adiabatic cooling results in the positive values of u_z seen in Figure 4. In the case of v_z , the amplitudes are in the order of ~ 10 m/s/km. Besides, there is considerable day-to-day variability, although no clear pattern or structure in time can be noticed.

The mean vertical wind (w_0) exhibits large amplitudes, with values in the order of ~ 6 -11 m/s. This study is not the first one to report large amplitudes in the mean vertical winds estimated using SMRs. For example, Babu et al. (2012) and Egito et al. (2016) reported that the vertical winds at low latitudes may reach magnitudes of 6-10 m/s. In our study, the maximum amplitudes are observed mostly above 90 km of altitude. Notice that below 90 km, w_0 is predominantly positive during the summer, and only becomes predominantly negative after 25 January, when eastward values of u_0 start to dominate at all height levels above 81 km. It is important to stress again that, by estimating w_0 together with the first-order derivatives, possible contamination by the horizontal gradients is removed. In Figure 4, we only present the statistical uncertainties of w_0 , because they happen to be the largest ones (for the other estimates, the largest values of the *relative* σ are in the order of 5-7 %). Note that only for altitudes above 96 km and below 79 km, where the number of meteor detections is lower, $\sigma(w_0)$ exceeds values of 1 m/s.

It is of interest mentioning that mostly above 90 km of altitude, and during certain days of the summer, the 4-hour w_0 estimates (without computing the daily mean) exhibit a clear diurnal pattern (not shown here). Similar diurnal oscillations in w_0 , and also in the horizontal gradients are reported in Chau et al. (2020). Exploring diurnal and other period oscillations that may be present in the vertical winds will be the focus of a future study.

4.3 Gravity-wave-driven momentum flux

We now present and discuss the momentum flux estimates obtained after subtracting the mean winds calculated following the gradient method, i.e., the u_0 , v_0 and w_0 that were estimated together with the gradients. In Figure 4, it can be seen that the horizontal wind gradients and the vertical mean wind exhibit considerable variability, both in time and altitude. Part of this variability is leaked into the mean horizontal winds when one solves Equation 1 without including the gradients and w_0 in the wind vector \mathbf{u} . This implies that the u_0 , v_0 , w_0 obtained using the gradient method constitute a better representation of the real mean wind, provided that enough meteors are detected in order to accomplish a robust wind estimation. Besides, it is important to have in mind that the subtracted winds were estimated in bins of 4 hour and 4 km (in altitude). This means that the corresponding momentum flux estimates are representative mostly of waves with temporal and vertical scales of less than 4 hour and 4 km, respectively.

In Figure 5, we present 28-day averages of the momentum flux estimates that result from subtracting the \hat{f} that was calculated using the u_0 , v_0 , w_0 discussed in the previous paragraph (see Equation 4). The averages were calculated over 28 days in order to obtain estimates that are statistically more significant. The upper panels present U

Daily mean of 4-h winds and gradients

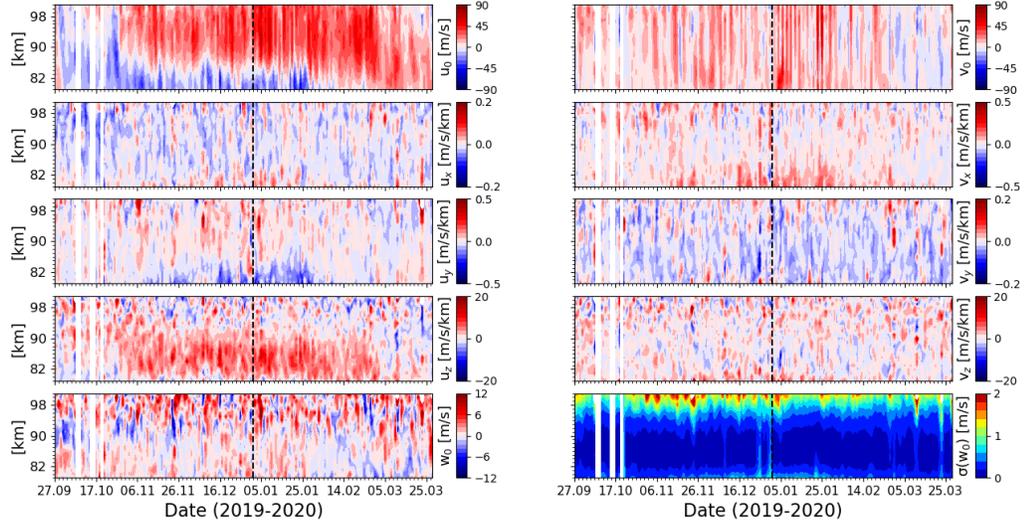


Figure 4. Daily averages of: (first panel) mean zonal (u_0) and meridional (v_0) winds, (second panel) eastward zonal (u_x) and meridional (v_x) derivatives, (third panel) northward zonal (u_y) and meridional (v_y) derivatives, and (fourth panel) upward zonal (u_z) and meridional (v_z) derivatives. The mean vertical wind (w_0) and its statistical uncertainty ($\sigma(w_0)$) are also included (bottom panel). All parameters were estimated using a 4 hour, 4 km (in vertical) bin, shifted by 30 min and 1 km, respectively. The vertical black dashed line indicates 1 January 2020.

379 and V , which correspond to the 28-day averages of the mean zonal and meridional winds,
 380 respectively. The middle panels are used to show the 28-day mean horizontal momen-
 381 tum vertical fluxes, $\langle u'w' \rangle$, $\langle v'w' \rangle$. The statistical uncertainties of $\langle u'w' \rangle$ and $\langle v'w' \rangle$,
 382 i.e., $\sigma(u'w')$ and $\sigma(v'w')$, respectively, are shown in the bottom panels.

383 Momentum flux estimations based on SMR observations were first presented by Hocking
 384 (2005). Since then, several studies have investigated momentum fluxes using meteor radar
 385 winds (e.g., Fritts et al., 2010; Placke et al., 2011b; Fritts et al., 2012b; Andrioli et al.,
 386 2013; Placke et al., 2015, and references therein). More recently, de Wit et al. (2016) ob-
 387 served a modulation by the quasi-biennial oscillation (QBO) of the momentum fluxes
 388 over Tierra del Fuego, an island south of Santa Cruz province. Typical amplitudes of $\langle u'w' \rangle$
 389 and $\langle v'w' \rangle$ reported in these previous studies rarely reach values larger than $60\text{--}70 \text{ m}^2/\text{s}^2$.
 390 Using a multistatic meteor radar network over southern Australia, Spargo et al. (2019)
 391 observed values for $\langle u'w' \rangle$ and $\langle v'w' \rangle$ in the order of $40\text{--}50 \text{ m}^2/\text{s}^2$, which are larger
 392 than momentum flux absolute values obtained from satellite measurements (e.g., Ern et
 393 al., 2011; Trinh et al., 2018). The latter is most likely due to observational features in-
 394 herent to satellites.

395 From Figure 5, it can be seen that our momentum flux estimates have amplitudes
 396 in the order of $40\text{--}60 \text{ m}^2/\text{s}^2$. Some large amplitudes of more than $100 \text{ m}^2/\text{s}^2$ can be ob-
 397 served above 97 km of altitude, where the corresponding statistical uncertainties are in
 398 the order of $50 \text{ m}^2/\text{s}^2$ (see bottom panels). Both momentum flux estimates exhibit vari-
 399 ability in time and height, although it is more evident in the case of $\langle u'w' \rangle$. Notice that
 400 despite the 28-day averaging, the magnitudes of $\langle u'w' \rangle$ and $\langle v'w' \rangle$ are considerable
 401 throughout the entire period of time analyzed in this study. Momentum flux estimates
 402 averaged over 20 days presented by Placke et al. (2015) exhibit maximum amplitudes

of 6-8 m²/s² during the summer of the northern hemisphere. The largest amplitudes of ten-day momentum flux estimates shown by Spargo et al. (2019) reach values of 30-40 m²/s² in the beginning of the spring of the southern hemisphere. Compared to these studies, it is clear that the wave activity in the MLT region over the Patagonian sector is very strong. The latter is consistent with previous studies based on satellite measurements (e.g., Ern et al., 2011), and numerical simulations (e.g., Lund et al., 2020).

As GWs propagate upwards, they transfer momentum and energy into the mean flow. Consequently, a decrease in the vertical flux of zonal momentum should correspond with an increase in the zonal wind speed (e.g., Fan et al., 1991). In other words, the mean zonal wind and $\langle u'w' \rangle$ should have opposite signs, i.e., $\langle u'w' \rangle$ positive in summer and negative during the winter. Our results are consistent with this reasoning, since for altitudes below 90-92 km, $\langle u'w' \rangle$ is mainly positive until ~ 24 February. Furthermore, these positive $\langle u'w' \rangle$ amplitudes exhibit a vertical gradient: as the altitude increases, they progressively decrease from values of around 90 m²/s² below 82 km, to values of ~ 10 m²/s² above 90 km. The latter is also consistent with the results on u_z presented in the previous section. Besides, note that after 22 February, negative values of $\langle u'w' \rangle$ start to develop below 90-91 km. At that time of the year, U has become eastward at all altitudes observed by SIMONE, a condition that allows westward propagating GWs to reach higher altitudes and, most likely, induce the aforementioned $\langle u'w' \rangle$ negative values.

In the case of $\langle v'w' \rangle$, an upward movement of southward momentum dominates mostly below ~ 91 km, from the beginning of October until the end of March. During the last week of December 2019, positive values of $\langle v'w' \rangle$ start to dominate above 92-94 km, an altitude range that had been dominated by negative values of $\langle v'w' \rangle$ since the beginning of November 2019. During the latter, and approximately above 96 km of altitude, an upward movement of eastward momentum can be noticed again. These positive values of $\langle v'w' \rangle$ develop very abruptly around 26 November, and remain dominant above 95 km for more than 25 days. We wonder if this might be an indicative of eastward momentum deposition by GWs that were in-situ generated at altitudes above 90 km.

To finalize, we discuss the procedure followed to estimate $\sigma(u'w')$ and $\sigma(v'w')$, and the reason for calculating 28-day averages. Unless one knows with some degree of certainty that a given wave event has occurred, the effects of GWs should be treated as stochastic processes. In other words, the mean momentum flux estimates are highly dependent on the effects of the geophysical variability. Kudeki and Franke (1998) showed that in order to obtain statistically significant momentum flux estimates at mesospheric heights, one must consider averaging intervals of more than 25 days. Specifically, they found that the statistical uncertainty of $\langle u'w' \rangle$ can be approximated with:

$$\sigma(u'w') = \sqrt{\frac{\langle u'u' \rangle \langle w'w' \rangle}{T/\tau}}, \quad (6)$$

where, $\langle u'u' \rangle$ and $\langle w'w' \rangle$ are averaged over the interval of time T ; and τ is equal to half of the mesosphere Brunt-Väisälä period (~ 7 min). $\sigma(v'w')$ is obtained using same Equation 6, but replacing $\langle u'u' \rangle$ by $\langle v'v' \rangle$. Selection of an averaging window $T = 28$ days resulted in the values presented in Figure 5. In this way, the $\langle u'w' \rangle$ and $\langle v'w' \rangle$ obtained from our study must be understood as representatives of a monthly mean momentum flux due to waves with periods of 4 h or less, and horizontal scales less than 400 km. Besides, for those estimates corresponding to altitudes lower than ~ 98 km, statistical uncertainties between 2-3 and 35 m²/s² should be taken into consideration.

The momentum flux estimates are also affected by the correlated Doppler shift errors. In other words, because the Doppler shift uncertainties are squared when introduced into Equation 4, the resulting momentum flux estimates are in fact an overestimation

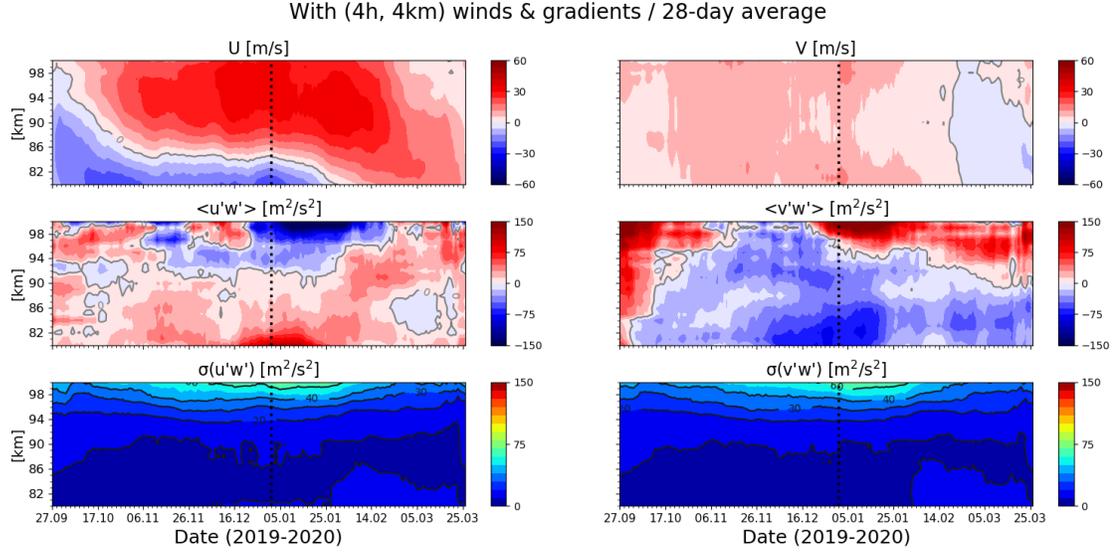


Figure 5. 28-day averages of 4-h, 4-km mean zonal and meridional winds (U and V , respectively), horizontal momentum vertical fluxes ($\langle u'w' \rangle$, $\langle v'w' \rangle$) and their corresponding statistical uncertainties ($\sigma(u'w')$, $\sigma(v'w')$). 4 hour, 4 km (in altitude) mean horizontal and vertical winds estimated in combination with the gradients (i.e., the u_0 , v_0 , w_0 described in Section 3.2) were subtracted before estimating the momentum flux. The vertical black dashed line indicates 1 January 2020.

452 of the real $\langle u'w' \rangle$ and $\langle v'w' \rangle$ (Vierinen et al., 2019). However, given that the Doppler
 453 shift uncertainties we obtained are small and that the amplitudes of $\langle u'w' \rangle$ and $\langle v'w' \rangle$
 454 are large, the effects of the correlated errors were neglected in this study. We plan to fur-
 455 ther investigate this issue by extending our analysis to include non-zero lag second-order
 456 statistics of the wind velocities.

457 5 Concluding remarks

458 The first multistatic meteor radar based studies of mean winds, tides, gradients and
 459 momentum flux over the southern part of Patagonia have been presented in this paper.
 460 By doing this, we have demonstrated the ability of SIMONe Argentina to obtain not only
 461 information on typical MLT parameters such as mean winds and tides, but also to suc-
 462 cessfully estimate previously little investigated parameters, such as horizontal and ver-
 463 tical gradients of the horizontal winds. Using the latter, one can estimate, e.g., the hor-
 464 izontal divergence and the relative vorticity, parameters from which global circulation
 465 models can benefit and in this way help to further the understanding of MLT dynam-
 466 ics.

467 Our results show a strong positive vertical gradient in the zonal wind during the
 468 summer, in agreement with the residual mean meridional circulation. Besides, the hor-
 469 izontal gradients of the zonal and meridional winds seem to indicate possible effects of
 470 the southern hemisphere polar vortex. The mean vertical wind (w_0) has also been es-
 471 timated, but only when the horizontal and vertical gradients were taken into account.
 472 The amplitudes obtained for w_0 are larger than expected (~ 10 - 15 m/s).

473 Momentum fluxes, $\langle u'w' \rangle$ and $\langle v'w' \rangle$, have been estimated after removal of hor-
 474 izontal and vertical mean winds that were fitted together with the wind gradients. Com-

475 pared to some previous studies, our momentum flux estimates exhibit larger amplitudes,
 476 which indicates that the GW activity in the MLT over southern Patagonia is very strong.
 477 The statistical uncertainties of $\langle u'w' \rangle$ and $\langle v'w' \rangle$ were also included in the analysis.
 478 The latter reveal that in order to have statistically significant momentum flux estimates,
 479 one should consider averages of at least 28 days. In this way, our results must be con-
 480 sidered as representative of monthly mean momentum fluxes, driven by waves with pe-
 481 riods shorter than 4 hours, vertical wavelengths shorter than 4 km, and horizontal scales
 482 less than 400 km.

483 We are confident that SIMONe Argentina has also the potential to investigate non-
 484 zero lag second-order statistics of MLT wind velocities, e.g., by using correlation func-
 485 tion techniques such as those presented in Vierinen et al. (2019). This will be explored
 486 in future studies. Besides, we also plan to investigate momentum flux estimates with-
 487 out averaging over long periods of time, provided there is evidence of specific (determi-
 488 nistic) wave events occurring in the troposphere/stratosphere.

489 Acknowledgments

490 The authors deeply thank Nico Pfeffer, Matthias Clahsen, Jens Wedrich and Thomas
 491 Barth (at IAP) for their invaluable help and expertise during the installation of SIMONe
 492 Argentina; and Jonathan Quiroga and Nahuel Daz (CONICET) for helping with the lo-
 493 gistics prior to the installation. The authors would also like to thank Daro Godoy and
 494 Facundo Olivares (at Tres Lagos); Pablo Quiroz (at La Estela); Martn el griego Palopoli
 495 and National Parks staff (at El Chaltn); the Consejo Agrario de la Provincia de Santa
 496 Cruz (at Gob. Gregores and El Calafate); and the BAM-GAL (Argentina Air Force) for
 497 their support and help in maintaining the sites.

498 This work was partially supported by the Deutsche Forschungsgemeinschaft (DFG,
 499 German Research Foundation) under SPP 1788 (DynamicEarth) CH 1482/2, and by the
 500 Bundeministerium für Bildung und Forschung via project WASCLIM-IAP part of the
 501 ROMIC-II program.

502 The data used to produce the figures can be obtained in HDF5 format from [ftp://ftp.iap-](ftp://ftp.iap-kborn.de/data-in-publications/ConteESS2020)
 503 [kborn.de/data-in-publications/ConteESS2020](ftp://ftp.iap-kborn.de/data-in-publications/ConteESS2020).

504 References

- 505 Andrews, D. G., Holton, J. R., & Leovy, C. B. (1987). Middle atmosphere dynamics.
 506 *Academic Press, San Diego, California*. doi: 10.1002/qj.49711548612
- 507 Andrioli, V. F., Batista, P. P., Clemesha, B. R., Schuch, N. J., & Buriti, R. A.
 508 (2015). Multi-year observations of gravity wave momentum fluxes at low and
 509 middle latitudes inferred by all-sky meteor radar. *Ann. Geophys.*, *33*, 1183–
 510 1193. doi: 10.5194/angeo-33-1183-2015
- 511 Andrioli, V. F., Fritts, D. C., Batista, P. P., & Clemesha, B. R. (2013). Im-
 512 proved analysis of all-sky meteor radar measurements of gravity wave
 513 variances and momentum fluxes. *Ann. Geophys.*, *31*, 889–908. doi:
 514 10.5194/angeo-31-889-2013
- 515 Babu, V. S., Ramkumar, G., & John, S. R. (2012). Seasonal variation of plan-
 516 etary wave momentum flux and the forcing towards mean flow accelera-
 517 tion in the MLT region. *J. Atmos. Sol. Terr. Phys.*, *78-79*, 53–61. doi:
 518 10.1016/j.jastp.2011.05.010
- 519 Browning, K. A., & Wexler, R. (1968). The determination of kinematic properties of
 520 a wind field using Doppler radar. *J. Appl. Meteorol.*, *7*(1), 105–113.
- 521 Burnside, R. G., Herrero, F. A., Meriwether, J. W., & Walker, J. C. G. (1981). Op-
 522 tical observations of thermospheric dynamics at Arecibo. *J. Geophys. Res.*, *86*,
 523 5532–5540.

- 524 Chau, J. L., Hoffmann, P., Pedatella, N. M., Matthias, V., & Stober, G. (2015).
525 Upper mesospheric lunar tides over middle and high latitudes during sudden
526 stratospheric warming events. *J. Geophys. Res.: Space Phys.*, *120*(4), 3084–
527 3096. doi: 10.1002/2015JA020998
- 528 Chau, J. L., Stober, G., Hall, C. M., Tsutsumi, M., Laskar, F. I., & Hoffmann, P.
529 (2017). Polar mesospheric horizontal divergence and relative vorticity measure-
530 ments using multiple specular meteor radars. *Radio Sci.*, *52*, 811–828. doi:
531 10.1002/2016RS006225
- 532 Chau, J. L., Urco, J. M., Vierinen, J., Harding, B. J., Clahsen, M., Pfeffer, N., . . .
533 Erickson, P. J. (2020). Multistatic specular meteor radar network in Peru:
534 System description and initial results. *Earth and Space Science*.
- 535 Chau, J. L., Urco, J. M., Vierinen, J. P., Volz, R. A., Clahsen, M., Pfeffer, N.,
536 & Trautner, J. (2019). Novel specular meteor radar systems using co-
537 herent MIMO techniques to study the mesosphere and lower thermo-
538 sphere. *Atmospheric Measurement Techniques*, *12*(4), 2113–2127. doi:
539 10.5194/amt-12-2113-2019
- 540 Clemesha, B. R., Batista, P. P., Buriti da Costa, R. A., & Schuch, N. (2009). Sea-
541 sonal variations in gravity wave activity at three locations in Brazil. *Ann. Geo-*
542 *phys.*, *27*, 1059–1065. doi: 10.5194/angeo-27-1059-2009
- 543 Conde, M., & Smith, R. W. (1998). Spatial structure in the thermospheric horizon-
544 tal wind above Poker Flat, Alaska, during solar minimum. *J. Geophys. Res.*,
545 *103*, 9449–9472. doi: 10.1029/97JA03331
- 546 Conte, J. F., Chau, J. L., Laskar, F. I., Stober, G., Schmidt, H., & Brown, P.
547 (2018). Semidiurnal solar tide differences between fall and spring transi-
548 tion times in the Northern Hemisphere. *Ann. Geophys.*, *36*, 999–1008. doi:
549 10.5194/angeo-36-999-2018
- 550 Conte, J. F., Chau, J. L., & Peters, D. H. W. (2019). Middle- and high-latitude
551 mesosphere and lower thermosphere mean winds and tides in response to
552 strong polar-night jet oscillations. *J. Geophys. Res.*, *124*, 9262–9276. doi:
553 10.1029/2019JD030828
- 554 Conte, J. F., Chau, J. L., Stober, G., Pedatella, N., Maute, A., Hoffmann, P., . . .
555 Murphy, D. J. (2017). Climatology of semidiurnal lunar and solar tides at
556 middle and high latitudes: Interhemispheric comparison. *J. Geophys. Res.*,
557 *122*. doi: 10.1002/2017JA024396
- 558 de Wit, R. J., Janches, D., Fritts, D. C., & Hibbins, R. E. (2016). QBO modu-
559 lation of the mesopause gravity wave momentum flux over Tierra del Fuego.
560 *Geophys. Res. Lett.*, *43*, 4049–4055. doi: 10.1002/2016GL068599
- 561 Egito, F., Andrioli, V. F., & Batista, P. P. (2016). Vertical winds and momen-
562 tum fluxes due to equatorial planetary scale waves using all-sky meteor
563 radar over Brazilian region. *J. Atmos. Sol. Terr. Phys.*, *149*, 108–119. doi:
564 10.1016/j.jastp.2016.10.005
- 565 Ern, M., Preusse, P., Gille, J. C., Hepplewhite, C. L., Mlynyczak, M. G., Russell III,
566 J. M., & Riese, M. (2011). Implications for atmospheric dynamics derived
567 from global observations of gravity wave momentum flux in stratosphere and
568 mesosphere. *J. Geophys. Res.*, *116*. doi: 10.1029/2011JD015821
- 569 Fan, Y., Klostermeyer, J., & Rüster, R. (1991). VHF radar observation of grav-
570 ity wave critical levels in the mid-latitude summer mesopause region. *Geophys.*
571 *Res. Lett.*, *18*, 697–700.
- 572 Forbes, J. M. (1984). Middle atmosphere tides. *J. Atmos. Terr. Phys.*(46), 1049–
573 1067.
- 574 Fritts, D. C. (1984). Gravity wave saturation in the middle atmosphere: A review of
575 theory and observations. *Rev. Geophys.*, *22*, 275–308.
- 576 Fritts, D. C., & Alexander, M. J. (2003). Gravity wave dynamics and ef-
577 fects in the middle atmosphere. *Reviews of Geophysics*, *41*(1). doi:
578 10.1029/2001RG000106

- 579 Fritts, D. C., Iimura, H., Lieberman, R., Janches, D., & Singer, W. (2012a). A con-
580 jugate study of mean winds and planetary waves employing enhanced meteor
581 radars at Rio Grande, Argentina (53.8°S) and Juliusruh, Germany (54.6°N). *J.*
582 *Geophys. Res.*, *117*. doi: 10.1029/2011JD016305
- 583 Fritts, D. C., Janches, D., & Hocking, W. K. (2010). Southern Argentina Agile Me-
584 teor Radar: Initial assessment of gravity wave momentum fluxes. *J. Geophys.*
585 *Res.*, *115*. doi: 10.1029/2010JD013891
- 586 Fritts, D. C., Janches, D., Hocking, W. K., Mitchell, N. J., & Taylor, M. J. (2012b).
587 Assessment of gravity wave momentum flux measurement capabilities by me-
588 teor radars having different transmitter power and antenna configurations. *J.*
589 *Geophys. Res.*, *117*(D10108). doi: 10.1029/2011JD017174
- 590 He, M., Chau, J. L., Stober, G., Hall, C. M., Tsutsumi, M., & Hoffmann, P. (2017).
591 Application of Manley-Rowe relation in analyzing nonlinear interactions be-
592 tween planetary waves and the solar semidiurnal tide during 2009 sudden
593 stratospheric warming event. *J. Geophys. Res.: Space Phys.*, *122*(10), 10,783-
594 10,795. doi: 10.1002/2017JA024630
- 595 Hines, C. O. (1988). A modeling of atmospheric gravity waves and wave drag gener-
596 ated by isotropic and anisotropic terrain. *J. Atmos. Sci.*, *45*, 309–322.
- 597 Hocking, W. K. (2005). A new approach to momentum flux determinations using
598 SKiYMET meteor radars. *Annales Geophysicae*, *23*(7), 2433–2439. doi: 10
599 .5194/angeo-23-2433-2005
- 600 Hocking, W. K., Fuller, B., & Vandeppeer, B. (2001). Real-time determination of
601 meteor-related parameters utilizing modern digital technology. *J. Atmos. Sol.-*
602 *Terr. Phy.*(63), 155–169. doi: 10.1016/S1364-6826(00)00138-3
- 603 Hocking, W. K., & Thayaparan, T. (1997). Simultaneous and collocated observa-
604 tions of winds and tides by MF and meteor radars over London, Canada (43°N,
605 81°W), during 1994/1996. *Radio Science*, *32*(2), 833–865.
- 606 Hoffmann, P., Becker, E., Singer, W., & Placke, M. (2010). Seasonal variation of
607 mesospheric waves at northern middle and high latitudes. *J. Atmos. Sol.-Terr.*
608 *Phys.*, *72*(14-15), 1068–1079. doi: 10.1016/j.jastp.2010.07.002
- 609 Holdsworth, D. A., Reid, I. M., & Cervera, M. A. (2004). Buckland park allsky in-
610 terferometric meteor radar. *Radio Sci.*, *39*(5). doi: 10.1029/2003RS003014
- 611 Jacobi, C., Portnyagin, Y. I., Solovjova, T. V., Hoffmann, P., Singer, W., Fahrut-
612 dinova, A. N., ... Meek, C. E. (1999). Climatology of the semidiurnal tide
613 at 5256°N from ground-based radar wind measurements 1985/1995. *J. Atmos.*
614 *Sol.-Terr. Phys.*, *61*(13), 975–991. doi: 10.1016/S1364-6826(99)00065-6
- 615 Jia, M., Xue, X., Gu, S., Chen, T., Ning, B., Wu, J., ... Dou, X. (2018). Multiyear
616 observations of gravity wave momentum fluxes in the midlatitude mesosphere
617 and lower thermosphere region by meteor radar. *J. Geophys. Res.*, *123*, 5684–
618 5703. doi: 10.1029/2018JA025285
- 619 Jones, J., Webster, A. R., & Hocking, W. K. (1998). An improved interferome-
620 ter design for use with meteor radars. *Radio Sci.*, *33*, 55–65. doi: 10.1029/
621 97RS03050
- 622 Kudeki, E., & Franke, S. J. (1998). Statistics of momentum flux estimation. *J. At-*
623 *mos. Sol.-Terr. Phys.*, *60*, 1549–1553.
- 624 Kumar, K. K., Subrahmanyam, K. V., Mathew, S. S., Koushik, N., & Ramkummar,
625 G. (2018). Simultaneous observations of the quasi 2-day wave climatology over
626 the low and equatorial latitudes in the mesosphere lower thermosphere. *Clim.*
627 *Dyn.*, *51*, 221–233. doi: 10.1007/s00382-017-3916-2
- 628 Laskar, F. I., Chau, J. L., Stober, G., Hoffmann, P., Hall, C. M., & Tsutsumi, M.
629 (2016). Quasi-biennial oscillation modulation of the middle- and high-latitude
630 mesospheric semidiurnal tides during August/September. *J. Geophys. Res.:
631 Space Phys.*, *121*(5), 4869–4879. doi: 10.1002/2015JA022065
- 632 Lindzen, R. S., & Chapman, S. (1969). Atmospheric tides. *Space Sci. Rev.*, *10*, 3–
633 188.

- 634 Liu, A. Z., Lu, X., & Franke, S. J. (2013). Diurnal variation of gravity wave mo-
635 mentum flux and its forcing on the diurnal tide. *J. Geophys. Res.*, *118*, 1668–
636 1678.
- 637 Liu, H.-L., & Roble, R. G. (2002). A study of a self-generated stratospheric sudden
638 warming and its mesospheric-lower thermospheric impacts using the coupled
639 TIME-GCM/CCM3. *J. Geophys. Res.*, *107*. doi: 10.1029/2001JD001533
- 640 Lund, T., Fritts, D., Kam, W., Laughman, B., & Liu, H.-L. (2020). Numerical sim-
641 ulation of mountain waves over the Southern Andes, Part 1: Mountain wave
642 and secondary wave character, evolutions, and breaking. *J. Atm. Sci.*.
- 643 McCormack, J. P., Coy, L., & Singer, W. (2014). Intraseasonal and inter-
644 annual variability of the quasi 2 day wave in the Northern Hemisphere
645 summer mesosphere. *J. Geophys. Res. Atmos.*, *119*, 2928–2946. doi:
646 10.1002/2013JD020199
- 647 Meriwether, J., Faivre, M., Fesen, C., Sherwood, P., & Veliz, O. (2008). New results
648 on equatorial thermospheric winds and the midnight temperature maximum.
649 *Ann. Geophys.*, *26*, 447–466.
- 650 Murphy, D. J., Forbes, J. M., Walterscheid, R. L., Hagan, M. E., Avery, S. K.,
651 Aso, T., . . . Vincent, R. A. (2006). A climatology of tides in the antarc-
652 tic mesosphere and lower thermosphere. *J. Geophys. Res.*, *111*. doi:
653 10.1029/2005JD006803
- 654 Offermann, D., Hoffmann, P., Knieling, P., Koppmann, R., Oberheide, J., Rigglin,
655 D. M., . . . Steinbrecht, W. (2011). Quasi 2 day waves in the summer meso-
656 sphere: Triple structure of amplitudes and longterm development. *J. Geophys.*
657 *Res.*, *116*. doi: 10.1029/2010JD015051
- 658 Orte, P. F., Wolfram, E., Salvador, J., Mizuno, A., Bègue, N., Bencherif, H., . . .
659 Quiroga, J. (2019). Analysis of a southern sub-polar short-term ozone varia-
660 tion event using a millimetre-wave radiometer. *Ann. Geophys.*, *37*, 613–629.
661 doi: 10.5194/angeo-37-613-2019
- 662 Pancheva, D., & Mukhtarov, P. (2011). Atmospheric tides and planetary waves:
663 recent progress based on SABER/TIMED temperature measurements
664 (20022007). *Aeronomy of the Earths Atmosphere and Ionosphere, IAGA*
665 *Special Sopron Book Series 2*. doi: 10.1007/978-94-007-0326-1_2
- 666 Piani, C., Durran, D., Alexander, M. J., & Holton, J. R. (2000). A numerical study
667 of three-dimensional gravity waves triggered by deep tropical convection and
668 their role in the dynamics of the QBO. *J. Atmos. Sci.*, *57*, 3689–3702. doi:
669 10.1029/2019EA000570
- 670 Placke, M., Hoffmann, P., Becker, E., Jacobi, C., Singer, W., & Rapp, M. (2011b).
671 Gravity wave momentum fluxes in the MLTPart II: Meteor radar investiga-
672 tions at high and midlatitudes in comparison with modeling studies. *J. Atmos.*
673 *Sol. Terr. Phys.*, *73*, 911–920. doi: 10.1016/j.jastp.2010.05.007
- 674 Placke, M., Hoffmann, P., Latteck, R., & Rapp, M. (2015). Gravity wave momen-
675 tum fluxes from MF and meteor radar measurements in the polar MLT region.
676 *J. Geophys. Res.*, *120*, 736–750. doi: 10.1002/2014JA020460
- 677 Placke, M., Stober, G., & Jacobi, C. (2011a). Gravity wave momentum fluxes in the
678 MLTPart I: seasonal variation at Collm (51.3°N, 13.0°E). *J. Atmos. Sol. Terr.*
679 *Phys.*, *73*, 904–910. doi: 10.1016/j.jastp.2010.05.007
- 680 Rossby, C.-G. (1939). Relation between variations in the intensity of the zonal cir-
681 culation of the atmosphere and the displacements of the semipermanent centers
682 of action. *J. Mar. Res.*, *2*, 38–55.
- 683 Sandford, D. J., Muller, H. G., & Mitchell, N. J. (2006). Observations of lunar tides
684 in the mesosphere and lower thermosphere at Arctic and middle latitudes. *At-*
685 *mos. Chem. Phys.*, *6*, 4117–4127.
- 686 Siddiqui, T. A., Maute, A., Pedatella, N., Yamazaki, Y., Lühr, H., & Stolle, C.
687 (2018). On the variability of the semidiurnal solar and lunar tides of the equa-
688 torial electrojet during sudden stratospheric warmings. *Ann. Geophys.*, *36*,

- 1545–1562. doi: 10.5194/angeo-2018-80
- 689 Smith, A. K. (2012). Global dynamics of the MLT. *Surveys in Geophysics*, *33*(6),
690 1177–1230.
- 691
- 692 Spargo, A. J., Reid, I. M., & MacKinnon, A. D. (2019). Multistatic meteor
693 radar observations of gravity-wavetidal interaction over southern Aus-
694 tralia. *Atmospheric Measurement Techniques*, *12*, 4791–4812. doi:
695 10.5194/amt-12-4791-2019
- 696 Stening, R. J., Forbes, J. M., Hagan, M. E., & Richmond, A. D. (1997). Exper-
697 iments with a lunar atmospheric tidal model. *J. Geophys. Res.*, *102*, 13,465–
698 13,472. doi: 10.1029/97JD00778
- 699 Stober, G., & Chau, J. L. (2015). A multistatic and multifrequency novel approach
700 for specular meteor radars to improve wind measurements in the mlt region.
701 *Radio Sci.*, *50*, 431–442. doi: 10.1002/2014RS005591
- 702 Thorsen, D., Franke, S. J., & Kudeki, E. (1997). A new approach to MF radar inter-
703 ferometry for estimating mean winds and momentum flux. *Radio Sci.*, *32*(2),
704 707–726.
- 705 Trinh, Q. T., Ern, M., Doornbos, E., Preusse, P., & Riese, M. (2018). Charac-
706 teristics of the quiettime hot spot gravity waves observed by GOCE over
707 the Southern Andes on 5 July 2010. *Ann. Geophys.*, *36*, 425–444. doi:
708 10.5194/angeo-36-425-2018
- 709 Urco, J. M., Chau, J. L., Milla, M. A., Vierinen, J., & Weber, T. (2018). Coherent
710 MIMO to improve aperture synthesis radar imaging of field-aligned irregulari-
711 ties: First results at Jicamarca. *IEEE Transactions on Geoscience and Remote
712 Sensing*, *PP*(99), 1–11. doi: 10.1109/TGRS.2017.2788425
- 713 Urco, J. M., Chau, J. L., Weber, T., & Latteck, R. (2019). Enhancing the spatio-
714 temporal features of polar mesosphere summer echoes using coherent MIMO
715 and radar imaging at MAARSY. *Atmospheric Measurement Techniques*, *12*,
716 955–969. doi: 10.5194/amt-12-955-2019
- 717 Vadas, S. L., & Becker, E. (2019). Numerical modeling of the generation of tertiary
718 gravity waves in the mesosphere and thermosphere during strong mountain
719 wave events over the Southern Andes. *J. Geophys. Res.*, *124*, 7687–7718. doi:
720 10.1029/2019JA026694
- 721 Vadas, S. L., Xu, S., Yue, J., Bossert, K., Becker, E., & Baumgarten, G. (2019a).
722 Characteristics of the quiettime hot spot gravity waves observed by GOCE
723 over the Southern Andes on 5 July 2010. *J. Geophys. Res.*, *124*, 7034–7061.
724 doi: 10.1029/2019JA026693
- 725 Vierinen, J., Chau, J. L., Charuvil, H., Urco, J. M., Clahsen, M., Avsarkisov, V.,
726 ... Volz, R. (2019). Observing mesospheric turbulence with specular meteor
727 radars: A novel method for estimating second-order statistics of wind velocity.
728 *Earth and Space Science*, *6*, 1171–1195. doi: 10.1029/2019EA000570
- 729 Vierinen, J., Chau, J. L., Pfeffer, N., Clahsen, M., & Stober, G. (2016). Coded con-
730 tinuous wave meteor radar. *Atmospheric Measurement Techniques*, *9*(2), 829–
731 839.
- 732 Vincent, R. A., Kovalam, S., Reid, I. M., & Younger, J. P. (2010). Gravity wave
733 flux retrievals using meteor radars. *Geophys. Res. Lett.*, *37*. doi: 10.1029/
734 2010GL044086
- 735 Wilhelm, S., Stober, G., & Brown, P. (2019). Climatologies and long-term
736 changes in mesospheric wind and wave measurements based on radar ob-
737 servations at high and mid latitudes. *Ann. Geophys.*, *37*, 851–875. doi:
738 10.5194/angeo-37-851-2019
- 739 Yamazaki, Y., Matthias, V., Miyoshi, Y., Stolle, C., Siddiqui, T., Kervalishvili, G.,
740 ... Alken, P. (2020). September 2019 antarctic sudden stratospheric warming:
741 Quasi-6-day wave burst and ionospheric effects. *Geophys. Res. Lett.*, *47*. doi:
742 10.1029/2019GL086577

Figure 1.

SIMONe Argentina sites

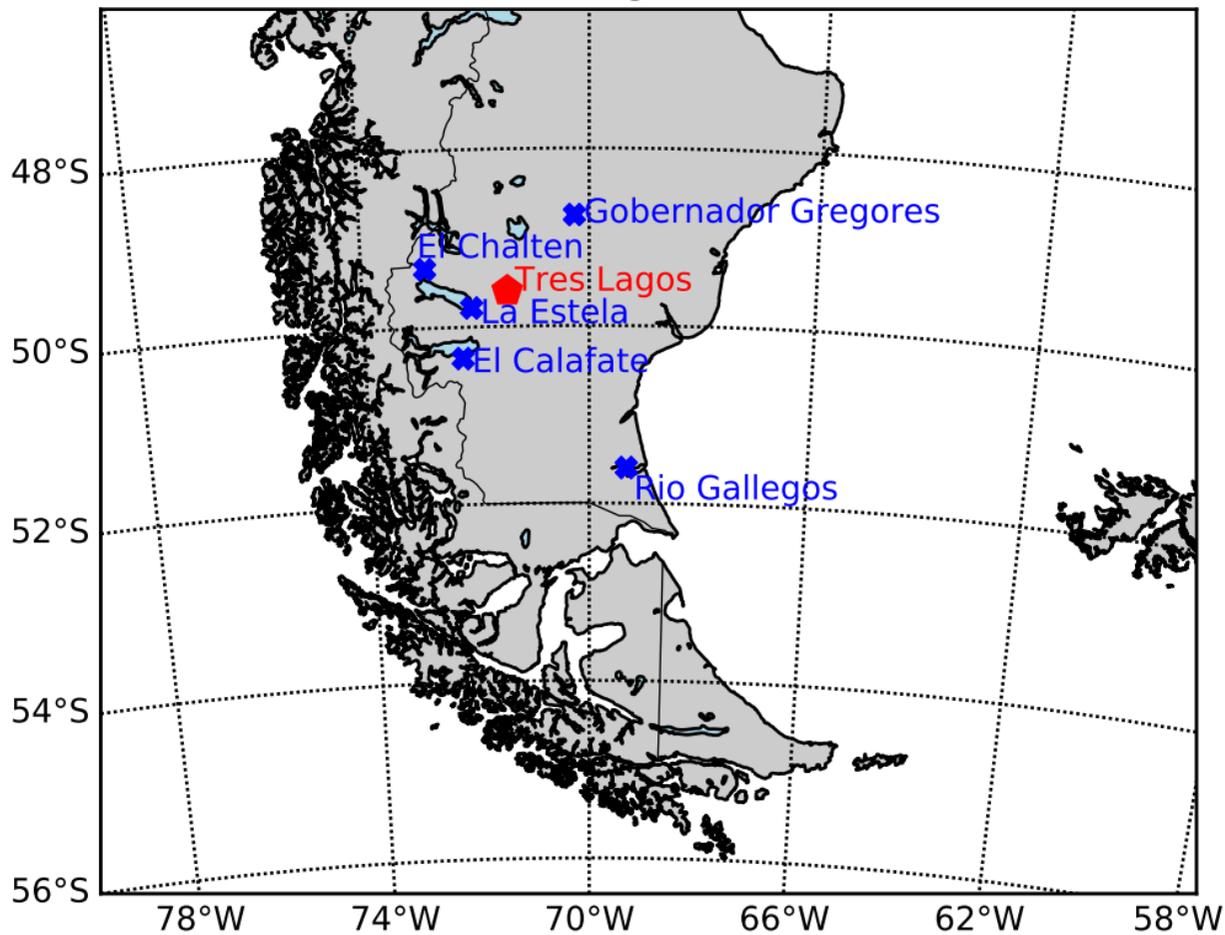


Figure 2.

SIMONE Argentina detection statistics

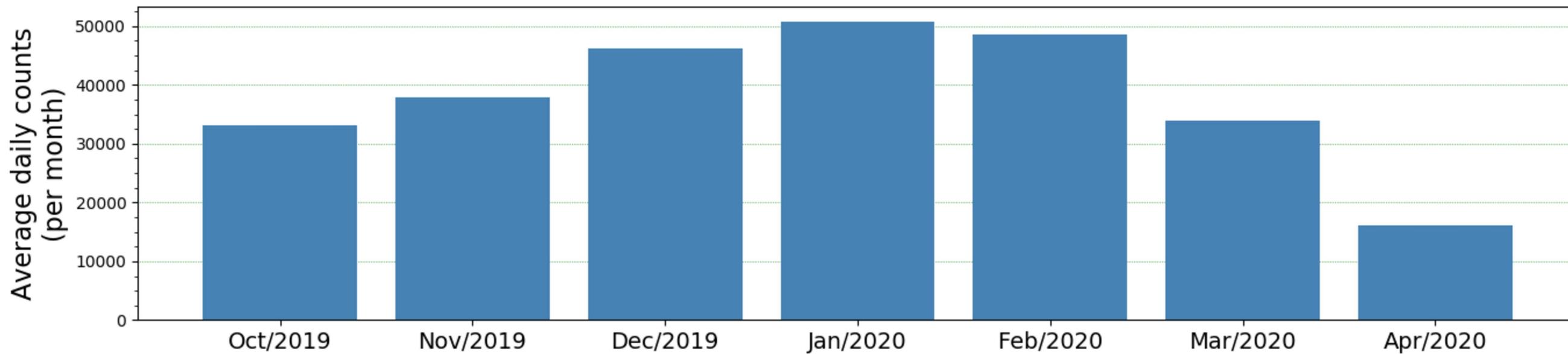
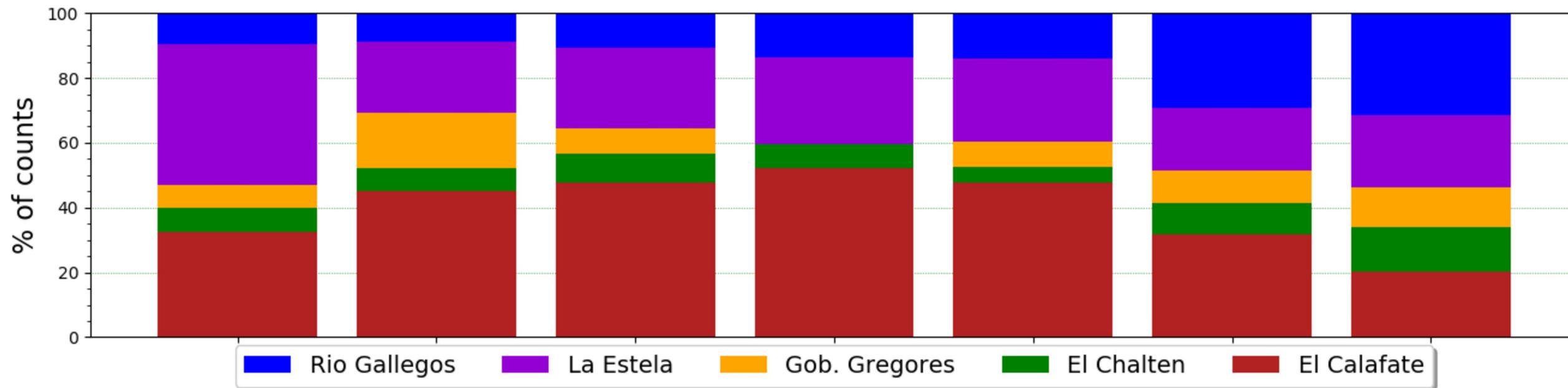


Figure 3.

SIMONE Argentina - Mean winds & tides

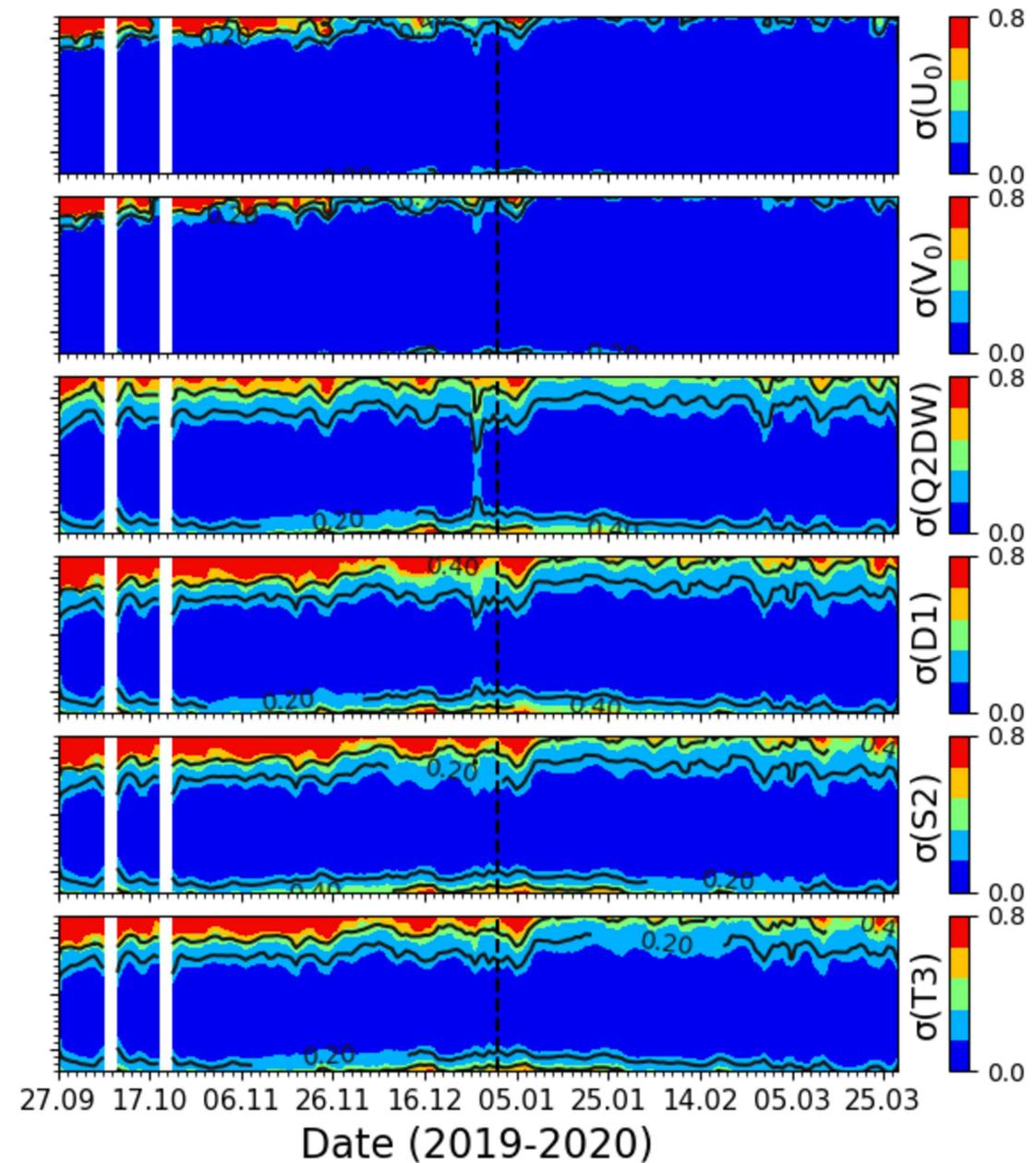
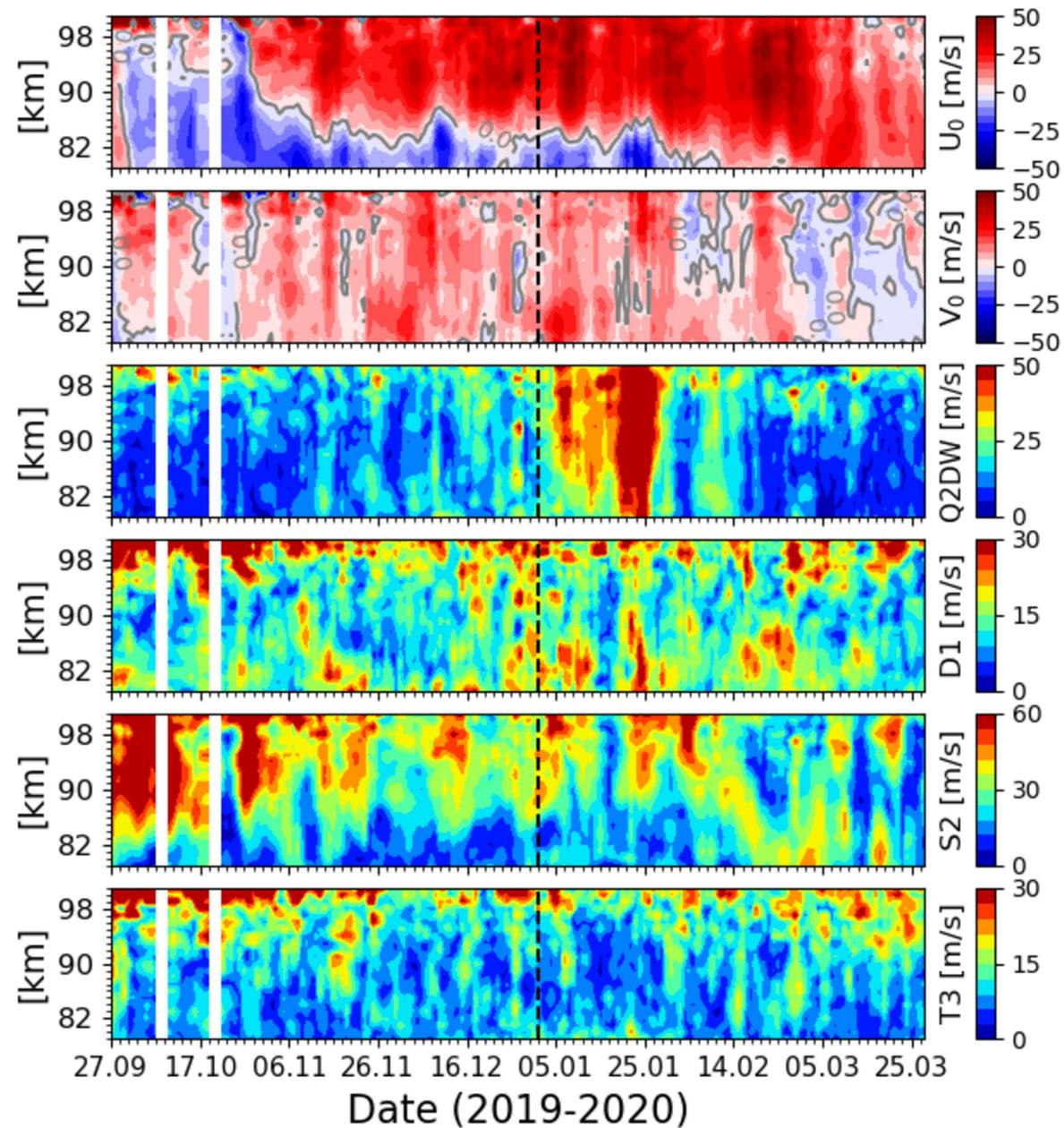


Figure 4.

Daily mean of 4-h winds and gradients

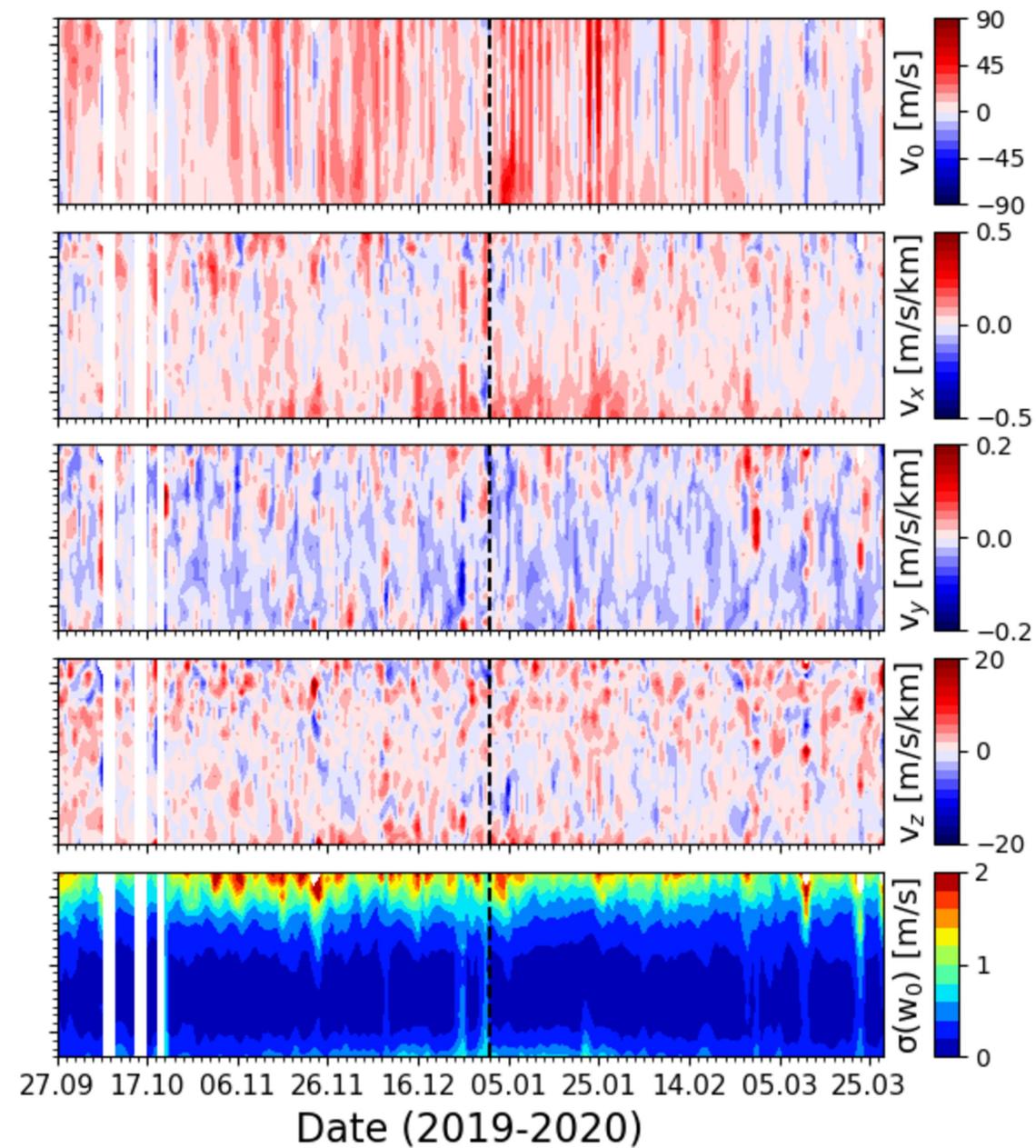
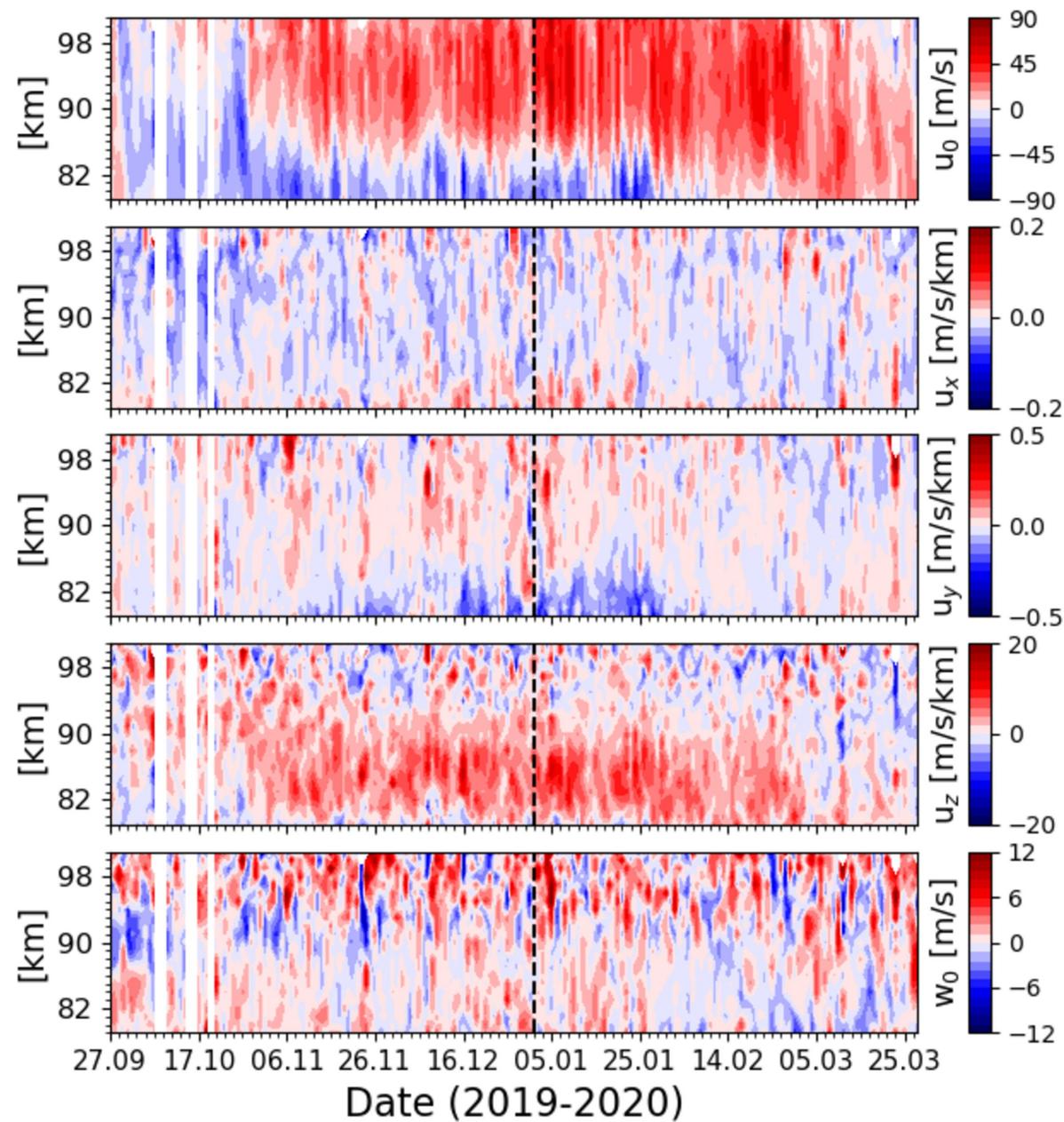


Figure 5.

With (4h, 4km) winds & gradients / 28-day average

