

Semi-automatic Ice Rafted Debris quantification with Computed Tomography

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Key Points:

- CT counting reproduces the known number of particles added to phantom archives.
- CT counting tracks the main trends in published IRD profiles from sediment archives.
- CT counting improves the sampling resolution to resolve higher-frequency variability.

Abstract

Sedimentary Ice Rafted Debris (IRD) provides critical information about the climate sensitivity and dynamics of ice sheets. In recent decades, high-resolution investigations have revealed ice rafting events in response to rapid warming: such reconstructions help us constrain the near-future stability of our planet's fast-changing cryosphere. However, similar efforts require laborious and destructive analytical procedures to separate and count IRD. Computed Tomography (CT) holds great promise to overcome these impediments to progress by enabling the micrometer scale visualization of individual IRD grains. This study demonstrates the potential of this emerging approach by **1)** validating CT counts in synthetic sediment archives (phantoms) spiked with a known number of grains, **2)** replicating published IRD stratigraphies, and **3)** improving sampling resolution. Our results show that semi-automated CT counting of grains in the common 150-500 μm size fraction reproduces actual particle numbers and tracks manually counted trends. We also find that differences between manual and CT-counted data are explained by image processing artifacts, offsets in sampling resolution and bioturbation. By acquiring these promising results using basic image processing tools, we argue that our work advances and broadens the applicability of ultra-high resolution IRD counting with CT to deepen our understanding of ice sheet-climate interactions on human-relevant timescales.

Plain Language Summary

Chunks of ice regularly break off glaciers floating in the ocean. These icebergs contain rock fragments picked up during the journey from land to water. As icebergs drift into warmer waters and melt, this rubble sinks to the bottom and settles on the ocean floor. Detection of these particles in marine sediments thus provide evidence that glacial ice reached down to sea-level. The flux of this ice rafted debris (IRD) gives researchers information about the past behavior of glaciers. As our planet warms, melting glaciers have become important drivers of sea-level rise. IRD studies can therefore help us better adapt to rising sea levels. But to do so on timescales relevant for humans, researchers have to extract thousands of samples from meters of sediment and sieve out IRD grains before manually counting them. Faster approaches would greatly ease the workload. In this study, we present a promising way to do so with the help from a medical technique: Computed Tomography (CT). Our findings show it is possible to semi-automatically count sand-sized grains from CT imagery without touching or destroying samples. We also show that this can be done with simple processing steps accessible to non-experts.

1 Introduction

Along glaciated margins, the calving and rafting of melting icebergs from marine-terminating glaciers deliver Ice Rafted Debris (IRD) to the open ocean (Ruddiman, 1977). The presence and concentration of IRD grains in marine sediment sequences provides critical information about ice sheet dynamics (Andrews, 2000). Over the past decades, such investigations have revealed enigmatic phases of millennial-scale ice sheet instability – notably Heinrich (H) events, Dansgaard-Oeschger (D-O) cycles and Bond events (Bond et al., 1992, Dansgaard et al., 1993, Heinrich, 1988) – which have attracted significant research activity. Greater spatial coverage and a higher sampling resolution of IRD reconstructions allow us to better understand the pattern, pace and causes of these extreme events to better assess future ice sheet stability (e.g. Hemming, 2004).

Such efforts are, however, hampered by the time-consuming laboratory work that is required to separate IRD grains from background sediments, and subsequently count individual particles. Typical steps include multiple rounds of manually weighing, and sieving material into different grain size fractions. In addition, size requirements often limit the sampling resolution of records, while the counting of split samples due to time constraints may introduce uncertainty (e.g. Van der Plas & Tobi, 1965). Evidently, (semi)-automated non-destructive approaches have significant potential to advance the field by **1)** reducing analysis time, **2)** improving sampling resolution, and **3)** preserving valuable core material for other analyses. Over the past decades, researchers have proposed various approaches to do so, and key examples include the use of semi-automated particle size counting or the investigation of 2-D and 3-D X-Ray images (e.g., Andrews et al., 1997, Becker et al., 2018, , Ekblom Johansson et al., 2020, Jennings et al., 2018, Grobe, 1987).

However, while the semi-automated approach (e.g. Becker et al. 2018) is destructive and requires a series of manual steps, others only target the coarsest size fraction or rely on 2-D imagery (e.g. Grobe, 1987) so that counts are not reported per weight or volume as is customary in the literature.

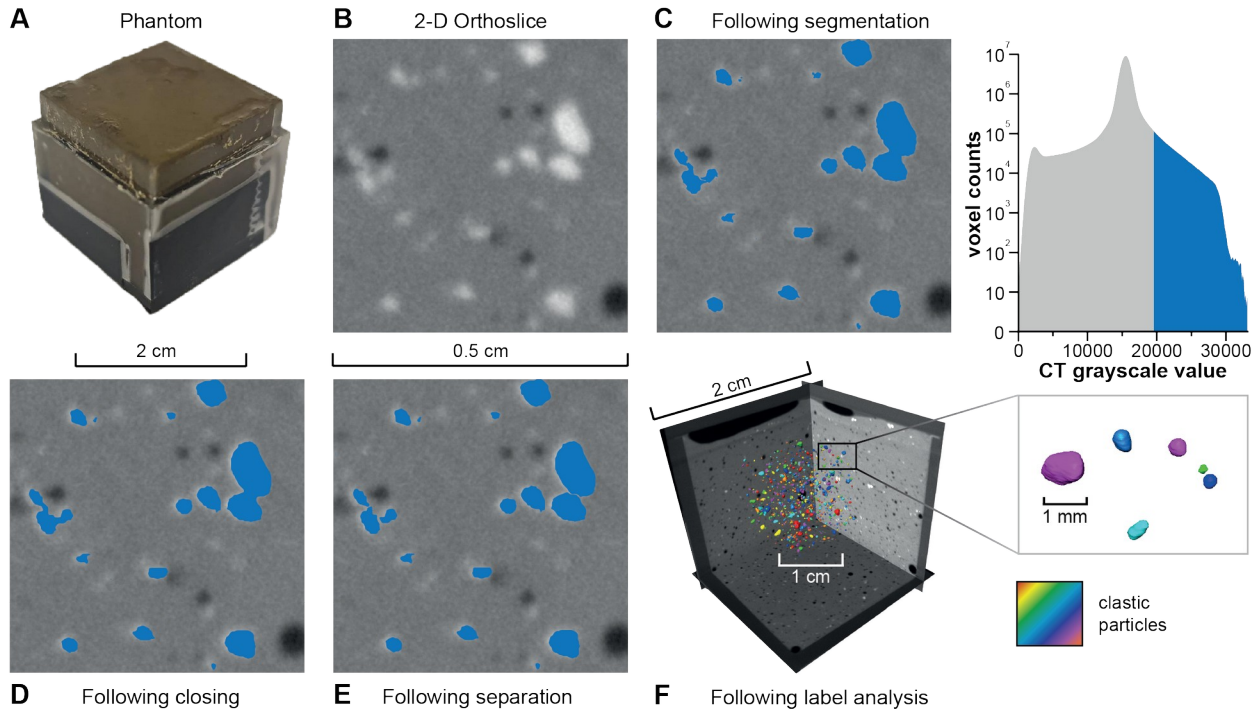
This study explores the potential of 3-D X-Ray Computed Tomography (CT) to overcome the foresaid limitations. This approach distinguishes grains from host sediment based on density differences. Recent increases in resolution and sampling size have shown great promise to detect and count barely visible particles in sediment volumes (e.g. Fouinat et al., 2017, Hodell et al., 2017, Røthe et al., 2018). Here, we advance the ability of CT to semi-automatically detect and count IRD particles by **1)** designing an experiment based on synthetic sediment records spiked with varying, but known, number of particles of the commonly analyzed 150-500 μm size fraction, **2)** validating our experimental findings by comparing CT and manual particle counts on published conventionally analyzed IRD records, and **3)** demonstrating that high-resolution CT counts capture high-frequency variability that is not captured by standard manual sampling protocols.

2 Materials and Methods

2.1 Experimental design

To explore the capability of CT to detect and count IRD particles within a sediment matrix, we designed a controlled experiment using synthetic sediment archives (phantoms). For this purpose, we filled 20 standard 8 cm^3 plastic cubes with a calculated number of 150-500 μm

119 grains mixed in a marine sediment matrix (Fig. 1A). To cover the typical range of IRD
 120 concentrations identified in published reconstructions, we added circa 25, 100, 500, 1000 and
 121 2000 grains per gram of dry weight (g^{-1} dry sediment).



122 **Figure 1.** Imagery that highlights key steps of our experimental approach. (a) close-up of one of
 123 the created synthetic records - phantoms. (b) a raw 0.25 cm^2 2-D cross-section (orthoslice) from
 124 one of our phantoms – note how dense radiopaque (light) clastic particles stand out. (c)
 125 thresholded particles following iterative segmentation (see section 2.4). The histogram on the
 126 right shows the applied CT grayscale value thresholds. (d) Restoration of fuzzy object
 127 boundaries (see section 2.4). (e) Separation of adjoining particles. (f) Individually classified
 128 (colored) clastic particles in a 1 cm^3 3-D visualization (reconstruction) used for subsequent
 129 sieving and counting.

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132 Counting three extracts of a known weight in triplicate, with the help of a Leica MZ6 optical
 133 microscope under $\times 40$ magnification, allowed us to establish a robust relation between weight

and particle counts whilst quantifying the human counting error. To assess the effect of lithological differences on our results, we created four sets of phantoms that were each spiked with bedrock types that are commonly found in major IRD source areas: quartz (density of 2.65 g/cm³), basalt (density of 3.0 g/cm³), dolomite (density of 2.8 g/cm³), and a 1:1:1 mixture of these materials (e.g., Bond et al., 1992, Jullien et al., 2006). For this purpose, we respectively relied on commercial quartz sand, basaltic floodplain sediments from northern Iceland and a dolomite laboratory standard (see Fig. S1). The mineralogy of each material was ascertained using a Bruker D8 ADVANCE ECO X-Ray diffractometer, equipped with a 1.5418 Å Copper (Cu) source operated at 40 kV/25 mA (see Fig. S2). *Sensu* Hemming (2004), we identify grains larger than 150 µm as IRD-sized and consequently only added material retained on a 150 µm mesh. To remove large visible grains, we also sieved out clasts larger than 2 mm. To assess the ability of CT imagery to reproduce sample particle size distributions, we determined the particle size distribution of each IRD-sized lithology in triplicate using a Mastersizer 3000. Samples were measured for 20 seconds at a stirring speed of 2500 rpm with ultrasound applied for 40 seconds prior to measurement.

The matrix of our phantoms derives from a pelagic multi core (GS15-198-62MC-F) retrieved off the Iceland Plateau (70°01'N 13°33'W) at 1423 m water depth (Jansen & Cruise-Members, 2015). To avoid introducing noise to the experiment, all IRD-sized particles were removed *a-priori* by sieving the sediment through a 63 µm mesh. Further, treatment with 1M acetic acid at 50 °C (until reaction ceased) dissolved *in-situ* calcite shells. To each sample box, we added approximately 6 grams of matrix mixed with 5 ml water to emulate the properties of natural marine sediments. In addition to the aforementioned known number of IRD-sized grains,

phantoms were spiked with ~600 foraminifera shells of arbitrary species larger than 150 μm from the Norwegian Sea (H. Halfidason, pers. comm.) to assess whether ubiquitous calcite shells introduce noise to CT IRD counts. The potential error margin related to loss of material during mixing and transfer of material was estimated at 0.31 g ($2\sigma = 0.61\text{g}$), by weighing the box after finalizing it. Finally, we ascertained the Dry Bulk Density (DBD) of our phantoms following the approach of Dean Jr (1974) to convert CT-counted particles per volume data to particles per gram of dry weight conform most studies.

2.2 Natural marine sediment cores

To further test the potential of CT-based IRD-sized particle detection and counting, we applied the insights gained from our phantom experiment on two published conventionally analyzed IRD stratigraphies (2 cm counting resolution). These encompass two segments of North Atlantic calypso cores that were extracted on-board the R/V G.O. Sars (Dokken & Cruise-Members, 2016, Jansen & Cruise-Members, 2015): 1) the 454-488.5 cm segment from core GS16-204-22CC-A ($58^\circ 2.830'\text{N}$, $47^\circ 2.360'\text{W}$: 3160 m water depth), which was previously investigated by Griem et al. (2019), and 2) the 231-281 cm section of GS16-204-18CC ($60^\circ 1.840'\text{N}$, $40^\circ 33.450'\text{W}$: 2220 m water depth) (Rutledal et al., 2020). As detailed in section 3.2, we manually re-counted the 267.5-280 cm interval of the latter archive continuously at 0.5 cm resolution. Both cores were primarily selected because they have been analyzed using standardized IRD counting methods, show distinct variability, focus on the 150-500 μm size range, and the number of counted particles falls within the range of our experimental design ($<2000\text{ IRD g}^{-1}$ dry sediment). To optimize scanning resolution by minimizing the distance between source and detector (see van der Bilt et al. 2021), we extracted 2 cm wide u-channels from both sediment cores for CT

scanning. As with our phantoms (see section 2.2), we relied on DBD measurements after Dean Jr (1974) to convert CT-derived counts per volume to particles per gram dry weight. To this end, we extracted one sample near the top and bottom of the homogenous scanned section from GS16-204-22CC, while extracting four samples from the investigated segment of GS16-204-18CC due to a lithological change at 259 cm core depth as reported in Dokken & Cruise-Members (2016).

2.3 CT scanning

Fundamentally, Computed Tomography (CT) can resolve objects based on differences in X-Ray absorption: X-Ray photons penetrate light (black; radiolucent) materials with ease, while radiation is absorbed by dense (white; radiopaque) matter like bone (Röntgen, 1896), or clastic particles (Fig. 1B). The degree of X-Ray attenuation is captured by grayscale values, which typically reflect material density (higher is denser). By rotating objects or an X-Ray source and detector, CT scanners generate large numbers of 2-D radiographs known as orthoslices from various angles. These images can be reconstructed to create 3-D visualizations or reconstructions (e.g. Kalender, 2011). In contrast with more established 2-D X-Ray-based IRD detection approaches (e.g. Grobe, 1987), this allows characterization and counting of particles per volume.

For this study, CT scanning was performed using a ProCon CT-ALPHA-CORE system located at the Earth Surface Sediment Laboratory (EARTH LAB) of the University of Bergen that is customized for whole-core (max. 150 cm) analysis (see e.g. van der Bilt et al., 2018). This one-of-its-kind 16-bit scanner is fitted with a 240 kV microfocus X-Ray source and 9 MP detector that move vertically while the scanned object rotates. All presented scans were scanned at 800 μ A and 100 kV with an exposure time of 334 ms to generate 1600 projections per rotation. This

relatively high current helps us minimize the imprint of photoelectric effect (Duliu, 1999). A physical 0.5 mm Cu filter was applied to reduce beam hardening effects (see Brooks & Di Chiro, 1976), as well as ring artifact correction and median filtering. Using 2 cm wide u-channels and boxes allowed us to optimize scanning resolution by minimizing the distance between source and detector, producing imagery at $\sim 21 \mu\text{m}$ isotropic voxel size.

2.4 CT processing

After scanning, CT projections were reconstructed for 3-D visualization with the Fraunhofer Voxel X-Ray Office software. To further minimize the imprint of CT artifacts like beam hardening or edge effects (e.g. Barrett & Keat, 2004 and section 2.4), we cropped $\sim 1 \text{ cm}^3$ volumes near the center of scanned boxes and 1 cm^2 wide sections of the u-channels. This step was performed in duplicate (henceforth referred to as samples A and B) to assess the representativeness of these 3-D cutouts. All subsequent image analyses were executed using version 9.1.1 of Thermo Scientific Avizo. To broaden the applicability of our approach, we relied on basic image processing techniques that are accessible to most geoscientists (see Fig. S3). All applied tools and modules are highlighted in *italics* below and briefly described to help users execute the same steps in other often-used image processing suites like ImageJ or VGStudio Max. We first applied an iterative routine using the *Colormap* editor to highlight clastic particles from background host sediments. This simple approach fundamentally relies on the subtle but measurable density differences between both materials and the shape of the clastic particles; as can be seen in Fig. 1C, the porous (water-soaked) matrix is significantly lighter (darker) than dense (white) clasts. We then isolated the designated CT density range using the *Interactive Threshold* segmentation tool. As can be seen in Fig. 1B, this binary image does not

adequately resolve the edge of clasts – a prerequisite when counting specific size fractions for IRD analysis. The observed noise is introduced by partial volume effects: the inter-voxel blurring of CT greyscale values along the steep density gradient between different materials (e.g. Glover & Pelc, 1980, Schlüter et al., 2010). To overcome this issue, we restored object boundaries with a combination of dilation and erosion using the *Closing* module as shown in Fig. 1D. Next, the *Separate Objects* module was applied to make sure that adjoining or coagulating particles are split as can be seen in Fig. 1E.

Following the above steps to detect and resolve particles, we individually characterized them for analysis with the *Label Analysis* module (Fig. 1F). During this step, the equivalent diameter and shape properties of each object was calculated using the *EqDiameter* and *Shape_Va3D* measures, along with the coordinates (*BaryCentre*) of grains. We used the *Shape_Va3D*-measurement to account for the fact that non-spherical objects may pass through a sieve mesh that is larger than their equivalent diameter if oriented towards their smallest projection (see e.g. Retsch, 2009). To do so, we normalized our 150 μm size threshold (see sections 2.1-2) against the degree of non-sphericity reflected by *Shape_Va3D* values >1 . 'Digital sieving' was performed using the *Sieve Analysis* module before summing up particle counts for each phantom and at 1 mm depth intervals in scanned core sections. Finally, we performed basic geostatistical analyses like re-sampling, correlation and linear regression using version 16 of the StataSE software.

3 Results and discussion

3.1 Experimental findings

3.1.1 Particle size analysis

The correct identification of a known number of particles within a specific size range is of fundamental importance to this study. Therefore, we compared CT particle counts on all phantoms spiked with ~ 1000 grains (>150 μm) per g^{-1} dry sediment to laser diffraction measurements of pure extracts of each lithology used to spike these synthetic archives (section 2.1). Intercomparability is aided by the fact that both these approaches calculate the equivalent diameter of a sphere with the same volume for each particle. CT counts were corrected with the *Shape_Va3D* measure to account for the possibility that non-spherical objects may pass through a sieve mesh that is larger than their equivalent diameter (see section 2.4 and Fig. S3). As can be seen in Fig. 2, there is close agreement between CT and laser-derived Particle Size Distributions (PSDs). These findings 1) strengthen our confidence that the CT processing steps applied in this study accurately constrain the distribution of size fractions commonly targeted for IRD analysis – a prerequisite for automatic counting, 2) open doors for future venture into non-destructive CT-based particle size analysis, and 3) highlight differences between PSDs of the lithologies used to spike our phantoms to help contextualize possible counting offsets in the following paragraphs.

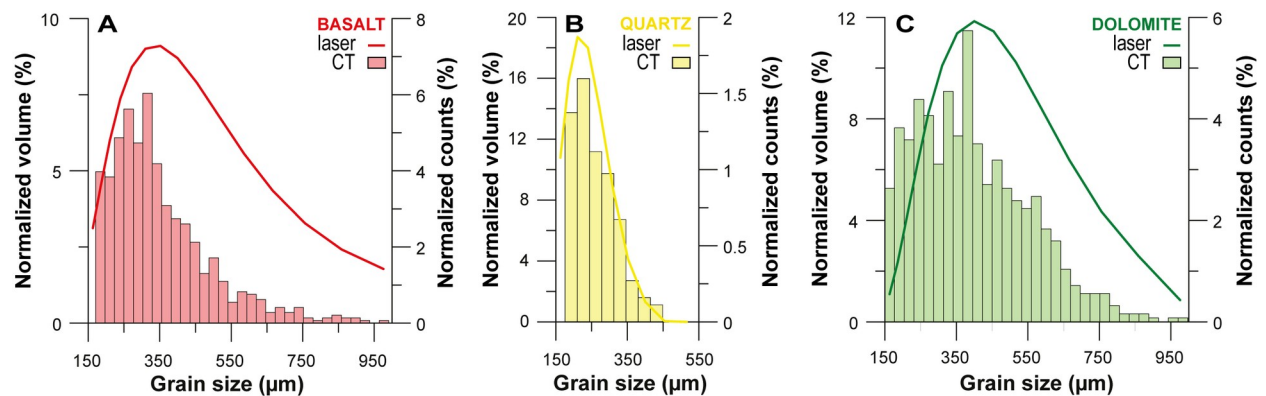


Figure 2. Comparing laser diffraction (Mastersizer: line) and CT-derived (bars) particle size distributions for IRD-sized grains of each bedrock type used to spike our phantoms: basalt (a),

quartz (**b**) and dolomite (**c**). The former data are expressed as a normalized volume (%), while the latter are calculated as normalized counts (%). See section 2.1 for additional details.

3.1.2 Lithic grain counting

As shown in Fig. 3, all linear regression fits between manual and CT counts in our phantoms are highly significant ($R^2 = 0.96-0.99$, $p = 0.00$), regardless of the lithology of added grains. Besides demonstrating the potential of CT scanning to automatically count IRD-sized particles, these findings also allay concerns that calcite shells introduce noise: the high reproducibility of lower counts in particular show that the ~600 foraminifera shells added to each phantom were not CT-counted. We attribute this to partial volume effects: while the density of calcium carbonate (2.7 g/cm^3) is near-identical to that of the rock types of added grains ($2.65-3 \text{ g/cm}^3$), voxel blurring with air-, water or matrix-filled chambers of foraminiferal tests yield a lower density (section 2.4).

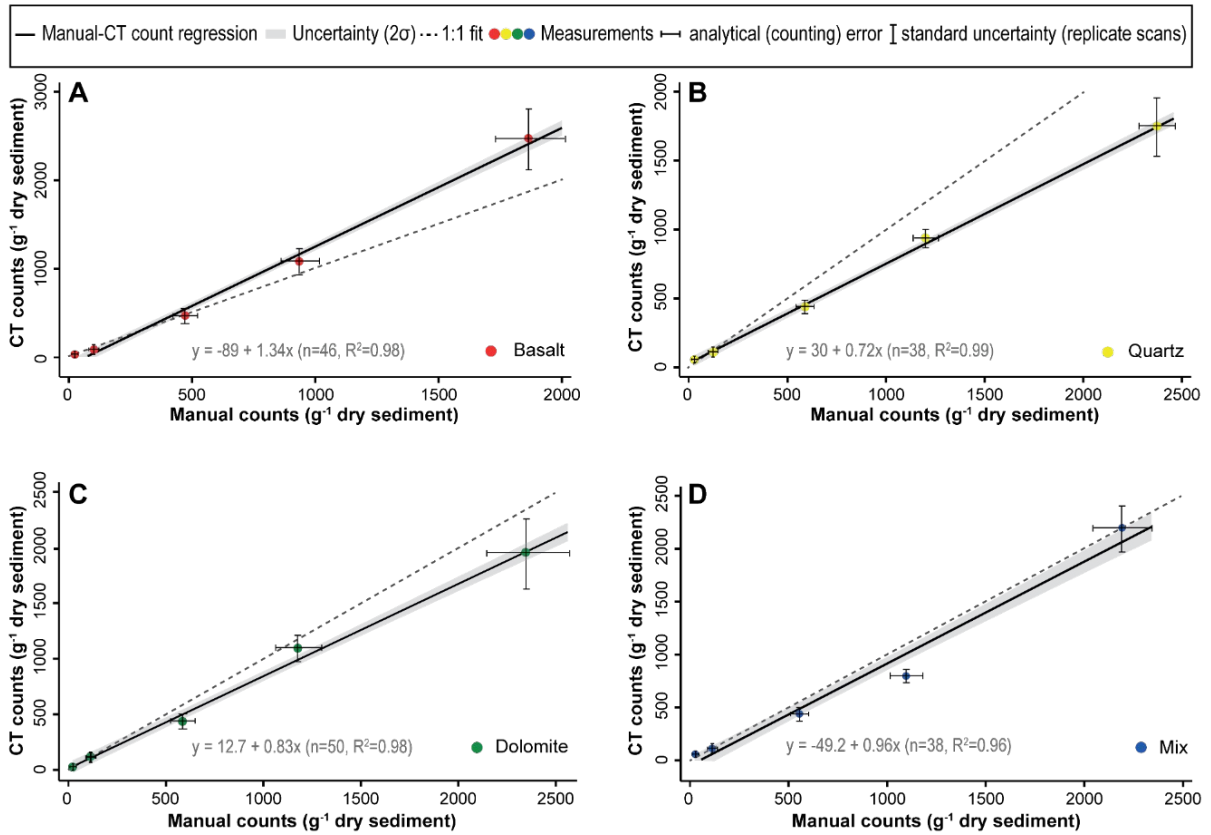


Figure 3. Linear regression fits (and summary statistics) between manual and CT counts of IRD-sized (150-500 μm) particles in synthetic sediment records (phantoms) spiked with basalt (a), Quartz (b), Dolomite (c) and a 1:1:1 mixture of each (d). See section 2.1 for additional details.

Our experimental findings compare favorably with previous efforts to (semi)-automatically count particles. Fouinat et al. (2017), for example, applied a similar CT-based approach to count larger mm-scale particles in a silty matrix, but derive a poorer fit ($R^2 = 0.66$, $p = 0.015$). We argue that this weaker correspondence can be primarily attributed to a lower scanning resolution of 0.25 mm versus 21 μm in our study (section 2.3). Becker et al. (2018) employed an approach based on automated microscopy to derive a marginally lower goodness-of-fit ($R^2 = 0.94$). However,

296 this approach still requires destructive (lower resolution) sampling and time-consuming wet
297 sieving.

298
299 However, while our experimental findings are promising, systematic offsets exist between
300 manual and CT counts. While highly significant, the slopes of all fits deviate from a 1:1 relation:
301 as seen in Fig. 3, these offsets often exceed calculated counting errors (sections 2.1 and 2.4) for
302 all lithologies, but especially for quartz. Here, we tentatively attribute these errors to a number of
303 analytical sources. Firstly, differences in the PSD of lithic grains. Assuming a unimodal
304 distribution (to estimate the proportion of sieved-out particles $<150\text{ }\mu\text{m}$ – see section 2.1), a
305 significant percentage of grains may be included or excluded when object boundaries (and thus
306 diameters) are incorrectly resolved during CT processing (see section 2.4 and Figs. 1C-E). This
307 source of error may well explain why offsets are largest for quartz as **1**) the median PSD of this
308 lithology sits closest to our $150\text{ }\mu\text{m}$ cut-off (Fig. 3B) so that small errors generate large count
309 uncertainties, and **2**) density differences with host sediments are smallest, which complicates our
310 efforts to accurately resolve object boundaries based on CT greyscale values (see section 2.4). In
311 addition, image processing may also impact CT particle counts by erroneously splitting
312 irregularly shaped grains into multiple objects with the *Separation* module (see Fig. 1E). This
313 source of error may help explain the observed overestimation of basaltic grains by CT counting
314 (Fig. 3A). This notion is supported by **1**) visual evidence of the irregular shape of these particles
315 (see Fig. S1), and **2**) their comparatively large size (Fig. 2A), which increases the probability that
316 erroneously split particles are included in the counted $>150\text{ }\mu\text{m}$ fraction. The applied *Closing*
317 *module* (see section 2.4) might also have exacerbated this effect as it may expand the size of
318 particles by smoothing uneven surfaces or filling in hollow particles (see Figs. 1C-D).

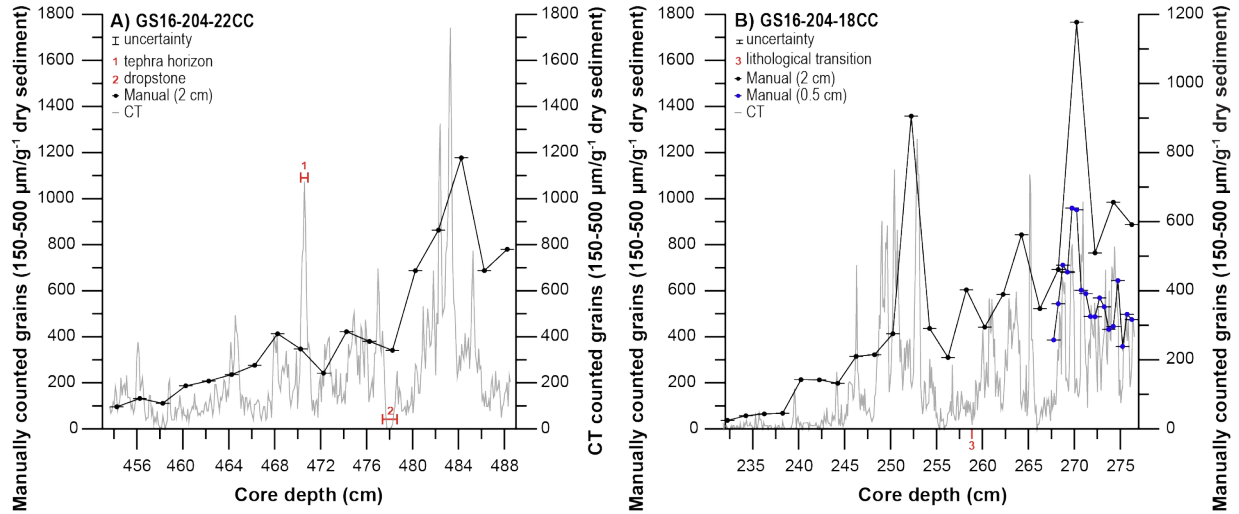


Figure 4. A comparison of manual and CT-derived counts of IRD-sized particles in two marine sediment segments. For (a) GS16-204-22CC and (b) GS16-204-18CC. Highlighted uncertainty intervals (grey) are based on the average offset in CT-derived grain counts between replicate 3-D samples A and B (see section 2.4). Horizontal bars on manual count symbols mark the 0.5 cm sampling width (see section 2.2). Red numerals indicate the stratigraphic position of marked features in the scanned core segments – 1) the NAAZ II tephra marker (Rutledal et al., 2020), 2) a cm-scale drop stone, and 3) a visible lithological transition (Dokken & Cruise-Members, 2016).

3.2 Application on manually counted natural sediment archives.

As can be seen in Fig. 4, our CT-based approach to count 150-500 μm particles capture most of the main IRD peaks in the manually counted records. The strength of this relation is confirmed by positive Spearman ρ values of 0.75 ($n=18$, $p=0.0003$) for core GS16-204-22CC and 0.63 ($n=25$, $p=0.0007$) for core GS16-204-18CC – all calculated on evenly (0.5 cm) resampled data. These findings clearly demonstrate the potential of our CT-based approach to semi-automatically

detect the 150-500 μm -sized particles that are typically targeted in IRD studies, even at comparatively low concentrations (max. 1800 grains/gr).

However, while certainly encouraging, the presented results also reveal substantial disparities. These can partly be explained by differences in sampling resolution: grains were CT-counted at 0.1 cm intervals, while 0.5 cm wide samples were taken every 2 cm for manual counts – smoothing out high-frequency (mm-scale) variability. As can be seen in Fig. 4B, an improved 0.5 cm sampling resolution greatly improves the agreement between manual and CT counts. Indeed, correlation of both datasets using the most similar CT-derived grain numbers within the 0.5 cm sampling width of manual samples yields a Spearman ρ of 0.96 ($n=18$, $p=0.0000$) – a result that equals the robustness of our experimental findings (see section 3.1.2). In addition, our scanned u-channel from GS16-204-22C contains two features that are also highlighted using the applied segmentation approach (see section 2.4) due to their highly similar density: the basaltic-component of a NAAZ II tephra deposit and a large drop stone (see Fig. 4A: 1 and 2). The latter is not CT-counted as its size falls outside our specified 150-500 μm grain range (see section 2.1), but its size simply leaves less space for other particles within the 0.1 cm^3 sample slice - creating a distinct minimum in counted particles. The tephra is captured by a sharp peak in the CT-counted IRD record. To remedy this, the characteristically high concentration of particles (ash shards) in tephra deposits may be highlighted using down-core variations in CT grayscale values as outlined by van der Bilt et al. (2021). The structural offset between CT and manual counts, which particularly affects GS16-204-18CC as seen in Fig. 4B, is more difficult to account for. As both cores were counted by the same analyst and derive from the same area (see section 2.1), we preclude differences in human counting error and lithology-specific analytical errors (see section

3.1.2) as plausible explanations. Because the bedrock geology of proximal IRD source areas in the region is dominated by quartz-rich metamorphic bedrock types (Dawes, 2009), it is worth noting that the offset between evenly sampled CT and manual counts in GS16-204-22CC is identical (28%) to the difference found in our quartz-spiked phantoms. But why this mismatch far greater in GS16-204-18CC, where our CT-based approach captures just 40% of manually counted grains (Fig. 4B)? We argue that the dissimilarity between both datasets may be attributed to disturbance introduced by bioturbation. In recent years, numerous researchers have harnessed various imaging techniques to demonstrate that burrowing may extensively modify the sediment structure and blur IRD signals (e.g., Dorador et al., 2014, Hodell et al., 2017). Indeed, Rutledal et al. (2020) relied on the same threshold-based segmentation routine presented in section 2.4 to highlight the presence of air-filled burrows in GS16-204-18CC and GS16-204-22CC. As can be seen in Fig. S4, these features are particularly extensive in the section of GS16-204-18CC scanned for this study. Furthermore, closer inspection of the appended X-Ray images shown in Fig. S4 reveals additional deformational structures that may represent infilled trace fossils or burrows. As these features are distributed both horizontally and vertically, the lateral offset between manually counted samples and scanned u-channels could have a major impact on down-core IRD profiles. To test this, we compared our CT data from GS16-204-18CC to higher-resolution manual counts performed on the same u-channel. As can be seen in Fig. 4B, the offset between these data is significantly smaller and is similar (39%) to the difference found between counts and scans in our quartz-spiked phantoms (28%).

4 Conclusions

This work underscores the potential of CT scanning for semi-automated and non-destructive counting of IRD-sized (150-500 μm) grains in sediment archives. Notwithstanding analytical errors that we ascribe to image processing artefacts, our experimental findings show that CT numbers capture more than 95% of grain count variability in homogenous phantoms. Also, by spiking each of these synthetic samples with a known number of foraminiferal tests, we allay concerns that (often-ubiquitous) calcite shells of a similar size and density affect CT IRD counts. Despite evidence of bioturbation and differences in sampling resolution, CT-derived counts strongly correlate ($\rho = 0.63\text{-}0.75$) with manual IRD profiles in both scanned core sections. Moreover, quadrupling our manual counting resolution on CT-scanned u-channels minimizes offsets between both datasets ($\rho = 0.96$). This somewhat surprising result suggests that mm-scale CT variations capture a signal rather than noise and highlights how bioturbation may modify IRD profiles. Importantly, all our results were acquired using basic image processing techniques that can be quickly mastered by most geologists. Following from the above, we argue that the presented CT-based counting approach significantly benefit IRD investigations by preserving material, improving sampling resolution, and optimizing lab workflows. By enabling faster detection of higher-frequency IRD events, these advances have significant potential to deepen our understanding of climate-ice sheet interactions on human-relevant timescales.

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