

Solar wind modeling with the Alfvén Wave Solar atmosphere Model driven by HMI-based Near-Real-Time maps by the National Solar Observatory

2 NISHTHA SACHDEVA ¹, WARD B. MANCHESTER ¹, IGOR SOKOLOV ¹, ZHENGUANG HUANG ¹,
3 ALEXANDER PEVTSOV ², LUCA BERTELLO ², ALEXEI A. PEVTSOV ², GABOR TOTH ² AND BART VAN DER HOLST ¹

5 ¹*Department of Climate and Space Sciences and Engineering, University of Michigan, Ann Arbor, MI 48109, USA*

6 ²*National Solar Observatory, 3665 Discovery Drive, 3rd Floor Boulder, CO 80303 USA*

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8 ABSTRACT

9 We explore model performance for the Alfvén Wave Solar atmosphere Model (AWSoM) with near-
10 real-time (NRT) synoptic maps of the photospheric vector magnetic field. These maps, produced by
11 assimilating data from the Helioseismic Magnetic Imager (HMI) onboard the Solar Dynamics Obser-
12 vatory (SDO), use a different method developed at the National Solar Observatory (NSO) to provide
13 a near contemporaneous source of data to drive numerical models. Here, we apply these NSO-HMI-
14 NRT maps to simulate three Carrington rotations (CRs): 2107 -2108 (centered on 2011/03/07 20:12
15 CME event), 2123 (integer CR) and 2218–2219 (centered on 2019/07/2 solar eclipse), which together
16 cover a wide range of activity level for solar cycle 24. We show simulation results, which reproduce
17 both extreme ultraviolet emission (EUV) from the low corona while simultaneously matching *in situ*
18 observations at 1 au as well as quantify the total unsigned open magnetic flux from these maps.

19 1. INTRODUCTION

20 Global estimates of the solar photospheric magnetic
21 field in the form of synoptic and synchronic maps are
22 the fundamental empirical data product that allow for
23 simulation and prediction of the three-dimensional (3D)
24 structure of the solar corona, solar wind, and the helio-
25 sphere (Mikic et al. 1999; Roussev et al. 2003; Usmanov
26 & Goldstein 2003; Cohen et al. 2007; van der Holst et al.
27 2010; Lionello et al. 2013; Sokolov et al. 2013; van der
28 Holst et al. 2014; Riley et al. 2014; Feng et al. 2014, 2015;
29 Riley et al. 2019; van der Holst et al. 2019). These maps
30 of the photospheric magnetic field are constructed from
31 a time-series of full-disk magnetograms (collected over
32 a solar rotation period of 27 days or more), which are
33 then modified and assembled to simultaneously cover the
34 entire solar surface. Photospheric full-surface maps be-
35 came available shortly following the routine production
36 of full-disk magnetograms, beginning with the Global
37 Oscillation Network Group (GONG) (see, e.g. Donald-
38 son Hanna & Harvey 2002). The Stanford approach for

39 producing synoptic maps from the Helioseismic Mag-
40 netic Imager (HMI) data is to use only 2 degree strips
41 of data at the solar central meridian from each full disk
42 magnetogram and stitch these strips together to form
43 synoptic maps.

44 Perhaps the most advanced system for producing
45 global photospheric maps is the Air Force Data Assim-
46 ilative Photospheric flux Transport (ADAPT) model,
47 which is a flux transport model that makes use of data
48 assimilation for incorporating magnetic field data. In
49 ADAPT, the photospheric magnetic flux is transported
50 by differential rotation, meridional flows and convection-
51 driven diffusion while observational data-driven updates
52 to the model are made using data assimilation tech-
53 niques (Arge & Pizzo 2000; Arge et al. 2013). ADAPT
54 maps are routinely used in numerical simulations, in-
55 cluding our model validation work (Sachdeva et al. 2019,
56 2021) and simulations of Parker Solar Probe encounters
57 (van der Holst et al. 2019, 2022) using AWSoM. Cur-
58 rently NSO provides ADAPT maps driven with GONG
59 magnetograms (<https://gong.nso.edu/adapt/maps/>).

60 While the use of HMI and GONG informed ADAPT
61 maps has been extremely successful, these data products
62 are not suitable for use in near-real-time (NRT) simula-
63 tions as a result of the significant time delay in produc-

ing the maps. For space weather forecasting, accurate maps with minimum delay from the moment the magnetic fields are observed are required. For this purpose, the National Solar Observatory embarked on a mission to produce NRT synoptic maps specifically designed as input for numerical models to forecast the coronal space environment. The synoptic map data products are available via [doi: 10.25668/nw0t-b078](https://doi.org/10.25668/nw0t-b078).

The NSO approach is to speed map creation by using the full disk vector magnetogram, and weight pixel contribution based on its distance from the central meridian (see, [Bertello et al. \(2014\)](#)). The maps (hereafter referred to as NSO-HMI-NRT) are a product of this NSO approach applied to HMI full-disk magnetograms. SOLIS/VSM vector data may also be used for NRT maps. While HMI and SOLIS/VSM produce different results for weaker fields and sometimes show the opposite orientation in transverse fields ([Pevtsov et al. 2021b](#); [Liu et al. 2022](#)), the two instruments agree very well in strong field regions ([Pietarila et al. 2013](#); [Riley et al. 2014](#)). The disagreement in weak magnetic field regions is not the result of disambiguation, but mostly due to differences in noise levels and magnetic fill factor (fraction of magnetized and non-magnetized plasma contribution to a single pixel) ([Pevtsov et al. 2021b](#)).

In this work, we explore the performance of AWSoM driven with NSO-HMI-NRT maps. For this goal, we choose three Carrington rotations (CRs): 2107–2108 (centered on 2011/03/07 20:12 UT, a CME event), 2123 (integer CR) and 2218–2219 (centered on 2019/07/2, total solar eclipse), which cover the ascending phase, solar maximum and solar minimum of the solar cycle. For simplicity, hereafter we refer to these synoptic maps using their nearest integer rotation number (i.e., CR2107, CR2123, and CR2219) although two of them straddle more than one Carrington rotation. We then make direct comparisons to observed data to provide a measure of model fidelity, first for coronal images made in the extreme ultra violet, and second with *in situ* time series data extracted near Earth. This two-type data comparison is less thorough than previous model validation efforts ([Cohen et al. 2007](#); [Jin et al. 2012](#); [Sachdeva et al. 2019, 2021](#)), but will serve the purpose of demonstrating model performance with the maps designed for space weather forecasting. In Sections 2 and 3, we briefly describe the AWSoM model and the NSO-HMI-NRT maps while Section 4 describes the simulation design. Sections 5 and 6 describe simulation results and summarize this work.

2. ALFVÉN WAVE SOLAR ATMOSPHERE MODEL (AWSOM)

AWSoM ([van der Holst et al. 2014](#); [Sokolov et al. 2013](#)) within the Space Weather Modeling Framework (SWMF; [Tóth et al. \(2012\)](#)) is a self-consistent, 3D global magnetohydrodynamic (MHD) model with its inner boundary at the base of the transition region (upper chromosphere) extending into the solar corona and the heliosphere. It is driven by the radial component of the photospheric magnetic field at the inner boundary. Like most solar corona models, this input comes from the solar synoptic/synchronic magnetic field maps, which is essential for reliable predictions. AWSoM incorporates the low-frequency Alfvén wave turbulence as a consequence of the non-linear interaction of forward and counter propagating Alfvén waves, which is based on well established theories describing the evolution and transport of Alfvén turbulence, (e.g., [Hollweg \(1986\)](#); [Matthaeus et al. \(1999\)](#); [Zank \(2014\)](#); [Zank et al. \(2017\)](#)). The AWSoM phenomenological approach self-consistently describes the heating and acceleration of the solar wind in response to turbulence while not yet including many higher-order physical effects. Several other extended MHD coronal models have been developed ([Usmanov et al. 2000](#); [Suzuki & Inutsuka 2005](#); [Lionello et al. 2014](#)), which also include Alfvén wave turbulence. AWSoM is distinguished from other global MHD models by including proton temperature anisotropy (perpendicular and parallel ion temperature), isotropic electron temperature, heat conduction and radiative cooling. The wave dissipation heats the solar wind plasma and the (thermal and nonthermal) pressure gradients accelerate the solar wind ([Meng et al. 2015](#)). The full set of MHD equations using the Block Adaptive Tree Solarwind-Roe-Upwind Scheme (BATS-R-US; [Powell et al. \(1999\)](#)) numerical scheme are solved within AWSoM. A detailed description of the model equations and their implementation is available in [van der Holst et al. \(2014\)](#). The energy partitioning scheme in AWSoM has been significantly improved and recently validated against the data from Parker Solar Probe ([van der Holst et al. 2019, 2022](#)). These improvements include using the critical balance formulation of [Lithwick et al. \(2007\)](#) and implementation of the alignment angle between the counter-propagating Alfvén waves in the energy cascade.

AWSoM has been meticulously validated by comparing the simulated results with a variety of observations. Near the Sun, the modeled density and temperature structure of the solar corona is compared to extreme ultraviolet (EUV) observations from *Solar Terrestrial Relations Observatory* (STEREO, [Howard et al. \(2008\)](#)), *Solar Dynamics Observatory* (SDO, [Pesnell et al. \(2012\)](#))/*Atmospheric Imaging Assembly* (AIA, [Lemen et al. \(2012\)](#)) and *Solar and Heliospheric Ob-*

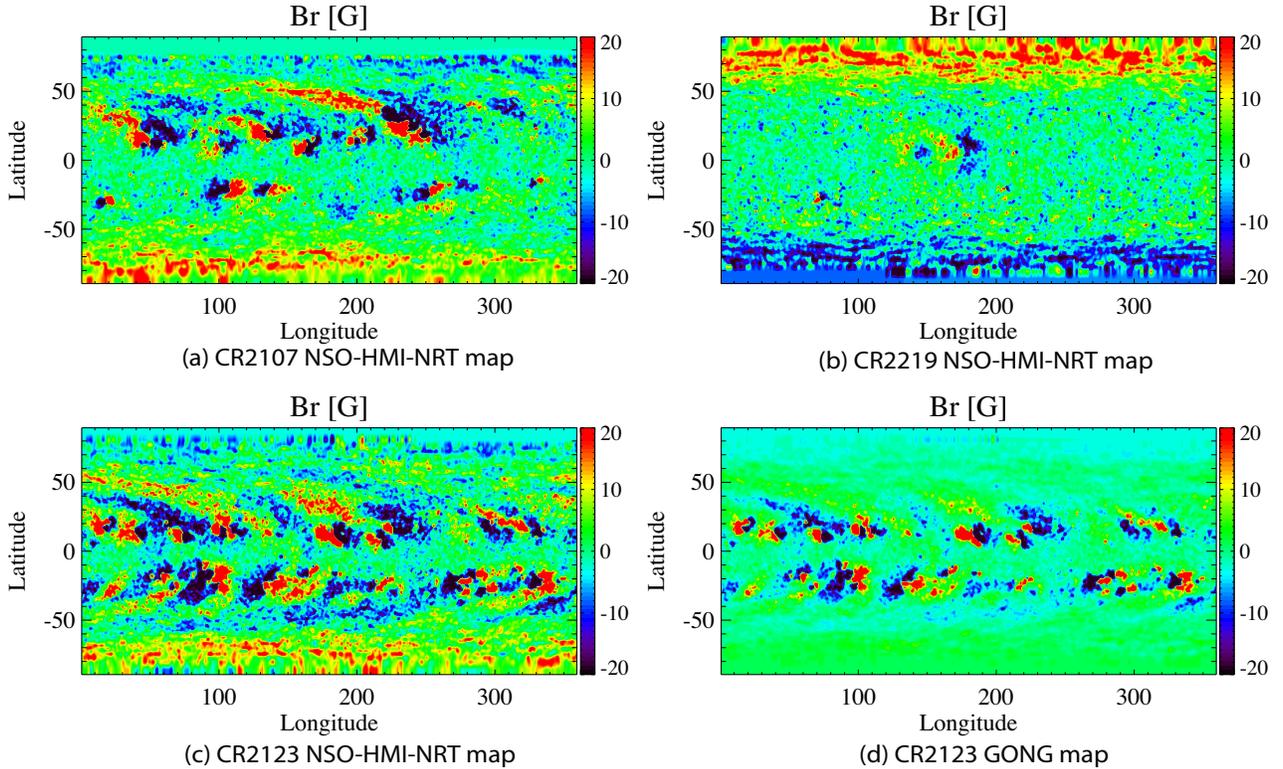


Figure 1. NSO-HMI-NRT and GONG synoptic maps showing the observed radial photospheric magnetic field. The B_r component from NSO-HMI-NRT maps are shown for Carrington Rotations 2107, 2219 and CR2123 in panels a, b, and c respectively. Panel d shows the B_r field from GONG synoptic map for CR2123. The B_r field range of ± 20 G is chosen to highlight the features on the map.

168 *servatory*(SOHO)/*Large Angle and Spectrometric Coro-*
 169 *nagraph* (LASCO, Brueckner et al. (1995)). In the
 170 low corona, AWSoM results have been compared with
 171 the tomographic reconstructions of electron density and
 172 temperature using EUV and visible-light observations
 173 (Lloveras et al. 2017, 2020, 2022; Vázquez et al. 2022). In
 174 the inner heliosphere, AWSoM successfully reproduces
 175 the velocity observations of InterPlanetary Scintillation
 176 (IPS) data (Jackson et al. 1998) and the solar wind
 177 plasma parameters at 1 au (WIND observations) (Jin
 178 et al. 2017; Sachdeva et al. 2019)). AWSoM has been
 179 successful in simulating observed solar wind properties
 180 during both solar minimum and maximum conditions
 181 (Sachdeva et al. 2019, 2021).

182 3. NSO-HMI-NRT MAPS

183 Here, we discuss the methodology of creating the
 184 NSO-HMI-NRT maps used in this paper for prescribing
 185 the magnetic field at the AWSoM inner boundary. Syn-
 186 optic maps are constructed over a full solar rotation by

187 adding new observations of the solar disk as they rotate
 188 into the observer’s view. Assuming that the Sun rotates
 189 as a solid body, the synoptic maps cover the entire solar
 190 surface over a whole Carrington Rotation (~ 27.27
 191 days) (for a general description see, e.g. Pevtsov et al.
 192 2021a, Section 7). Here we limit the discussion of charts
 193 representing full vector magnetic field, B_r (radial, or
 194 up-down), B_φ (zonal, East-West), and B_θ (meridional,
 195 North-South).

196 Various factors contribute to uncertainties in the syn-
 197 optic charts. Posing the largest challenge are the lim-
 198 ited contemporaneous observations of the solar surface,
 199 particularly in the polar regions. Instrumental noise,
 200 conditions of observations and similar factors are also
 201 inherited in the maps. Smearing of solar features due
 202 to differential rotation is also a potential issue, par-
 203 ticularly when creating high-spatial resolution synoptic
 204 maps. Because the differential component of solar ro-
 205 tation increases with latitude, the smearing effect will
 206 mostly be prominent within the polar regions (above 60

degrees). A detailed description of this problem and a possible solution is given in Ulrich & Boyden (2006). However, due to the low resolution of the synoptic maps used in this study, this correction is not required.

As a first step in the creation of a synoptic map, full disk images are remapped from sky-plane coordinates to heliographic coordinates. If remapping is required, the image resolution is reduced to match the resolution of the synoptic map (e.g., 1 by 1 degree in solar latitude and longitude). Next, these sampled or remapped images are added to a synoptic map based on heliographic coordinates of pixels. One method used by SDO/HMI team employs vertical strips of about 2 solar degrees wide centered at the solar central meridian. It produces the so-called diachronic maps. This method is quick and easy since it does not require excessive additional processing. However, it requires sufficiently high cadence in observations as the strips are simply added one after the other, similar to a picket fence. Therefore, any gap in observations results in a gap in the diachronic map. Furthermore, such a map fails to correctly represent any features that emerge or drastically evolve after passing the central meridian.

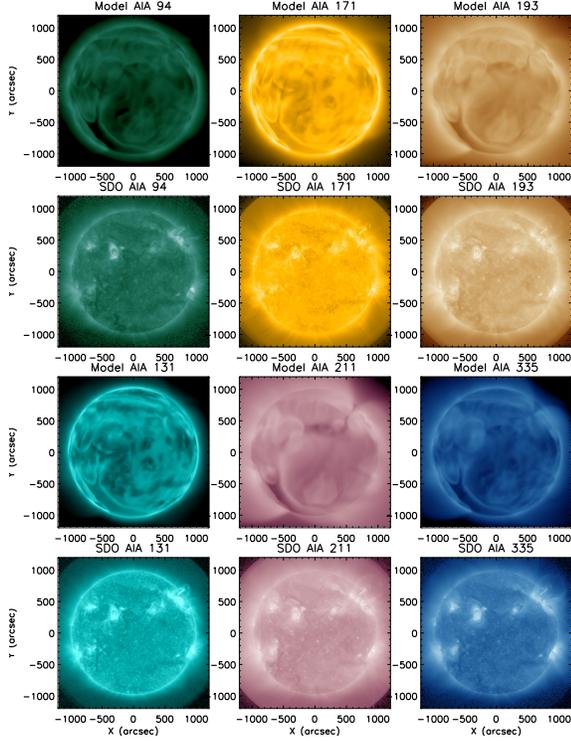
The NSO-HMI-NRT synoptic maps used in this paper are generated using a different approach which incorporates the use of the full disk SDO/HMI magnetograms to build a synoptic map (Bertello et al. 2014). This method can be computationally expensive if the cadence of full-disk magnetograms is too high. However, the maps created with this technique include all the magnetic features regardless of when they appear during a rotation or whether they evolve significantly before and/or after passing the central meridian. Each NSO-HMI-NRT magnetic field map in this paper incorporates approximately 43 (8+27+8) days of observations. That is, in addition to the 27 days of a Carrington rotation, the synoptic maps cover eight days (each) before and after the rotation. While adding 8 days before/after the start/end of a Carrington rotation is not necessary, it allows for a better equalisation of weight (or number of contributing full disk observations) for each heliographic pixel. Without this contribution, the first (last) 8 days of a Carrington rotation map will see a gradual increase (decrease) in a normalized number of contributing points from just a few percent at the leading (trailing) edge to 100% in the center of the map. Because of this difference in weights, without additional 8 days, the noise level would be slightly higher for the beginning and end parts of each map. A similar procedure is adopted in creating the synoptic maps of pseudo-radial field using GONG observations. Our past experience with HMI vector observations have shown that the observation on

the 48th minute of every hour provides the best coverage and data quality for the one-hour cadence that is used here. The selection of the 48th minute is not critical, and has no impact on the results of our project. Nevertheless, it may yield the synoptic maps of a slightly better quality.

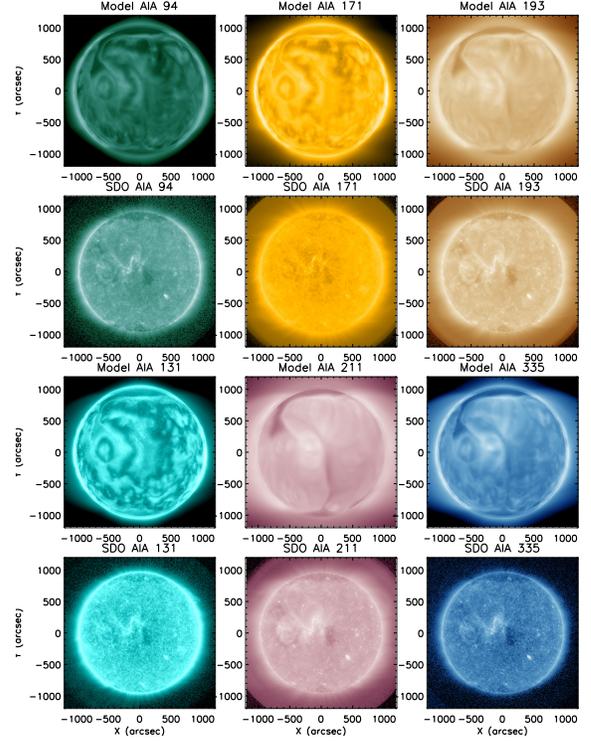
The SDO/HMI data is acquired from the Joint Science Operations Center (JSOC, <http://jsoc.stanford.edu/>). This process begins by using a custom python program utilizing an http ‘get request’ to the JSOC server to query the data that is available for download at that given time. This request is used to verify and record the availability status of all five Data Record Management System (DRMS) segments per observation that can be used to build a given synoptic map. Within JSOC these segments are identified as field, inclination, azimuth, disambig, and conf_disambig (map of the confidence in each pixels disambiguation) for the Full-Disk Milne-Eddington inversion data series (hmi.B.720s), each covering 720 s of observation. A one-hour cadence or 5 segments per hour are used to generate a synoptic map. If all five DRMS segments are available, they are separated into 5 lots of up to 9 days of data or $9 \times 24 \times 5 = 1080$ segments each. However, full block availability is rare and usually a few gaps of missing data occur every few days. Even with these gaps, this sums up to about 70 GB of observational data (before processing). This large amount of data requires parallelized workloads of each lot to reduce the computational time needed for the next steps. The data acquisition is followed by pre-processing of the SDO/HMI images of the photospheric magnetic field into a single coordinate system transformed package.

Pre-processing includes the coordinate transformation from the image (sky) plane to heliographic (solar latitude-longitude) coordinates, re-imaging the full disk data to larger pixels used for construction of a synoptic map, and applying \cos^4 of central meridian distance weighting function. For additional details see Bertello et al. (2014); Hughes et al. (2016). After the pre-processing, the data are used to assemble a complete synoptic map by averaging the contribution to a corresponding synoptic map pixel from all contributing pre-processed images.

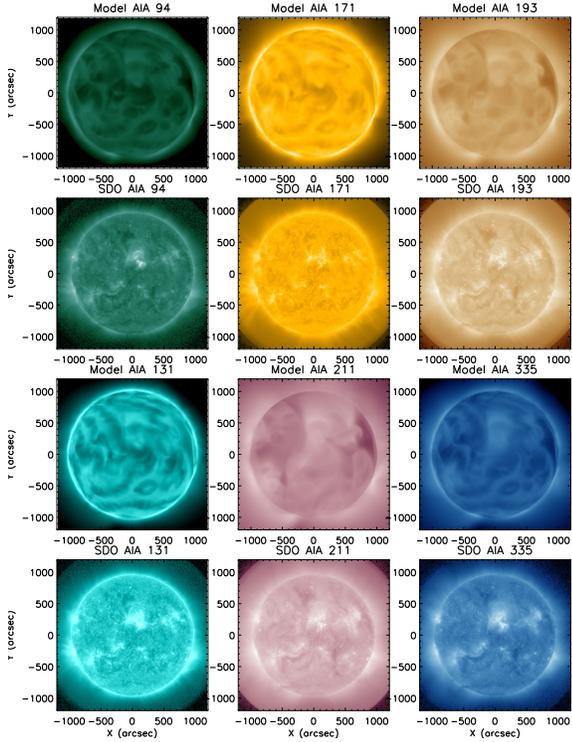
A well-known problem in constructing a full-surface synoptic map is the limited visibility of the polar fields of the Sun from near the Earth. The tilt angle between the Earth’s ecliptic plane and the solar rotation axis varies between about $\pm 7.25^\circ$ each year. The poles can therefore only be observed from the ecliptic plane with a large ($> 80^\circ$) viewing angle. Moreover, each pole is not observable from near-Earth for more than six months



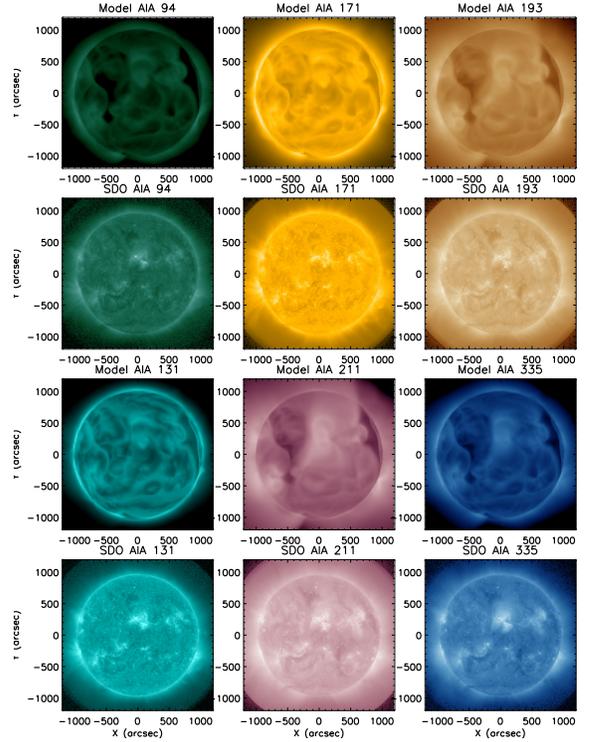
(a) CR2107 with NSO-HMI-NRT map



(b) CR2219 with NSO-HMI-NRT map



(c) CR2123 with NSO-HMI-NRT map



(d) CR2123 with GONG map

Figure 2. Comparison of synthetic EUV images with SDO/AIA observations. Panels a, b and c show model-data comparison for CR2107, CR2219 and CR2123 modeled using the NSO-HMI-NRT maps. Panel d shows the same for CR2123 modeled using GONG magnetogram. In each panel, the first and third rows represent the modeled AIA output and the second and fourth rows show the SDO/AIA observations. The comparison is shown in six wavelength channels (94, 171, 193, 131, 211 and 335 Å).

in a year. The unobserved polar fields are therefore required to be modeled. A simple approach, adopted here, is to fill the pixels corresponding to the unobserved polar data using a cubic surface fit to the observational data from the neighboring latitudes. Numerous studies suggest that the polar fields are approximately radially directed (Svalgaard et al. 1978; Petrie 2015). Ulrich & Tran (2013) argued for a slight $\approx 6^\circ$ poleward inclination of magnetic field in polar areas. Pevtsov et al. (2021a) applied a similar technique, and found about 3° equatorward inclination when using SOLIS/VSM data. Virtanen et al. (2019) also found a small ($\leq 10^\circ$) equatorward inclination at high ($\sim 75^\circ$) latitudes. Based on the results of these previous studies, we set the unobserved B_θ and B_ϕ values in the polar regions to zero and fill B_r using cubic surface fit. Finally, this process provides a complete synoptic map.

4. NUMERICAL SIMULATION SET-UP

We use the NSO-HMI-NRT magnetic field maps for three Carrington rotations (CRs) to drive the background solar wind simulations. The solar corona (SC) and inner heliosphere (IH) components of the SWMF are used via AWSoM (Section 2). The NSO-HMI-NRT synoptic maps used this in study for Carrington rotations CR2107 (2011-02-16 to 2011-03-16), CR2123 (2012-04-28 to 2012-05-25), and CR2219 (2019-06-29 to 2019-07-26) are shown in panels a, b and c of Figure 1. Each of these maps show the structure of the photospheric magnetic field obtained from the SDO/HMI images followed by the procedure described in the previous Section 3. In a previous work, Jin et al. (2017) used the *Global Oscillation Network Group* (GONG) synoptic magnetogram to simulate the solar wind conditions for CR2107 using AWSoM and compared the simulation results with OMNI data at 1 au and along the trajectory of STEREO-A. van der Holst et al. (2014); Meng et al. (2015) showed an improvement in the results for this rotation when they use the magnetic field obtained from SDO/HMI instrument (Scherrer et al. 2012).

The 2D photospheric magnetic field from the synoptic maps is used to reconstruct the 3D magnetic field using the Potential Field Source Surface Model (PFSSM). The radial component of the observed magnetic field is used as the boundary condition for the potential solver while the longitudinal and latitudinal components are allowed to relax to a solution. We use the spherical harmonics solution for the PFSSM with the source surface at $2.5 R_\odot$. At the inner boundary, the initial temperature for both isotropic electron and perpendicular and parallel proton temperature is set to 50,000 K. The proton number density at these temperatures

is overestimated to provide a ready source to replenish the plasma, which maybe depleted due to chromospheric evaporation (Lionello et al. 2009; van der Holst et al. 2014). AWSoM has a limited number of free parameters that may be varied to improve the results when compared to observations of the solar wind. The energy density of the outward propagating Alfvén waves is set using the Poynting Flux (S_A) of the wave which is proportional to the magnetic field strength at the inner boundary (B_\odot) (Fisk 2001; Fisk et al. 1999; Sokolov et al. 2013). A recent study by Huang et al. (2022) using AWSoM showed that the quantity $(S_A/B)_\odot$ needs to be varied based on the phase of the solar cycle. During phases of stronger magnetic activity, the amount of energy of the outward propagating Alfvén wave is reduced by reducing the $(S_A/B)_\odot$ parameter to avoid deposition of excess energy density into the chromosphere and high density peaks at 1 au (Sachdeva et al. 2021). For CR2107, CR2123 and CR2219 the optimal value of the $(S_A/B)_\odot$ parameter in the model is set to 0.35, 0.3 and 1.0 in units of $10^6 \text{ Wm}^{-2}\text{T}^{-1}$, respectively. The Alfvén wave correlation length (L_\perp), which is transverse to the magnetic field direction is proportional to $B^{-1/2}$ (Hollweg 1986) and is set to $1.5 \times 10^5 \text{ m} \sqrt{T}$.

The SC component uses a 3D spherical grid extending from 1 - $24 R_\odot$ and is coupled with the IH component which uses a Cartesian grid that extends from -250 to $250 R_\odot$. The SC and IH components are coupled with a buffer grid extending from 18-21 R_\odot to transfer the SC solution to the IH domain in the steady-state run. The SC domain is decomposed into $6 \times 8 \times 8$ grid blocks while IH has $8 \times 8 \times 8$ grid blocks. The computation includes Adaptive Mesh Refinement (AMR) in SC which provides an angular resolution of 1.4° below $1.7 R_\odot$ and 2.8° in the remaining domain. The number of cells in SC and IH are 4.2 million and 12.2 million respectively. The cell size in IH ranges between $0.48 R_\odot$ and $7.8 R_\odot$. Using local time stepping, the SC component is run for 80000 iterations and coupled with IH for one step followed by 5000 steps in IH to get the steady state solution. Additional AMR is done below $1.7 R_\odot$ along with the 5th order shock-capturing scheme (Chen et al. 2016) to produce high resolution line of sight synthetic EUV images for comparison with observations.

5. RESULTS

We use the magnetic field from the NSO-HMI-NRT maps to obtain steady-state solar wind solutions for three CRs 2107, 2123 and 2219. Figure 1 shows the NSO-HMI-NRT synoptic maps depicting detailed features of the active regions as well as the polar regions as described in Section 3. These maps represent dif-

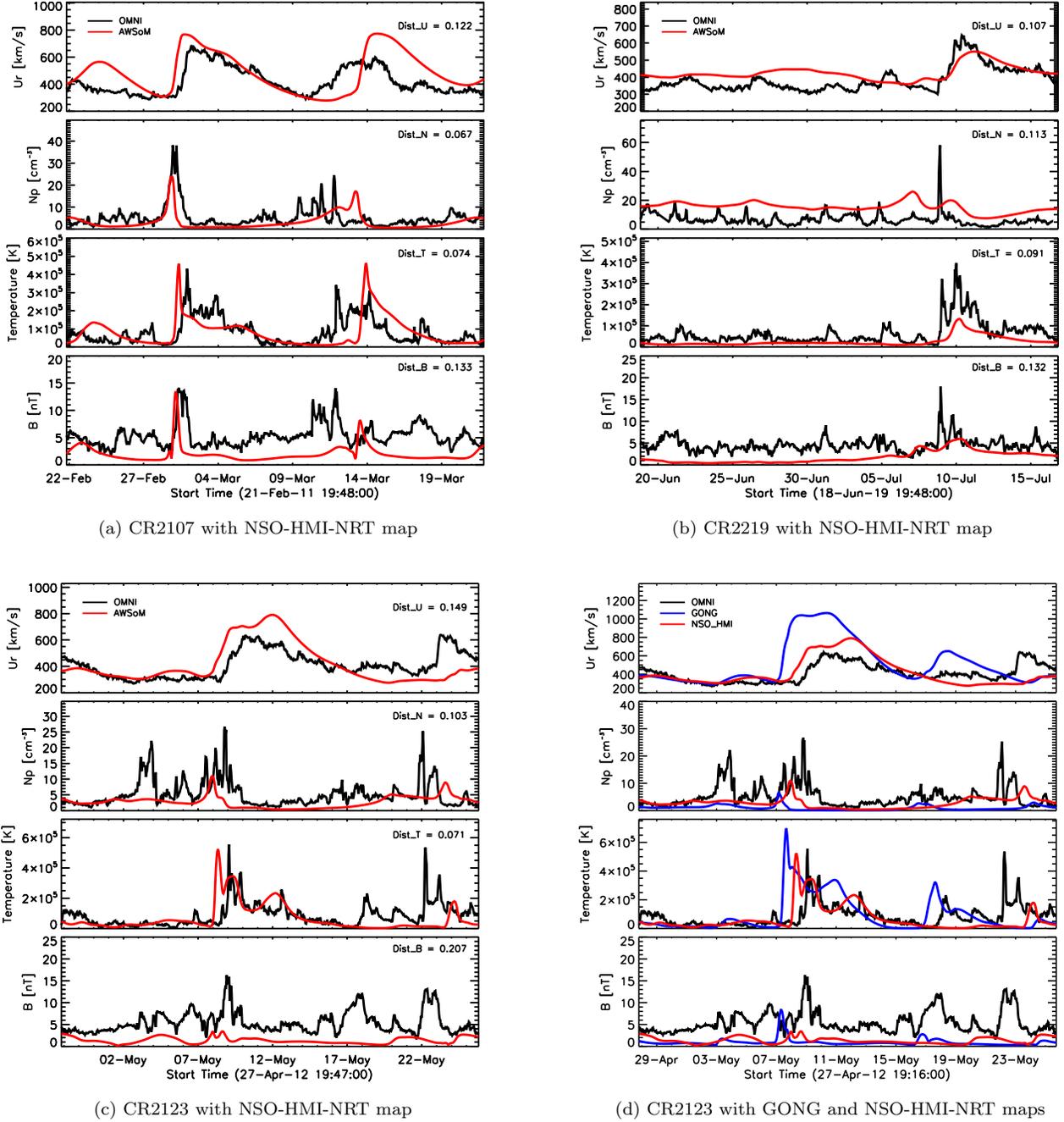


Figure 3. Comparison of AWSoM simulated 1 au solar wind plasma parameters with the 1-hr averaged OMNI observations for the three CR's. Model results are in red and data is in black. Panels a, b and c correspond to data-model comparisons for CR2107, CR2219 and CR2123 respectively where the simulations were driven by the NSO-HMI-NRT maps (shown in Figure 1). Panel d shows the comparison of simulation results from AWSoM driven by the NSO-HMI-NRT (red) and GONG (blue) maps for CR2123.

413 ferent phases of the solar cycle. CR2107 and CR2123
 414 correspond to higher solar activity with stronger mag-
 415 netic field and more active regions in comparison to
 416 solar minimum conditions found in CR2219. Panels
 417 c and d of Figure 1 show the comparison between
 418 the magnetic field maps from NSO-HMI-NRT (left)

419 and GONG (right) for CR2123. More small-scale fea-
 420 tures are present in the NSO-HMI-NRT map as well as
 421 stronger magnetic fields in the active regions. In the
 422 GONG map, the polar magnetic fields are weaker and
 423 smoother, a distinct difference which will impact the
 424 speed of the modeled solar wind.

425 The simulation domain for AWSoM covers the low
 426 corona making it ideal to obtain synthetic extreme ul-
 427 traviolet (EUV) images that can be compared with cor-
 428 responding observations. Figure 2 shows the synthetic
 429 line-of-sight images from the AWSoM simulation results
 430 compared with corresponding SDO/AIA observations in
 431 six wavelength channels (94, 171, 193, 131, 211 and
 432 335 Å). Here, panels a, b and c show model-data com-
 433 parisons for CR2107, CR2219 and CR2123 respectively,
 434 modeled using the NSO-HMI-NRT maps. The first and
 435 third row of images in each panel show the model syn-
 436 thesized AIA images and the second and fourth rows
 437 of images show the corresponding SDO/AIA observa-
 438 tions. Panel d of Figure 2 shows the same model-data
 439 comparison for CR2123 modeled using the GONG mag-
 440 netogram. The simulation results compares well with
 441 observations in matching the overall brightness and the
 442 location of the major active regions for the three rota-
 443 tions. This fact suggests that AWSoM can reproduce
 444 the 3D structure of the density and temperature in the
 445 low solar corona. The modeled coronal holes appear to
 446 be darker in comparison to observations but match in
 447 their location and extent. Although the average bright-
 448 ness matches well in all channels, the bright active
 449 regions can best be seen accurately in 193, 211 and 335
 450 Å channels. For example, for CR2219 (Panel b), dur-
 451 ing solar minimum, the model AIA images reproduce
 452 the major bright active region which can be seen clearly
 453 in these wavelength channels. For CR2123, synthetic
 454 AIA images obtained from driving the model using the
 455 NSO-HMI-NRT and GONG maps show major differ-
 456 ences in the overall brightness of the active regions and
 457 the coronal holes. In particular, the modeled coronal
 458 hole from the GONG-driven simulation appears to be
 459 much darker. Additional refinement of the grid with the
 460 AWSoM model can further improve model comparisons
 462 by producing brighter active regions (Shi et al. 2022).

463 To compare the simulated solar wind with *in situ* ob-
 464 servations of plasma parameters at L1 we extract (from
 465 the 3D result) the model solution along the trajectory
 466 of the Earth. Figure 3 shows the AWSoM output along
 467 the Earth’s trajectory in red color and the OMNI data
 468 in black for all three rotations. We see that overall,
 469 the model when driven by NSO-HMI-NRT maps suc-
 470 cessfully reproduces the observed solar wind plasma. In
 471 particular, we see that for both CR2107 and CR2123 the
 472 solar wind solution matches quite well with the observa-
 473 tions for all quantities. For CR2107, the model predicts
 474 the co-rotating interaction region (CIR) on March 1,
 475 2011. The solution matches the significant jump in the
 476 radial speed (U_r), proton density (N_p), ion temperature
 477 and the absolute magnetic field (B). For both rotations

478 that represent the near solar maximum phase (CR2107
 479 and CR2123), the features in the solar wind plasma pa-
 480 rameters are well matched by the model solution. In
 481 both case, however, we find that the magnetic field is
 482 under-predicted and the peak model speed is overesti-
 483 mated by about 11 % and 23 % for CR2107 and CR2123
 484 respectively. As a result of the higher speed, the CIR
 485 in both cases arrive slightly earlier in the model as com-
 486 pared to the observations. CR2219 is a period of re-
 487 duced activity for which the model overestimates the
 488 solar wind speed and density. However, the CIR speed
 489 in the model matches well with the observations. To
 490 further quantify the model-data comparison, we calcu-
 491 late a distance measure $Dist$ listed in each plot, which
 492 informs us of how well the model matches observations.
 493 Described in detail in Sachdeva et al. (2021), the quan-
 494 tity $Dist$ is a measure of the distance between two curves
 495 independent of the coordinates. Smaller values indicate
 496 a better fit.

497 Panel d of Figure 3 shows the OMNI observations in
 498 black and the model results for CR2123 driven by NSO-
 499 HMI-NRT and GONG maps (Panels c and d of Figure
 500 1) are shown in red and blue color, respectively. The
 501 modeled solutions differ significantly at 1 au, which is a
 502 direct result of the different initial magnetic field condi-
 503 tions from the two maps. All other model parameters
 504 are kept the same for both simulations. This demon-
 505 strates that the observational magnetic field input driv-
 506 ing the solar corona models significantly impacts the so-
 507 lar wind properties.

508 Figure 4 represents the radial magnetic field (B_r) at
 509 the source surface radius of $2.5 R_\odot$ obtained from the
 510 PFSSM using spherical harmonics with order 180 for
 511 each of the rotations. In panels a, b and c, the field
 512 B_r is obtained from the NSO-HMI-NRT maps for CRs
 513 2107, 2219 and 2123. For comparison, panel d shows
 514 the B_r field obtained from the GONG map for CR2123.
 515 Both maps for CR2123 (Panels c and d) are shown
 516 on the same scale to highlight the differences between
 517 them. The field obtained from the NSO-HMI-NRT map
 518 is much more pronounced in the coronal holes as well as
 519 the polar regions in comparison to the GONG map.

520 To quantify this effect, we also calculate the total un-
 521 signed open magnetic flux for all the maps at $2.5 R_\odot$.
 522 This quantity is an integral of the absolute value of the
 523 radial magnetic field, $|B_r|$, over the source surface. The
 524 total unsigned open magnetic flux at $2.5 R_\odot$ is found to
 525 be 10.3, 15.9, and 6.9 [Gauss R_\odot^2] for the NSO-HMI-
 526 NRT maps for CR2107, CR2219 and CR2123, respec-
 527 tively. For the GONG map for CR2123, the total un-
 528 signed magnetic flux obtained is 3.1 [Gauss R_\odot^2] at 2.5
 529 R_\odot . The scaling law by Pevtsov et al. (2003) relates

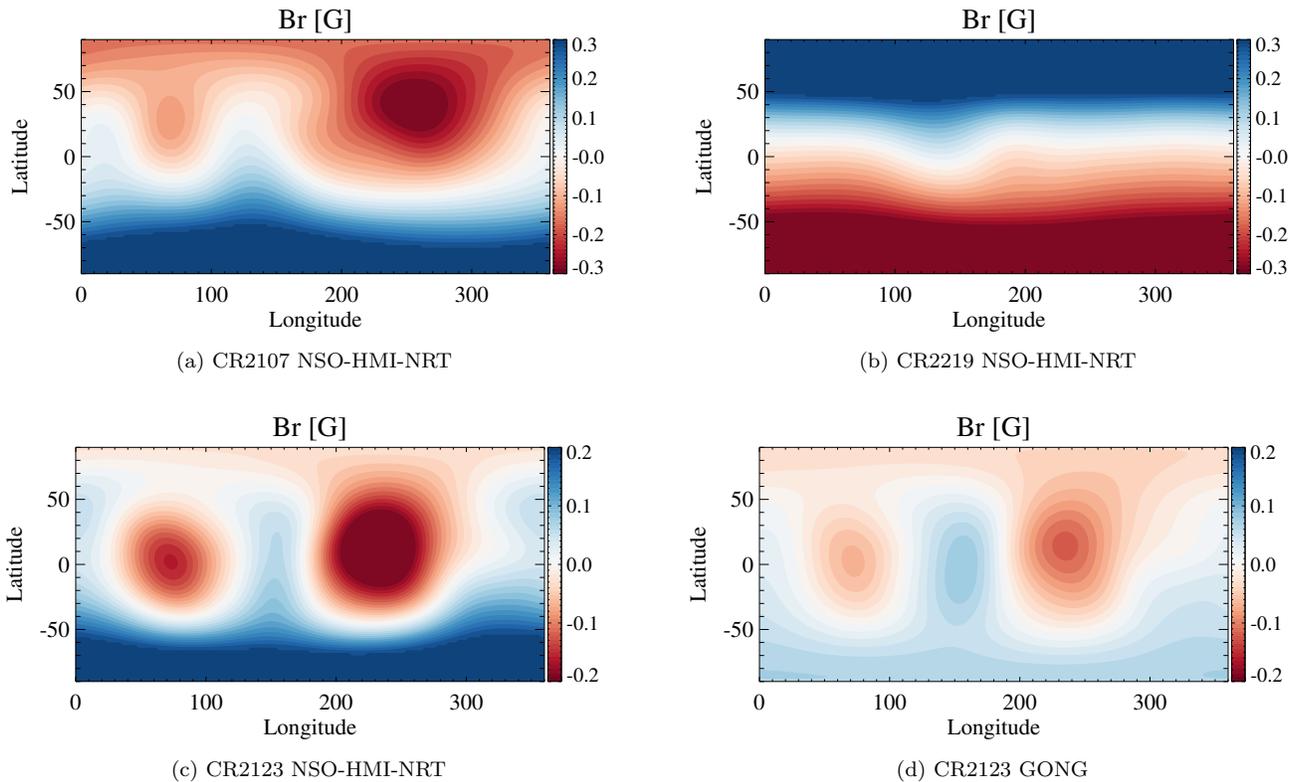


Figure 4. Radial magnetic field at the source surface radius ($2.5 R_{\odot}$). Panels a, b and c show the B_r magnetic field for CR2107, CR2219 and CR2123 respectively at $2.5 R_{\odot}$ calculated from the PFSSM using the NSO-HMI-NRT maps. Panel d shows the same for CR2123 using the GONG synoptic map. The source surface in the PFSSM is set to $2.5 R_{\odot}$ for all the maps.

530 the total unsigned flux to the energy deposition in the
 531 solar corona, therefore, a stronger total unsigned open
 532 flux leads to more energy, which accelerates and powers
 533 the solar wind. In relation to AWSoM, the Poynting flux
 534 outgoing into to the solar wind is directly proportional to
 535 the unsigned open magnetic flux and the constant ratio
 536 of the Poynting flux to the magnetic flux is one of the input
 537 parameters of the model (S_A/B_{\odot}). Therefore, the
 538 stronger open flux from the NSO-HMI-NRT map provides
 539 more energy to the corona, which increases chromospheric
 540 evaporation, increasing the density of the solar wind
 541 while reducing its speed. The result is a better
 542 comparison with observations at 1 au compared to the
 543 model results made with the GONG map for CR2123.

544 6. SUMMARY AND DISCUSSION

545 In this work, we show the impact of the magnetic field
 546 conditions obtained from the NSO approach of creating
 547 near-real-time maps using the HMI magnetic field obser-
 548 vations (NSO-HMI-NRT maps) on the modeled solar
 549 wind. The methodology used for these maps includes
 550 using full-disk HMI magnetogram but with a weighted

551 pixel contribution and the unobserved polar regions are
 552 filled using a polynomial fit to neighbouring observa-
 553 tions.

554 We use the 3D MHD model AWSoM to simulate the
 555 Sun to Earth background solar wind for three Carrington
 556 rotations (2107, 2219 and 2123). AWSoM is driven
 557 by the magnetic field from the NSO-HMI-NRT maps to
 558 demonstrate their performance during varying periods
 559 of solar activity. We compare the AWSoM simulated
 560 solar wind solutions with observations in the low corona
 561 and find that for all three CR's modeled using the corre-
 562 sponding NSO-HMI-NRT maps as input, the large-scale
 563 properties of the solar corona including the extent and
 564 location of coronal holes as well as regions of enhanced
 565 activity (active regions) compare well with SDO/AIA
 566 observations.

567 Further away from the Sun, we compare the observed
 568 solar wind properties at 1 au with the model results.
 569 The 1 au observations are reproduced reasonably well
 570 by the NSO-HMI-NRT map driven AWSoM model for
 571 all three rotations. We find that the while magnetic field
 572 is underestimated, the solar wind speed, density and

CIR properties are reproduced in the simulations. For one of the rotations (CR2123), we obtain the solar wind solution using the GONG synoptic map as the initial condition for the photospheric magnetic field and compare the results with the NSO-HMI-NRT map driven solar wind conditions. We find that the solutions with the NSO-HMI-NRT map perform well both in the low corona and at 1 au when compared to the GONG map results. Finally, we show the radial magnetic field at $2.5 R_{\odot}$ from the PFSSM for each of the maps and compare the total unsigned open magnetic flux at the source surface. This quantity for the NSO-HMI-NRT map is larger by a factor of ≈ 2 in comparison to the GONG map for the same rotation.

It is well-known that numerical models of the solar corona are sensitive to the observed magnetic field inputs obtained from a variety of synoptic magnetograms available in the community. Here, we highlight the performance of NSO produced, HMI observation based, near-real-time maps (NSO-HMI-NRT) maps with our 3D extended MHD model (AWSOM) and show that the NSO-HMI-NRT maps are a valuable data product allowing for coronal/solar wind simulations of equal or better

quality than those obtained by standard synoptic maps such as GONG maps.

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