

1 **Title:** The Resilience of Habitable Climates Around Circumbinary Stars

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32 **Key Points:**

- 33
- 34 • Circumbinary systems are common in the Universe.
  - 35 • Earth-like planets in circumbinary systems are resilient against climate catastrophe,  
36 despite significant time-dependent variations in the incident stellar flux.
  - 37 • While ocean temperatures are not appreciably affected, land surface temperatures can  
38 exhibit an additional mode of circumbinary-seasonal variability.
  - 39 • Circumbinary systems should be considered as viable hosts for habitable worlds.
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43 **Abstract:**

44       Here we use a 3-D climate system model to study the habitability of Earth-like planets  
45 orbiting in circumbinary systems. A circumbinary system is one where a planet orbits around  
46 two stars simultaneously, resulting in large and rapid changes to both the stellar energy  
47 distribution and the total stellar energy received by the planet. We find that Earth-like planets,  
48 having abundant surface liquid water, are generally effective at buffering against these time-  
49 dependent changes in the stellar irradiation due to the high thermal inertia of oceans compared  
50 with the relatively short periods of circumbinary-driven variations in the received stellar flux.  
51 Ocean surface temperatures exhibit little to no variation in time, however land surfaces can  
52 experience modest changes in temperature, thus exhibiting an additional mode of climate  
53 variability driven by the circumbinary variations. Still, meaningful oscillations in surface  
54 temperatures are only found for circumbinary system architectures featuring the largest  
55 physically possible amplitudes in the stellar flux variation. In the most extreme cases, an Earth-  
56 like planet could experience circumbinary-driven variations in the global mean land surface  
57 temperature of up to  $\sim 5$  K, and variations of local daytime maximum temperatures of up to  $\sim 12$   
58 K, on monthly timescales. Still, habitable planets in circumbinary systems are remarkably  
59 resilient against circumbinary driven climate variations.

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61 **Plain Language Summary:**

62 Planets in circumbinary systems can experience significant variations in their received stellar  
63 flux over 10 to  $\sim 100$  day timescales. However, habitable planets are effective at buffering these  
64 changes. Even for the most extreme variations in stellar flux possible for circumbinary systems,  
65 Earth-like planets avoid any climate catastrophe and remain resiliently habitable.

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## 68 1. Introduction

69 Planets that orbit two stars at once are a common trope of science fiction. These are  
70 known as circumbinary systems, or Planet (P)-type systems (Dvorak, 1984). Indeed, stellar  
71 population statistics indicate that perhaps half of Sun-like stars exist in binary and higher  
72 multiple systems, and planets in circumbinary systems have already been discovered in transit  
73 light curves (Doyle et al. 2011; Welsh et al. 2012; Kostov et al. 2014; Welsh et al. 2018; Orosz et  
74 al. 2019; Kostov et al. 2020). A number of previous works have considered the climate and  
75 habitability of planets in circumbinary systems using the results from 1-dimensional climate  
76 models. These have included both analytic calculations that weight the relative spectral  
77 contribution of two stars into the total irradiance received by planets (Kane & Hinkel 2012;  
78 Haghhighipour & Kaltenegger, 2013; Wang & Cuntz 2019a, Wang & Cuntz 2019b; Moorman et  
79 al. 2020), 1-D radiative-convective climate calculations that prognostically blend stellar spectra  
80 within their radiative transfer calculations (Cukier et al. 2019), and time-marching 1-D energy  
81 balance models that explicitly account for the time-dependent changes to the total incident stellar  
82 radiation (Forgan 2014; May & Rauscher 2016; Haqq-Misra et al. 2019). Additionally, two  
83 previous studies have explored the climate and atmospheres of planets in circumbinary systems  
84 using 3-D general circulation models (GCMs). May & Rauscher (2016) used an idealized GCM  
85 to study Kepler-47 b, assuming it to be a Neptune-like world. Popp & Eggl (2017) used a 3-D  
86 terrestrial planet climate system model to study theoretical Earth-sized and completely ocean  
87 covered planets in the Kepler-35 system. In both of these studies, it was shown that global mean  
88 temperatures are not significantly affected by periodic variations in the stellar instellation in  
89 these specific circumbinary systems. The high thermal inertia of both a thick Neptunian  
90 atmosphere, and of a completely ocean covered planet surface, are each highly effective at  
91 buffering time-dependent oscillations in the total stellar irradiance, resulting in variations in

92 global mean temperatures of  $\sim 1.0\%$  or less in these cases. However, both the Kepler-47 and  
93 Kepler-35 systems feature binary periods and amplitudes of stellar flux variation that are  
94 relatively small compared to what is physically possible.

95 To drive more meaningful variations in circumbinary planetary climate, the magnitude  
96 and period of stellar irradiance changes experienced by the planet must be maximized, and the  
97 thermal inertia of the surface/atmosphere system must be minimized (Haqq-Misra et al. 2019).  
98 The circumbinary periods used in May & Rauscher (2016) and Popp & Eggl (2017) are 7.448  
99 and 22 Earth days respectively, and the normalized amplitude of variation in the incident stellar  
100 flux due to the circumbinary architecture were  $\sim 20\%$  and  $\sim 15\%$  respectively. Here, we consider  
101 the normalized amplitude of the stellar flux variation as the maximum received flux minus the  
102 minimum received flux, divided by the time-averaged flux received by the planet over many  
103 orbits. Note, that eclipsing binaries can cause drops in stellar irradiance by up to nearly  $\sim 50\%$   
104 during the eclipse if the stars are of equal radius and on a short period; however, eclipses last  
105 only a few hours and thus their time integrated effect on the total stellar irradiance is actually  
106 quite small. Thus, eclipses are not generally important for driving changes to planetary climate  
107 nor weather. The scenarios studied by May & Rauscher (2016) and Popp & Eggl (2017), while  
108 based on known systems, represent circumbinary system architectures where the variation in  
109 stellar forcing is relatively modest compared to what is physically possible and dynamically  
110 stable. Using a simplified orbital calculation, that assumed circular orbits for all bodies, Haqq-  
111 Misra et al. (2019) found that for a planet receiving a time-averaged stellar flux equal to that of  
112 the modern Earth ( $1360 \text{ Wm}^{-2}$ ) in a maximal scenario would experience a change in stellar flux  
113 of  $\sim 55\%$  over a synodic period of  $\sim 150$  Earth days. Note that this result is based on the orbital  
114 dynamical stability criteria of Holman & Wiegert (1999), who argued that such planets can

115 remain dynamically stable in circumbinary systems so long as their orbital period is at least  $\sim 3$   
116 times larger than the circumbinary orbital period (see also Quarles et al. 2018). The latitudinal  
117 energy balance model results of Haqq-Misra et al. (2019) also showed that land surface  
118 characteristics and planetary obliquity can play key roles in the climate response to circumbinary  
119 forcings. Land surfaces heat up and cool down much more rapidly than does a deep optically  
120 thick atmosphere or an ocean surface, a consideration that was overlooked in previous GCM  
121 works. The coupling of obliquity-driven seasons like on Earth, with circumbinary-driven  
122 seasons due to the variation in received stellar flux, has the potential to generate more extreme  
123 local weather conditions when the warm seasons coincide.

124         While the global mean temperature is often the first consideration in whether or not a  
125 planet may be deemed habitable, seasonal and local extremes could pose challenges for the  
126 evolution and maintenance of surface life as we know it. Here, we revisit the problem of climate  
127 and habitability for terrestrial planets orbiting in circumbinary systems by using a GCM  
128 developed for Earth-like planets, while assuming modern Earth continental configurations,  
129 atmospheric composition, rotation rate, and time-averaged received stellar insolation ( $1360 \text{ Wm}^{-2}$ ).  
130 First, we examine the full-phase space of potential circumbinary orbital configurations using  
131 the analytic orbital calculations of Georgakarakos and Eggl (2015) to identify system  
132 architectures that cause the most extreme variations in incident stellar flux received by the  
133 planet. Then we assess the impact on global, regional, and seasonal climates of these Earth-like  
134 planets in circumbinary systems using a 3-D climate system model. Ultimately, our results show  
135 that for physically plausible systems, land surface temperatures can exhibit circumbinary-driven  
136 seasonality, but global climate remains resilient against climatic catastrophe and habitability  
137 remains possible.

## 138 2. Methods

### 139 2.1. Description of the Circumbinary Treatment

140 To represent the variation in stellar irradiance received by planets in circumbinary  
141 systems, we use the analytical solutions for coplanar planets in circumbinary systems of  
142 Georgakarakos & Eggl (2015). This method was also used to describe the Kepler-35 system in  
143 the 3-D climate simulations of Popp & Eggl (2017). We use the Georgakarakos & Eggl (2015)  
144 analytical circumbinary calculations first independently to explore what circumbinary system  
145 architectures are possible (section 3), and then fully integrated into our 3-D climate system  
146 model to explore the climatic effects for specific cases (section 4). Note we have made some  
147 minor modifications to the interfacing routine in order to suit our particular operational needs;  
148 however, the circumbinary system solutions are unaltered. As inputs we define the masses of the  
149 circumbinary stars and planet, the separation of the circumbinary stars, and the desired time-  
150 mean incident stellar flux received by the planet ( $1360 \text{ Wm}^{-2}$ ). The modified code then solves  
151 for the planetary orbit that meets these conditions, if such an orbit is dynamically stable. Cases  
152 with unstable orbital dynamics are rejected. The Georgakarakos & Eggl (2015) model solves for  
153 the planetary orbit, and provides the incident stellar flux received by the planet from each star  
154 individually as a function of time, which includes the effect of eclipsing binaries, dual evolving  
155 zenith angles, and planetary eccentricity. The planet is assumed to start out on a circular orbit  
156 which evolves over time to an equilibrium state. The resulting short periodic and secular  
157 changes in the planetary eccentricity are tracked. Whether planets form on circular orbits or  
158 orbits with forced eccentricity is still a matter of debate (Mardling 2007; Silsbee & Rafikov  
159 2014). Circumbinary planets on forced orbits tend to experience less variation in insolation  
160 (Eggl, 2018). In this work we have chosen initially circular orbits, as we aim to explore how

161 planetary climates react to extreme variability. To describe the stars, we use mass, luminosity,  
 162 radius, and effective temperature relationships derived from Pecaut et al. (2012) and Pecaut &  
 163 Mamajek (2012). The choice of mass, luminosity, and radius relationship underpin the  
 164 circumbinary solutions of Georgakaraos & Eggl (2015), which affect orbital properties, stellar  
 165 fluxes at the planet, and magnitude of eclipse effects respectively. Note that using our derived  
 166 mass, luminosity, radius, and effective temperature relationships, a  $1 M_{\odot}$  star has an effective  
 167 temperature ( $T_{\text{eff}}$ ) of 5714 K, a radius of  $0.99 R_{\odot}$ , and a luminosity of  $0.958 L_{\odot}$ . Thus, our  
 168 nominal solar twin is slightly cooler, smaller, and dimmer than our Sun. For this study, we have  
 169 considered four different stars in shown in Table 1; G2V ( $1 M_{\odot}$ , 5714 K), K0V ( $0.85 M_{\odot}$ , 5229  
 170 K), M0V ( $0.55 M_{\odot}$ , 3824 K), and M5V ( $0.15 M_{\odot}$ , 3070 K).

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#	Stellar Type	Mass ( $M_{\odot}$ )	$T_{\text{eff}}$ (K)	Radius ( $R_{\odot}$ )	Luminosity ( $L_{\odot}$ )
1	G2V	1.0	5714	0.999	0.958
2	K0V	0.85	5229	0.815	0.448
3	M0V	0.55	3824	0.556	0.0578
4	M5V	0.15	3070	0.194	0.00298

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*Table 1: Stellar characteristics assumed in this study. All circumbinary systems studied are a combination of these stars.*

## 175 2.2. Description of the 3-D General Circulation Model

176 Here we use a 3-D climate system model to simulate Earth-like planets in a variety of  
 177 dynamically stable circumbinary systems. We use the Community Earth System Model version  
 178 1.2.1 from the National Center for Atmospheric Research (Neale et al. 2010), along with the  
 179 ExoCAM modeling package (<https://github.com/storyofthewolf/ExoCAM>), which facilitates  
 180 simulations of extrasolar planets. As noted in section 2.1, we have further modified ExoCAM to  
 181 incorporate the analytic circumbinary orbital system solutions of Georgakarakos and Eggl  
 182 (2015), which drive the variations in stellar flux received by planets in circumbinary systems.  
 183 For this study the 3-D model was further developed so that the light from each star is  
 184 individually resolved in the radiative transfer computations, which accounts for stellar spectra

185 and zenith angle of each star respectively. Thus, we are explicitly resolving the changes in  
186 received stellar flux driven by both stars simultaneously and self-consistently. To model the  
187 stellar energy distribution, we use BT-Settl spectra for each star in the circumbinary pair (Allard  
188 et al. 2007) assuming  $[\text{Fe}/\text{H}] = 0$  and  $\log(g) = 5$ , and interpolated to desired  $T_{\text{eff}}$  (Table 1). In all  
189 cases we assume that the planets receive a stellar flux of  $1360 \text{ Wm}^{-2}$  averaged over many orbits  
190 (including the effect of eclipses). In this study, one of the circumbinary stars is always assumed  
191 to be a Sun-like star with mass of  $1 M_{\odot}$ . The addition of a secondary star means that the  
192 combined intrinsic luminosity of the circumbinary pair is greater than our Sun alone, and thus  
193 our simulated planets' semi-major axes and orbital periods both must be larger than Earth's in  
194 order to ensure the time-average received stellar flux remains  $1360 \text{ Wm}^{-2}$ .

195 We assume Earth-like planets that have the same mass, radius, surface gravity,  
196 continental configuration, and diurnal period as the Earth. The planet's absolute rotation period  
197 (i.e. the sidereal period), is modified slightly in each case to ensure that the diurnal period  
198 remains fixed at 24 hours, despite changes to the planet's orbital period in each case, but these  
199 variations are generally negligible and unlikely to yield any discernible effects on climate  
200 (Charnay et al. 2013; Wolf & Toon, 2014) We focus primarily on cases where the planet has  $0^{\circ}$   
201 obliquity in order to isolate the effects driven by the circumbinary orbital configuration, but  
202 sensitivity tests are also conducted with a  $23.5^{\circ}$  obliquity. The planets are assumed to have a  
203 total atmospheric pressure of 1 bar, with 400 ppm of  $\text{CO}_2$  and the remainder as  $\text{N}_2$ . We do not  
204 consider  $\text{O}_2$ ,  $\text{O}_3$ , nor the effects of photochemistry in this study. The primary climatological  
205 contribution of  $\text{O}_2$  is through its contribution to the total atmospheric pressure, which causes  
206 scattering and pressure broadening. While  $\text{O}_2$  could have a significant effect on stratospheric  
207 temperatures, replacing  $\text{O}_2$  with  $\text{N}_2$  has a negligible effect on the surface climate (Wolf & Toon,

208 2013), which is our primary interest in this paper. Water vapor, liquid water clouds, and ice  
209 water clouds are variable and advected constituents in the model.

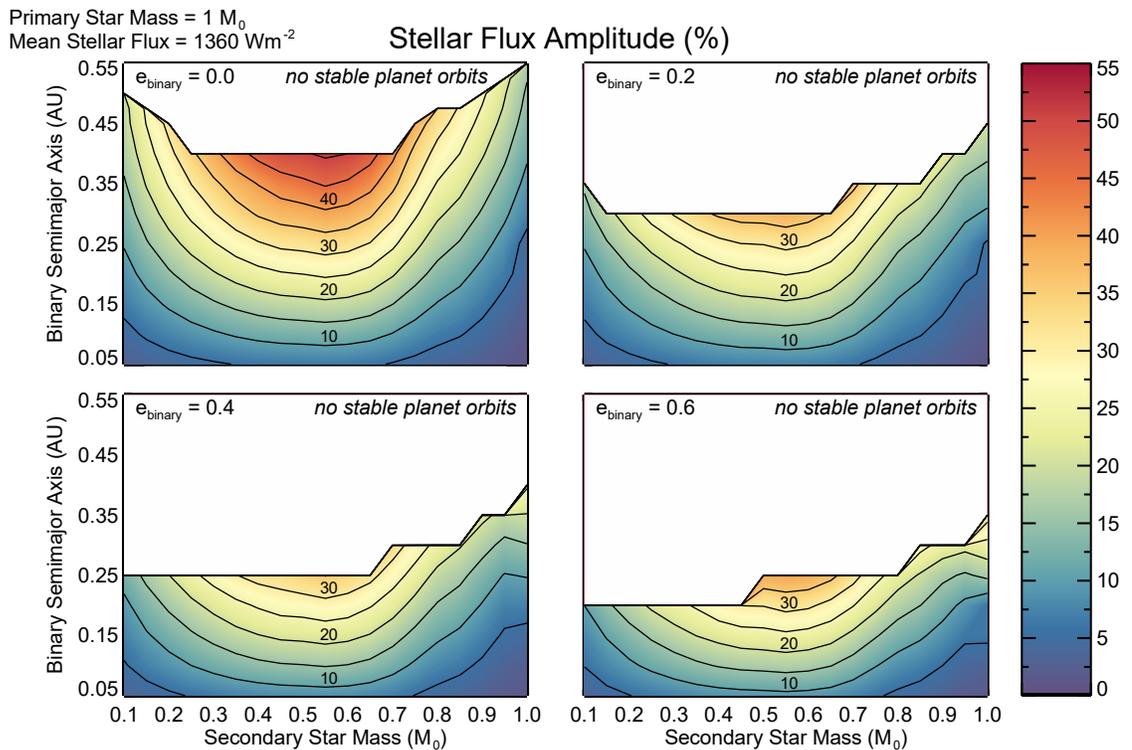
210 We assume the modern Earth continental configuration with the ocean treated as a 50-  
211 meter deep slab ocean. Implied ocean heat transport is included via heat flux convergence terms  
212 to approximate that of the modern Earth, where generally cold water is transported equatorward  
213 along western coastlines, and warm water is transported poleward along eastern coastlines (Bitz  
214 et al. 2012). We have removed seasonal dependencies from the heat flux convergence terms, and  
215 instead use the annual mean implied transport at all times. Simulations are run with  $4^\circ \times 5^\circ$   
216 horizontal resolution and 40 vertical layers extending from the surface up to  $\sim 1$  mbar pressures,  
217 using a finite-volume dynamical core (Lin & Rood, 1996). Land, ocean, snow, and ice albedos  
218 are treated with two bands, visible and near infrared split at  $0.77 \mu\text{m}$ . Open-ocean albedos are  
219 held fixed at 0.06 at all wavelengths. Sea ice albedos in visible (near-IR) are given as 0.68 (0.3),  
220 and snow albedos are given as 0.91 (0.63). Land albedos depend on the assumed surface type,  
221 surface color, and soil saturation. Permanent glacier ice has been removed from Antarctica and  
222 Greenland, however snow can accumulate and cover these regions. The specific heat capacity of  
223 sea water is  $3996 \text{ J kg}^{-1} \text{ K}^{-1}$ , and the volumetric heat capacity of land surfaces varies from  
224 between  $2 \times 10^6$  and  $2.5 \times 10^6 \text{ J m}^{-3} \text{ K}^{-1}$  depending on the surface type (i.e. soil types and vegetation  
225 types). Surface types are assumed to be identical to the modern Earth.

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### 227 **3. Analytic Solutions for Earth-like Worlds in Circumbinary Systems**

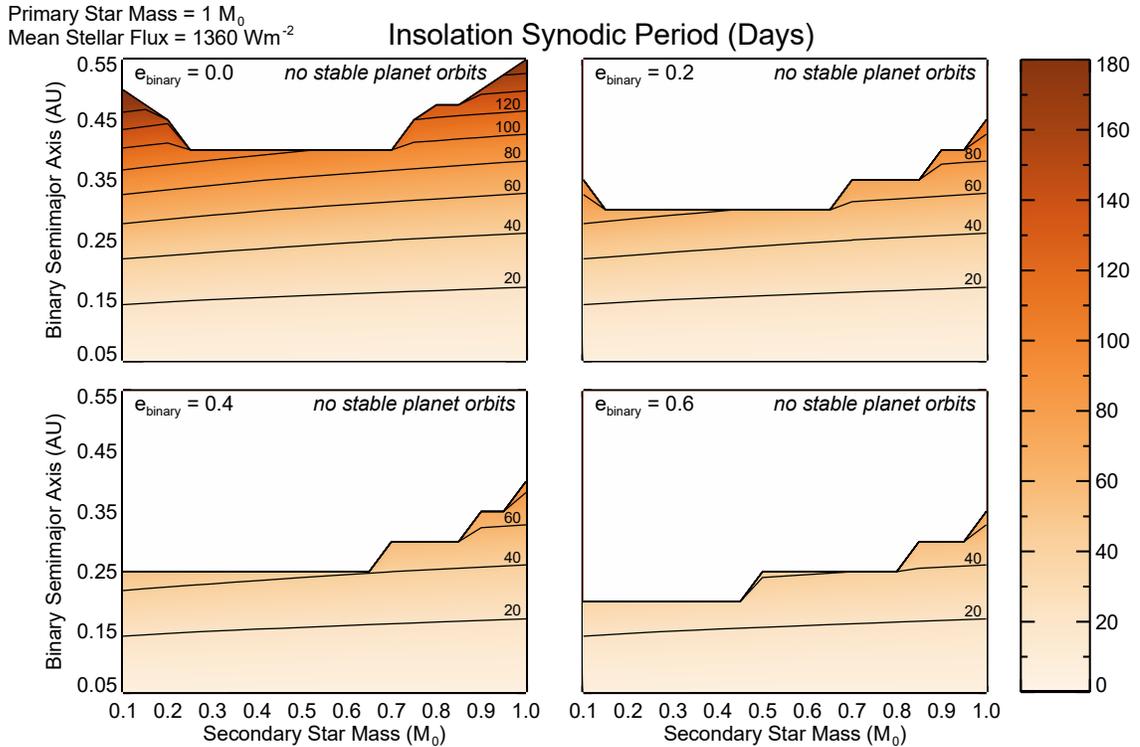
228 First, we examine which specific circumbinary system architectures can harbor  
229 dynamically stable planets at modern Earth insolation levels, and which architectures can drive  
230 the largest magnitude and longest period changes to the incident stellar flux. To do so, we

231 conduct an array of off-line calculations using our modified version of the Georgakarakos &  
 232 Ettl (2015) model, solving for planets receiving  $1360 \text{ Wm}^{-2}$  of stellar flux on time-average.  
 233 Figures 1, 2, and 3 describe the steady state properties of Earth-like planets in circumbinary  
 234 systems considering the stellar flux amplitude, the insolation synodic period, and the planet  
 235 eccentricity respectively. In all figures, we assume that the primary star has a mass of  $1 M_{\odot}$ , with  
 236 the secondary star mass given on the x-axis, and the separation of the binary stars in AU given  
 237 on the y-axis. The time mean stellar flux is  $1360 \text{ Wm}^{-2}$  in all cases. Binary eccentricities of 0,  
 238 0.2, 0.4, and 0.6 are shown and labeled in each panel. The white space at the top of each panel  
 239 are regions in phase space where no dynamically stable solution exists for our Earth-like planets  
 240 as described.



**Figure 1:** The stellar flux amplitudes possible for dynamically stable solutions for Earth-like planets in circumbinary systems derived from the Georgakarakos & Ettl (2015) model. The stellar flux amplitude is the percent change in the received stellar flux by the planet relative to the time-mean value ( $1360 \text{ Wm}^{-2}$ ), and is indicated by contours and the color bar. White space indicates regions in our phase space where no stable solutions exist. Shown are solutions for circumbinary pairs having eccentricities of 0, 0.2, 0.4, and 0.6. The primary star is always assumed to be  $1 M_{\odot}$  while the secondary star mass is given on the x-axis in units of stellar masses. The separation of the binary stars is given on the y-axis in units of AU.

249 Figure 1 shows the stellar flux amplitude (%), defined as the maximum received flux by  
250 the planet at any time, minus the minimum received stellar flux, divided by the time-average flux  
251 ( $1360 \text{ Wm}^{-2}$ ) and multiplied 100 to yield values in percentage. In this calculation we have  
252 excluded the effect of eclipses in the stellar flux amplitude, because the precipitous drop in  
253 stellar radiation during a binary eclipse only lasts for a few hours, and results in only a small  
254 time-integrated reduction in the total flux received by the planet with no apparent effect on  
255 planetary climate in our full climate calculations described in section 4. In Figure 1, note that  
256 the greatest stellar flux amplitude variation occurs when the secondary star has a mass of about  
257 half ( $\sim 0.55 M_{\odot}$ ) of the primary star. In this case, assuming the binary pair has zero eccentricity, a  
258 maximum stellar flux amplitude of  $\sim 50\%$  can occur when the binary separation is 0.4 AU. If the  
259 binary separation is increased further, the binary orbit becomes too wide, and no dynamically  
260 stable planet can exist in an orbit where it would receive  $1360 \text{ Wm}^{-2}$  of time-averaged flux.  
261 Although stable planets could still exist on larger orbits, they would necessarily receive less  
262 incident stellar flux than the present-day Earth. Likewise, as the eccentricity of the circumbinary  
263 pair is increased, the phase space where an Earth-like planet can remain dynamically stable is  
264 reduced. The elongated orbits of eccentric circumbinary pairs can effectively scatter away  
265 surrounding planets, fewer permutations of stable planet orbits exist in our phase space, and the  
266 maximum possible amplitudes of stellar flux variation are at most  $\sim 35\%$ .



267  
268 **Figure 2:** The insolation synodic period in Earth days derived from the Georgakarakos & Eggl (2015) model. The  
269 insolation synodic period is the period of variation in circumbinary driven insolation as viewed from the reference  
270 frame of the planet, and is indicated by contours and the color bar. The plot axis and conventions are identical to  
271 figure 1.

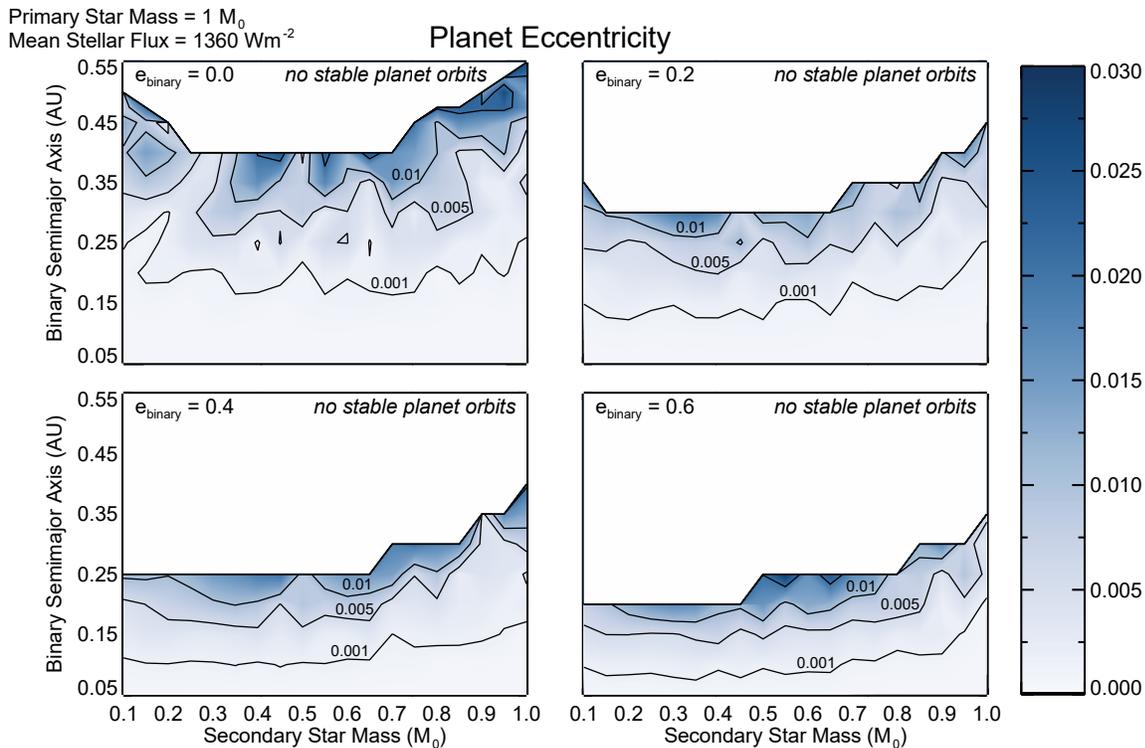
272 As noted, the time period of incident stellar flux variation is also critical to the problem  
273 of circumbinary driven climate variability. Figure 2 shows the insolation synodic period, which  
274 we define as the period of variation in the circumbinary driven insolation when viewed in the  
275 reference frame of the planet. The insolation synodic period is greater than the binary orbital  
276 period, due to the motion of the planet relative to the circumbinary pair, much like the diurnal  
277 period of a planet is greater than its sidereal period. For circumbinary pairs with zero  
278 eccentricity and separations greater than 0.3 AU, the insolation synodic period can exceed 75  
279 days, and reach as high as  $\sim 200$  days in certain corner cases of our phase space. However, for  
280 higher eccentricity binaries, the insolation synodic period generally remains less than about 50  
281 days, due to the constriction in semimajor axes possible which permit dynamically stable Earth-  
282 like planets. In Figure 3, we show the planet eccentricities, which are direct output of the

283 Georgakarakos & Eggl (2015) model. Planet eccentricities remain relatively small ( $< 0.03$ ) in all  
 284 cases studied; however, the highest planet eccentricities predicted cluster around the more  
 285 extreme circumbinary architectures with the widest binary separations. The planet's eccentricity  
 286 introduces another mode of variation in the received stellar by the planet, and as we will see  
 287 next, in some cases the stellar flux variations due to planetary eccentricity can sometimes  
 288 outweigh that driven by the circumbinary pair.

#	Stellar Pair	Binary Separation (AU)	Binary Eccentricity	Binary Period (days)	Planet Period (days)	Planet Eccentricity	Insolation Synodic Period (Days)	Planet Semi-major axis (AU)	Stellar Flux amplitude (%)	Net Transit Flux Reduction (%)
1	G2V	--	0	--	353.9	0	--	0.979	--	--
2	G2V-G2V	0.55	0	105.3	449.8	2.17e-2	137.5	1.448	16.06	0.174
3	G2V-G2V	0.3	0	42.4	428.9	1.48e-2	47.1	1.402	6.77	0.372
4	G2V-G2V	0.1	0	8.2	418.1	8.54e-5	8.3	1.378	0.89	1.224
5	G2V-K0V	0.45	0	81.1	368.3	5.00e-3	103.9	1.235	28.35	0.179
6	G2V-K0V	0.25	0	33.6	353.0	1.08e-2	37.1	1.200	14.63	0.371
7	G2V-M0V	0.4	0	74.2	313.2	5.85e-3	97.3	1.045	51.29	0.099
8	G2V-M0V	0.2	0	26.2	300.8	3.52e-3	28.7	1.017	25.80	0.227
9	G2V-M5V	0.45	0	102.8	340.5	2.18e-2	147.3	0.999	25.24	0.012
10	G2V-M5V	0.25	0	42.6	333.8	4.85e-3	48.8	0.987	14.55	0.022
11	G2V-G2V	0.35	0.6	53.5	439.4	7.74e-2	60.9	1.425	32.81	0.305
12	G2V-K0V	0.3	0.4	44.1	358.7	4.34e-2	50.3	1.213	28.80	0.298

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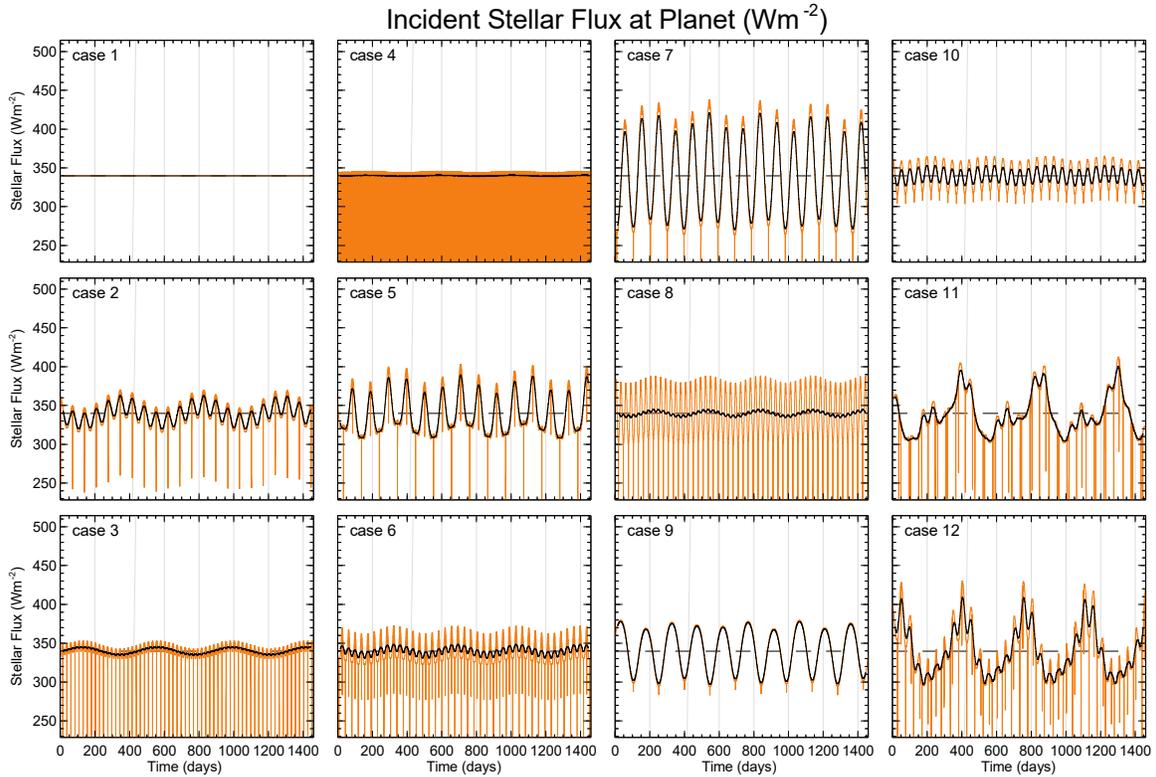
Table 2: Circumbinary system characteristics used in this study.



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Figure 3: The planet eccentricity derived from the Georgakarakos & Eggl (2015) model, indicated by contours and the color bar. The plot axis and conventions are identical to figure 1.

294 In Figure 4 we show the time-dependent behavior of the global mean combined stellar  
295 flux from both stars impinging on our Earth-like planets, which includes the effect of binary  
296 eclipses, manifested as vertical straight lines, and the effect of planet eccentricities. The orange  
297 lines indicate the instantaneous stellar flux received by the planet tabulated once per hour. The  
298 black lines indicate the 30-day running average of the incident stellar flux. We present 12  
299 different time-varying stellar flux patterns in Figure 4, corresponding to the system descriptions  
300 shown in Table 2. Figure 4 is labeled according to the cases listed in Table 2. Case 1 is our  
301 control case, where the planet orbits a single Sun-like star and the planet's orbital eccentricity is  
302 zero. Here, naturally, the global mean stellar flux received by the planet is constant in time, with  
303 a value of  $340 \text{ Wm}^{-2}$  (i.e.  $1360 \text{ Wm}^{-2}$  divided by 4, accounting for diurnal and geometric effects).  
304 All other cases shown are for a variety of permutations of physically plausible circumbinary  
305 systems. While we are most interested in the cases where the variation in the stellar received by  
306 the planet is of the largest amplitude and longest period, we also show some cases where the  
307 variations in the insolation synodic period and flux amplitude are small. Note that periodic  
308 variations on order  $\sim 100$  days or less are driven by the motion of the circumbinaries, while  
309 periodic variations on order of  $\sim 400$  days are driven by the eccentricity of the planets. Recall,  
310 that the planet eccentricity in these circumbinary systems is not a free parameter, but is a  
311 prediction outputted from Georgakarakos & Eggl (2015) model.



312

313 **Figure 4:** Variations in the stellar flux received by our test planet for the 12 scenarios described in Table 2. Note,  
 314 that case 1 is our single star control case, and the time-mean stellar flux in all cases is  $1360 \text{ Wm}^{-2}$ . The orange lines  
 315 are the instantaneous flux received by the planet tabulated once per hour, while the black lines are the 30-day  
 316 running average.

317

Case 7 illustrates a scenario of maximum amplitude stellar flux variation. This case

318

features a G2V-M0V star pairing with a binary separation of 0.4 AU, a stellar flux amplitude of

319

$\sim 51\%$ , and an insolation synodic period 97.3 days. In this case the amplitude of the stellar flux

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variation is dominated by the motion of the larger, brighter G2V star, which is 18 times more

321

luminous than the secondary M0V star (Table 1). Note also the subtle variation in the stellar flux

322

maxima in case 7, driven by the non-zero eccentricity of the planet. Case 8 shows the identical

323

star pairing (G2V-M0V), but with the binary separation cut in half to 0.2 AU. In this case the

324

stellar flux amplitude is reduced to  $\sim 25\%$ , while the insolation synodic period is reduced to 28.7

325

days. In this case the 30-day running average variation in incident stellar flux is now dominated

326

by the planet's eccentricity instead of the circumbinary effect. If we consider a smaller, dimmer,

327 secondary star (M5V), the stellar flux amplitudes decrease further, because the motion of the  
328 larger, brighter relative to the center of mass decreases for a constant binary separation (e.g.  
329 cases 9 and 10 versus cases 7 and 8 respectively).

330 Cases 2, 3, and 4 feature two equal mass stars (G2V-G2V pairings). In the case of two  
331 equal mass, equally luminous stars, the effective period of stellar flux variation at the planet is  
332 half that of the instellation synodic periodic. With each star being identical, the change in  
333 incident flux from each star is precisely balanced as each star approaches/recedes from the planet  
334 in tandem and at an equal velocity. The maximum incident flux occurs whenever one of the two  
335 stars is closest to the planet. The minimum incident flux occurs when the binary stars are at  
336 quadrature relative to the planet. For one complete orbit of the binary pair relative to the planet,  
337 thus, there are two cycles of insolation that occur. Similar as before, as the binary separation  
338 between the pair is reduced, the stellar flux amplitude and insolation synodic period decrease.  
339 Note in case 2, and case 3, that again the planet eccentricity imparts a significant signal on the  
340 total received insolation, rivaling the changes due to the circumbinary. Case 4 highlights when  
341 the effect of binary eclipsing is maximized. Note that the frequency of eclipses washes out the  
342 panel in Figure 4. Here, the circumbinary stars orbit each other every 8.2 days, and eclipses  
343 occur every 8.3 days. Because the stars are of equal size, each eclipse reduces the total flux  
344 received by the planet by ~46%. However, due to the brevity of eclipses, still, the net reduction  
345 in the stellar flux is only 1.2% compared to an idealized architecture where no eclipses occur.  
346 However, for all other cases studied, eclipses are infrequent and the net flux reductions or no  
347 more than a few tenths of a percent.

348 When the stellar masses are different by less than about 20% (case 5 and case 6), an  
349 irregular pattern of stellar insolation results due to the  $1/d^2$  dependence of radiative flux

350 combined with the  $\sim M^4$  relationship of stellar luminosities. Naturally, the stellar insolation is  
351 greatest when the larger star is closest to the planet, however the lesser star can contribute a  
352 meaningful fraction of the total irradiance ( $\sim 10\%$ ). As the larger star moves away from the  
353 planet, the total insolation decreases. The stellar flux reaches a minimum at quadrature, but then  
354 begins to increase as the lesser star moves closer to the planet. There is a secondary maximum  
355 that occurs when the secondary star is closest to the planet, but because the secondary is dimmer  
356 this stellar flux maximum is lesser. The brighter star still contributes to the largest fraction of  
357 stellar radiation, but now the secondary contributes up to 40%.

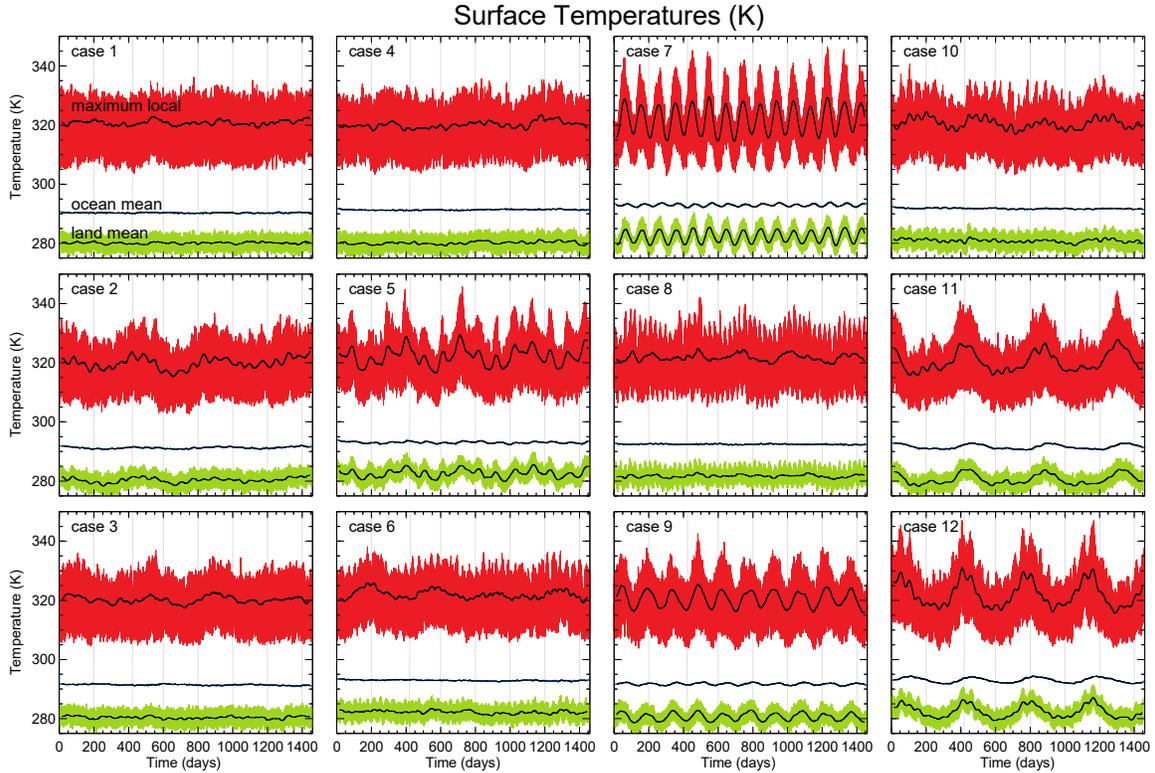
358 Finally, we also illustrate two scenarios (case 11 and 12) where the circumbinary pair has  
359 significant eccentricity. Note that again we have cherry-picked system architectures where the  
360 variation in stellar flux at the planet is largest. Here, complex interactions occur between star  
361 and planet positions yield complex time-dependent patterns of stellar insolation. Note that in  
362 both cases 11 and 12, the variation in incident flux is dominated by the planet's eccentricity, with  
363 the circumbinary signal laid over top. While the stellar flux variations in cases 11 and 12 (and  
364 others) can be dominated by the planet eccentricity rather than the circumbinary motions, recall  
365 that the non-zero eccentricity of these planets is a direct consequence of the complex orbital  
366 dynamical interactions present in circumbinary systems, and thus must be accounted for in any  
367 self-consistent circumbinary planets residing in the habitable zone (Georgakarakos & Eggl,  
368 2015; Popp & Eggl, 2017).

369

#### 370 **4. 3-D Climate modelling results**

371 Next, we conducted a variety of 3-D climate simulations to illustrate the possible changes  
372 to climate and weather for Earth-similar planets in circumbinary systems. Table 2 shows a list of

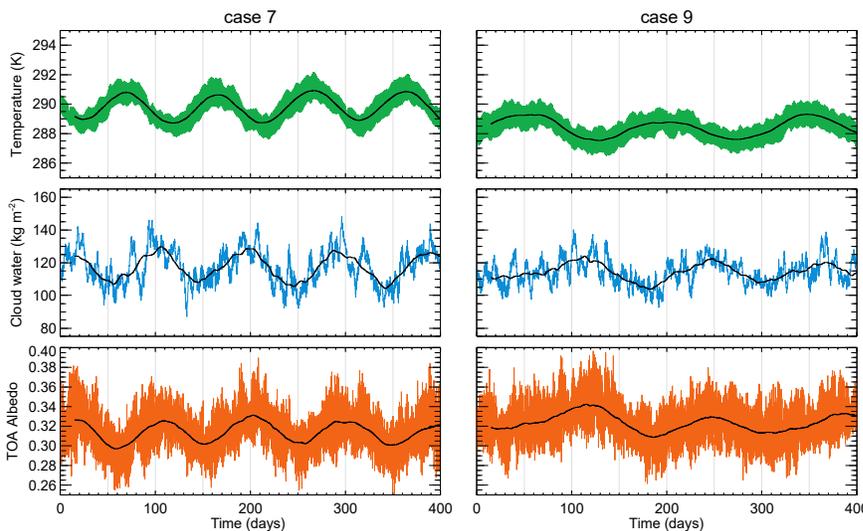
373 the circumbinary system architectures that were explored, along with relevant statistics that  
374 describe the systems. We explore 12 unique system configurations, with stellar flux evolution  
375 shown in Figure 4. For our primary set of simulations, we consider Earth-like planets with zero  
376 obliquity. Recall that in all cases the time-averaged stellar flux received by the planet is  $1360$   
377  $\text{Wm}^{-2}$ . Figure 5 illustrates surface temperature variations as a function of time for Earth-like  
378 planets with zero obliquity. Shown are the maximum local temperature, taken from anywhere on  
379 the planet (red), the global mean ocean temperatures, and the global mean land surface mean  
380 temperature (green). The solid black lines indicate the 30-day running mean temperature, while  
381 the colors show the instantaneous values sampled once every hour. Recall that case 1 is the  
382 single star control case. Due to the significant thermal inertia of the oceans, the instantaneous  
383 variations in ocean temperatures are virtually indistinguishable from the 30-day running mean  
384 value, and even under the most extreme circumbinary architectures ocean mean temperatures do  
385 not vary more than a few degrees Kelvin, whether driven by circumbinary forcings (e.g. case 7  
386 and case 9) or via planet eccentricity effects (case 11 and case 12). Thus, for ocean dominated  
387 planets in the habitable zones of circumbinary stars, even in the most extreme physically  
388 plausible causes, dramatic global changes in climate do not occur due to the thermal inertia of  
389 the oceans, combined with the fact that insolation synodic periods (Figure 2) are generally no  
390 more than  $\sim 150$  Earth days. Thus, the ocean does not have sufficient time to catastrophically  
391 cool or warm. This result agrees with the results of Popp & Eggl (2016) and Way et al. (2017).  
392 For evaluating the global mean climate of habitable circumbinary planets, the mean-flux  
393 approximation appears to hold true (Bolmont et al. 2016).



394 *Figure 5 shows the temperatures variations corresponding to our twelve system configurations, described in Table 2*  
 395 *and with stellar flux variations shown in Figure 4. We show the instantaneous maximum local temperature (red), the*  
 396 *land surface mean (green), and the ocean mean temperature. Black lines indicated the 30-day running mean, while*  
 397 *the colors show the instantaneous values at a 1-hour time cadence.*  
 398

399 Greater variations in the land surface temperatures compared to ocean temperatures are  
 400 evident in Figure 5, which agrees with the energy balance modeling of Haqq-Misra et al. (2019).  
 401 The relatively low thermal inertia of land surfaces allows continents to heat up and cool much  
 402 more quickly, thus signals of the circumbinary forcing and the planet eccentricity are more  
 403 apparent. Variations to land surface mean and local maximum temperatures are evident for  
 404 several cases. Cases 5, 7, and 9 indicate a relatively strong circumbinary signal in land surface  
 405 temperatures, implying a circumbinary pseudo-seasonality is induced in these systems. In our  
 406 most extreme case, case 7, the 30-day running mean land surface temperatures may vary by  $\sim 5$  K  
 407 between maximum and minimum insolation phases of the circumbinary motion, while the 30-  
 408 day running mean maximum local temperatures may vary by as much as  $\sim 12$  K. If we consider  
 409 that instantaneous maximum temperatures, in case 7 maximum insolation phases may drive local

410 land surface temperatures in the tropics to reach 345 K, whereas in our single star control  
 411 simulation, maximum local temperatures generally hover near 330 K. Note that these extreme  
 412 local instantaneous temperatures are also seen in case 11 and case 12, although these temperature  
 413 maximums are driven primarily by the planet eccentricity and not from the circumbinary forcing.  
 414 Still, note that we are highlighting cases of maximum plausible circumbinary driven variations in  
 415 stellar flux. Even in such cases of maximal stellar forcing, the heat stress thresholds for human  
 416 habitability based on wet bulb temperatures (Sherwood & Huber 2010) are not exceeded during  
 417 warm periods because the maximum temperatures indicated generally occur over regions of low  
 418 relative humidity (i.e. subtropical desert areas). For more moderate system architectures (e.g.  
 419 cases 3, 4, 6, 8) the effects of local climate are muted or nearly indistinguishable from the single  
 420 star control case.



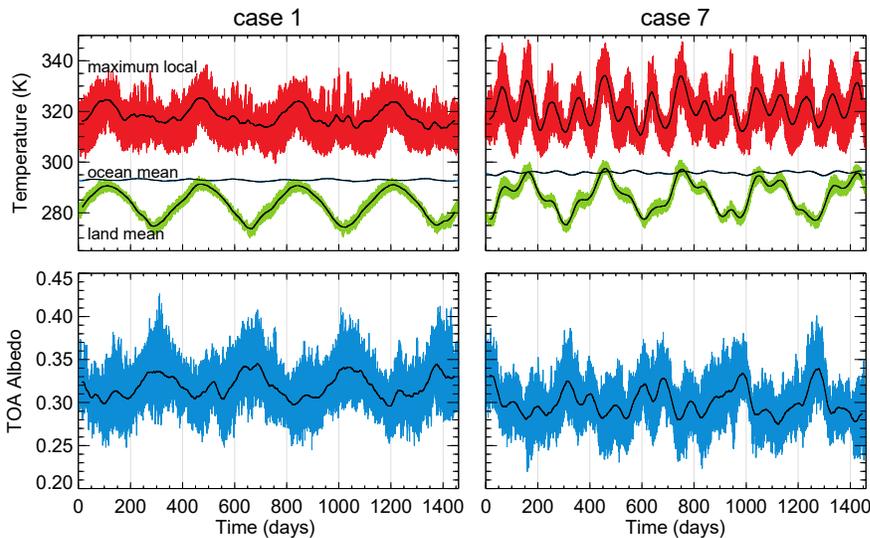
421 *Figure 6 illustrates that there is a subtle phase off-set between cloud water (middle row) and the variations in the*  
 422 *land surface temperatures (top row). While the TOA albedo, naturally, remains in phase with the cloud water, clouds*  
 423 *and albedo subtly lag the temperature maximum. Still, elevated clouds and albedo during warm phases act as a*  
 424 *negative feedback, stabilizing the climate against circumbinary driven effects.*  
 425

426 In Figure 6 we illustrate the land surface mean temperature, along with the global mean  
 427 cloud water column, and the TOA albedo, for case 7 and case 9 over several insolation synodic  
 428 periods. We see that there is a phase off-set between the land surface temperatures, and the

429 cloud water and albedo, respectively. Cloud water and albedo, as expected, remain in phase due  
430 to the strong contribution of clouds to modulating the planetary albedo. However, the maximum  
431 in cloud water and the planetary albedo lag behind the land surface temperature maximum.  
432 Indeed, the warmer the air is, the more vapor it can hold before condensing into cloud droplets.  
433 As the planet temperature warms, more water vapor is evaporated into the atmosphere. When  
434 temperatures begin to cool, after the insolation maximum has passed, the cooling temperatures  
435 decrease the ability of the atmosphere to hold water vapor, and so water condenses to form  
436 clouds. As this process is not instantaneous a slight offset is visible between the minimum cloud  
437 water and maximum temperature. Still, the generally trend of rising cloud water and albedo at  
438 and just past the insolation peak, represent a negative feedback at work that acts to mute  
439 circumbinary driven climate changes.

440 We have also conducted sensitivity tests considering planets with a  $23.5^\circ$  obliquity. In  
441 the majority of our cases, climate changes driven by obliquity seasons completely swamp any  
442 circumbinary driven climate signals (not shown). In Figure 7 we show the maximum local,  
443 ocean mean, and land surface mean temperatures for case 1 (single star control) and case 7 (our  
444 most extreme circumbinary architecture), while assuming the planet has a  $23.5^\circ$  obliquity. Note  
445 that for the control case, seasonal variations in land surface temperatures are a result of the  
446 asymmetric distribution of continents on Earth. Most of the land mass is located in the northern  
447 hemisphere, land surfaces have a lower thermal inertia, and thus the northern hemisphere  
448 summer exhibits warmer temperatures. When non-zero obliquity is coupled with an extreme  
449 circumbinary forcing, temperature variations can be either strengthened or weakened depending  
450 on the alignment of the eccentricity seasons with the circumbinary variations. When northern  
451 hemisphere summer aligns with a circumbinary maximum insolation phase, then land surface

452 temperatures can be more extreme. However, if conversely the circumbinary minimum occurs  
453 during the northern hemisphere summer, climatological effects are muted. Because the planet  
454 orbital periods do not divide evenly into the circumbinary periods, there may be longer term  
455 trends not explored here as the mutual planet and circumbinary orbits may only precisely repeat  
456 on many decades to century scale timescales.



457 *Figure 7 illustrates the effect of non-zero obliquity. Recall case 1 is the single star control case. For all but the most*  
458 *extreme circumbinary scenarios, seasonal effects from non-zero obliquity coupled with asymmetric north-south land*  
459 *distributions have dominate over the circumbinary effects.*  
460

461

#### 462 4. Discussion and Conclusions

463 Our results indicate that there is no climate catastrophe that precludes the habitability of  
464 Earth-like planets in the HZs of circumbinary systems. Earth-like planets in circumbinary  
465 systems are as likely to remain as habitable as planets in single star systems. Circumbinary  
466 driven variations in surface climate are insignificant in all but the most extreme system  
467 architectures. In the most extreme physically plausible circumbinary system architectures, land  
468 surface temperatures can exhibit an additional mode of variability with predictable periods of  
469 more intense heat waves; however, the thermal inertia of habitable planets (e.g. ocean-

470 dominated) is such that radical climate changes are prevented. Aspects such as the stochastic  
471 nature of planet formation, the acquisition and availability of volatiles, the orbital dynamical  
472 stability of planets, and the appropriate balance of time-averaged stellar flux and atmospheric  
473 greenhouse effect, are all more important arbiters of habitability than are periodic variations in  
474 stellar insolation found in circumbinary systems.

475         There are several caveats that should come with this conclusion. Here, we have limited  
476 our study to Earth-twin planets, and we have explored the effects of other planetary parameter  
477 variations such as planetary rotation period, orbital eccentricity, surface pressure and  
478 atmospheric composition. For instance, desiccated planets with exceedingly thin atmospheres  
479 and no oceans would be expected to show greater degree of temperature variations compared to  
480 the Earth-like planets studied here. Conversely, planets with thicker atmospheres, or larger  
481 fractions of ocean compared to land, should experience even less climate variability than is  
482 shown here (May & Rauscher, 2016). Results from slow rotating planets indicate that their  
483 climate sensitivity to changing stellar fluxes is much less than for rapidly rotating planets (Yang  
484 et al. 2013; Way et al. 2019), and thus we do not expect changes in the planetary rotation rate to  
485 significantly change the resiliency of habitable climates to circumbinary forcings. Furthermore,  
486 larger assumed planetary eccentricity or obliquity would swamp any circumbinary signals, as is  
487 already evident in Figure 5 (cases 11 and 12) and Figure 7. If we find that HZ planets are as  
488 common in circumbinary systems as in single star systems, we can conclude that they have an  
489 equal chance of hosting-life. Considering that perhaps half of Sun-like stars exist in binary and  
490 higher multiple systems, the resilience of habitable worlds in these systems essentially doubles  
491 the chances of life-hosting planets in our galaxy.

492           The prospect of finding habitable Earth-like planets in circumbinary systems raises  
493 further speculation as to how a circumbinary biosphere might differ from a single-star one.  
494 Several studies have suggested that the trajectory of biological evolution on Earth has been  
495 driven in part by orbital Milankovitch cycles and the resulting changes in regional insolation and  
496 glaciation (Bennett 1990; Jansson & Dynesius 2002; Lister 2004). The “seasonality hypothesis”  
497 even suggests that the evolution of traits such as animal body size could be connected with  
498 environmental pressures associated with climate cycles (Troost et al. 2009). Such ideas suggest  
499 that the effect of circumbinary seasons, as well as circumbinary Milankovitch forcing (Forgan  
500 2016), could similarly act as an evolutionary selection mechanism. Spectral biosignatures on  
501 circumbinary planets could likewise exhibit unique features that depend upon such such  
502 variations; for example, the observability of a planet’s photosynthetic “red edge” (Seager et al.  
503 2005; Turnbull et al. 2006; Kiang et al. 2007) could depend upon the phase of its circumbinary  
504 seasons. Seasonal and orbital variations provide a selection pressure for life on Earth, so we can  
505 reasonably expect that such pressures should apply to any life that develops on circumbinary  
506 planets.

507           We can even speculate further as to the implications of habitable circumbinary planets for  
508 the development of technology. DeVito (2013) discusses the connection between human  
509 mathematical systems and our physical spatial environment; for example, the model of natural  
510 numbers could have been extrapolated from observing the consistency of the day-night cycle on  
511 Earth. The period of insolation experienced by circumbinary planets poses a more complex  
512 pattern than a single-star system: the inference of natural numbers from the setting stars might be  
513 more difficult in such a system, although perhaps such added complexity would enhance the  
514 evolution of technology. The “proportional evolutionary time” hypothesis (Haqq-Misra 2019)

515 suggests that the evolutionary timescale for complex life depends upon the availability of free  
516 energy between 200 and 1000 nm. The incident radiation for an Earth-like circumbinary planet  
517 with different spectral host stars (e.g., a G-dwarf primary and K-dwarf secondary) will show an  
518 increase in photons within this wavelength range compared to a planet orbiting a single-star at an  
519 equivalent orbital distance. Thus, the proportional evolutionary time hypothesis suggests that  
520 such circumbinary systems are at least as promising candidates as single star systems to search  
521 for evidence of complex life or technology. Finally, the possibility that life, and even  
522 technology, could exist around circumbinary systems raises an important philosophical question:  
523 “Why do we find ourselves around a single star instead of a binary pair?” If binary systems are  
524 as likely or more likely to support life, then this might give us reason to suspect that the  
525 circumstances of life on Earth are less typical than we might otherwise assume (e.g., Haqq-Misra  
526 et al. 2018). The answer to this question, and many of our other speculations, will require  
527 detection and spectral characterization of a statistically significant sample of circumbinary  
528 planets in order to begin estimating the prevalence of biosignatures and technosignatures in such  
529 systems.

530

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543 enable modeling circumbinary planets, along with necessary initial files for implementation into  
544 CESM, are available on lead-authors Github page; <https://github.com/storyofthewolf/ExoCAM/>.  
545 In particular, refer to the “circumbinary” configuration on this github page. The primary model  
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