

1     **Impacts of the January 2022 Tonga volcanic eruption on**  
2             **the ionospheric dynamo: ICON-MIGHTI and Swarm**  
3             **observations of extreme neutral winds and currents**

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11     **Key Points:**

- 12     • Extreme thermospheric winds and ionospheric currents were observed in coordi-  
13     nated space/ground-based measurements, ten hours post-eruption
- 14     • The westward electrojet current when the Lamb wave reaches the dayside is con-  
15     sistent with recent studies of the wind-driven electrojet
- 16     • Observations of linked dynamo processes provide direct evidence of the space-weather  
17     impacts of acute lower atmospheric forcing

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**18 Abstract**

19 The eruption of the Hunga Tonga-Hunga Ha’apai volcano on 15 January 2022 trig-  
20 gered atmospheric waves at all altitudes. The NASA Ionospheric Connection Explorer  
21 (ICON) and ESA Swarm satellites were well placed to observe its impact on the iono-  
22 spheric wind dynamo. After the Lamb wave entered the dayside, Swarm A observed an  
23 eastward and then westward equatorial electrojet (EEJ) on two consecutive orbits, each  
24 with magnitudes exceeding the 99.9th percentile of typically observed values. ICON si-  
25 multaneously observed the neutral wind (90–300 km altitude) at approximately the same  
26 distance from Tonga. The observed neutral winds were also extreme (>99.9th percentile  
27 at some altitudes). The covariation of EEJ and winds is consistent with recent theoret-  
28 ical and observational results, indicating that the westward electrojet is driven by strong  
29 westward winds in the Pedersen region (~120–150 km). These observations imply that  
30 the dynamo is a key mechanism in the ionospheric response to the Tonga disturbance.

**31 Plain Language Summary**

32 The January 2022 Tonga volcanic eruption caused atmospheric impacts around the  
33 world. As a natural experiment, it can be used to test our understanding of how the lower  
34 atmosphere affects space weather. Researchers are only beginning to document the chain  
35 of events post-eruption, and this paper focuses on its impact on the generator that drives  
36 electric fields in near-Earth space, a key part of space weather. This generator is driven  
37 by the atmosphere pushing charged particles across Earth’s magnetic field. This usually  
38 creates a strong eastward current above the equator. When the Swarm A satellite co-  
39 coincided with the wave from Tonga, it observed that this current strengthened dramati-  
40 cally, then reversed. Although reversals are not unusual, this was the strongest rever-  
41 sal observed by Swarm since its 2013 launch, except for one large geomagnetic storm in  
42 2015. Another satellite, the Ionospheric Connection Explorer, was luckily at the right  
43 time and place to observe related motions of the upper atmosphere, which were similarly  
44 extreme. These observations are shown to be consistent with our theoretical understand-  
45 ing of the generator. This study is important because it represents a critical test of atmosphere-  
46 space interactions and implies that the Tonga eruption caused a major space weather  
47 event.

**48 1 Introduction**

49 Isolated disturbances such as earthquakes, tsunamis, and solar eclipses, as well as  
50 explosions from volcanoes, nuclear detonations, and meteor air bursts can offer discrete  
51 tests for models of atmosphere-ionosphere coupling and variability (Aryal et al., 2020;  
52 Astafyeva, 2019; Inchin, Snively, et al., 2020; Li et al., 2021; Zettergren & Snively, 2019).  
53 The Hunga Tonga-Hunga Ha’apai (hereafter Tonga) volcanic eruption on 15 Jan 2022  
54 generated atmospheric disturbances from the ground to the ionosphere (Adam, 2022; Wright  
55 et al., 2022). A typical wave mode excited by impulsive events in the lower atmosphere  
56 is the Lamb wave, an acoustic nondispersive edge wave (Garrett, 1969; Bretherton, 1969;  
57 Nishida et al., 2014). The study by Wright et al. (2022) presented data from the tropo-  
58 sphere, stratosphere, and mesosphere, showing a coherent wave propagating at 318 m/s  
59 around the globe multiple times, identified as a Lamb wave. Although the Lamb wave  
60 propagates in the troposphere, energy leaks into higher altitudes, exciting other wave modes,  
61 in which the amplitude of wind, temperature, and pressure fluctuations can grow with  
62 altitude (Nishida et al., 2014; Wright et al., 2022). As such, the ionosphere, readily ob-  
63 served by ground-based instruments, can function as a sensitive monitor of atmospheric  
64 disturbances.

65 Indeed, initial Total Electron Content (TEC) observations have reported Travel-  
66 ling Ionospheric Disturbances (TIDs) propagating globally for many hours and even days

67 after the Tonga eruption (Aa et al., 2022; Lin et al., 2022; Themens et al., 2022; Zhang  
 68 et al., 2022). Estimates of the horizontal wavelength of TIDs in the far field (i.e., at dis-  
 69 tances >3000 km from the eruption) range from 300 to 1000 km (Wright et al., 2022;  
 70 Zhang et al., 2022). Additionally, Soares et al. (2022) reported oscillations of the geo-  
 71 magnetic field observed by a ground-based magnetometer 835 km from Tonga, which are  
 72 attributed to short-period modulation (3-5 min) of ionospheric currents. No studies have  
 73 yet reported data connecting the homosphere with these ionospheric signatures. Wright  
 74 et al. (2022) presented ionospheric TIDs with phase speeds, horizontal wavelengths, and  
 75 arrival times inconsistent with the Lamb wave, speculating that the observed TEC sig-  
 76 natures likely arrived by indirect paths from Tonga.

77 The mechanisms through which signals from the lower atmosphere are transmit-  
 78 ted and create observable effects in the ionosphere are numerous, and understanding their  
 79 complex interplay is critical for interpreting and predicting ionospheric signals. These  
 80 mechanisms include those resulting from direct propagation of the wave or waves to iono-  
 81 spheric  $F$ -region heights, modifying ion drag and/or plasma loss rates. Another mech-  
 82 anism is indirect, mediated by electric fields resulting from the neutral wind dynamo,  
 83 which can carry signatures along magnetic field lines from the  $E$  region to the  $F$  region.  
 84 The signal can also be transmitted to the opposite hemisphere, which has been proposed  
 85 to explain the appearance of TIDs over Japan ahead of the Lamb wave (Lin et al., 2022).  
 86 Conjugate effects were also suggested by Themens et al. (2022). In this study we report  
 87 on two aspects of the Tonga disturbance: neutral winds and ionospheric dynamo signa-  
 88 tures.

89 Specifically, we report extreme perturbations in the equatorial electrojet (EEJ) ob-  
 90 served by Swarm and extreme perturbations in neutral winds from 90 to 300 km alti-  
 91 tude observed by the Michelson Interferometer for Global High-resolution Thermospheric  
 92 Imaging (MIGHTI) on the Ionospheric Connection Explorer (ICON) (Immel et al., 2018).  
 93 The EEJ is an intense band of zonal electric current confined near the magnetic equa-  
 94 tor flowing in the daytime between  $\sim$ 90 and 120 km altitude (Yamazaki & Maute, 2017,  
 95 and references therein). Variations in the EEJ closely track those of the equatorial zonal  
 96 electric field (i.e., vertical plasma drift) which has widespread effects on the equatorial  
 97 ionosphere by modifying the production-loss-transport balance. Typically the EEJ flows  
 98 eastward, associated with an eastward zonal electric field, upward drift, and enhanced  
 99 equatorial fountain effect, but sometimes the EEJ flows westward, associated with the  
 100 opposite ionospheric conditions. In the absence of direct solar insolation, the EEJ dis-  
 101 appears at night.

102 ICON and Swarm have been operating simultaneously since ICON’s launch in 2019,  
 103 offering an unprecedented observational capability for studies related to the ionospheric  
 104 dynamo. On 15 Jan 2022, their orbits were unusually well-synced to provide complemen-  
 105 tary observations of the Tonga signature, as discussed below. This study does not at-  
 106 tempt to quantify properties or classifications of the waves excited by the Tonga explo-  
 107 sion, which will undoubtedly be a focus of future investigations. However, the unique  
 108 opportunity created by coincident observations of the neutral wind by MIGHTI and iono-  
 109 spheric currents by Swarm allows us to directly study the impact of these waves on the  
 110 ionospheric dynamo, which we report here. In addition, four magnetometer sites are uti-  
 111 lized to provide a ground-based perspective on the EEJ variation.

## 112 2 Data sources

### 113 2.1 ICON-MIGHTI neutral winds

114 This study uses neutral wind data from the MIGHTI instrument on the ICON space-  
 115 craft, which is in a  $27^\circ$  inclination orbit. Neutral wind profiles (ICON data product 2.2  
 116 v04) from 90 to 300 km altitude are derived by remote limb observations of the Doppler

117 shift of naturally occurring green 557.7 nm and red 630.0 nm (OI) airglow emissions. We  
 118 only use wind data from the dayside, which is when the airglow layers are wide enough  
 119 to permit wind estimates spanning the full 90 to 300 km altitude range, which fully en-  
 120 compasses the dynamo region. Nightside data have a gap from  $\sim 109$  to 210 km. Data  
 121 are available from the green and red channels with vertical sampling of  $\sim 3$  km and  $\sim 10$   
 122 km, respectively. For this study, to improve statistics especially in regions of dim airglow,  
 123 we preprocess the green channel by binning vertically by a factor of 2, yielding  $\sim 6$  km  
 124 sampling. Below 180 km altitude, we use samples from the green channel, and above 180  
 125 km, we use samples from the red channel. Data above 180 km have a negligible impact  
 126 on our conclusions regarding the wind dynamo but are included because the volcanogenic  
 127 waves are clearly evident up to 300 km. The reader is referred to previous publications  
 128 for more information on the MIGHTI instrument (Englert et al., 2017), processing al-  
 129 gorithms (Harding et al., 2017), and on-orbit validation (Harding et al., 2021; Makela  
 130 et al., 2021).

131 Although the focus is on two orbits on 15 Jan 2022, we also make use of the en-  
 132 tire dataset for background statistics. Specifically, we use all MIGHTI profiles from the  
 133 start of the mission until 14 Jan 2022 for which the variable “Wind\_Quality” is equal  
 134 to 1 (i.e., highest quality, 1,086,830 profiles in total). To generate these statistics, in ad-  
 135 dition to the altitude binning discussed above, the data were preprocessed with a 5-sample  
 136 median filter in time to remove outliers. Data obtained during geomagnetic storms are  
 137 included in these statistics. Statistics are presented in terms of percentiles; for example,  
 138 the 90% level for zonal wind represents a value such that 10% of samples have a zonal  
 139 wind larger than that level.

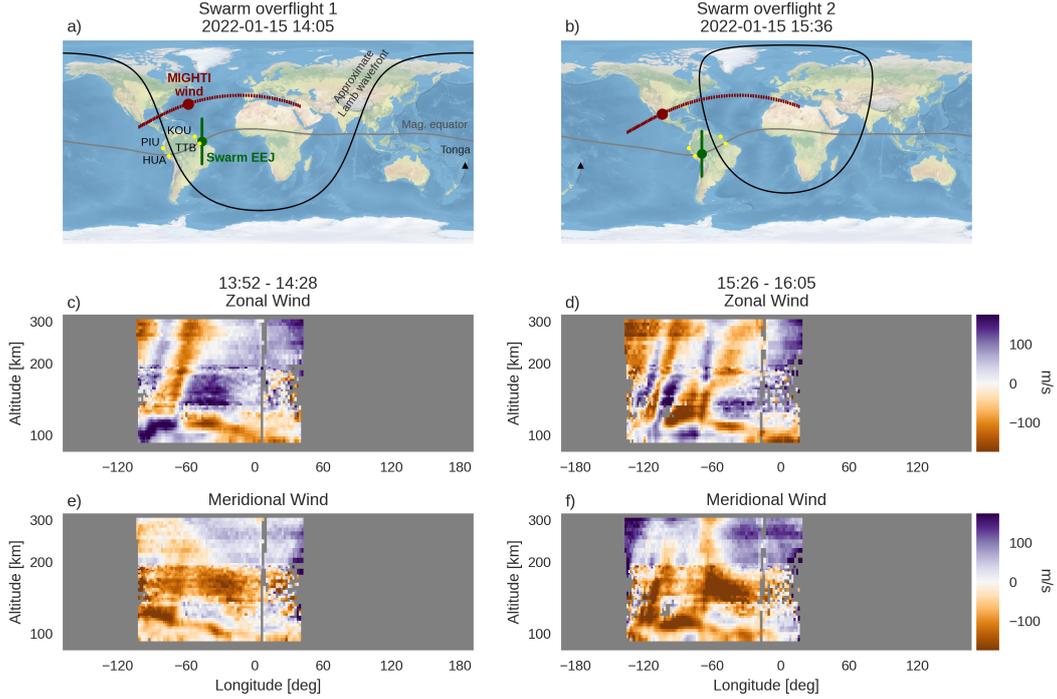
## 140 2.2 Swarm A EEJ current

141 The Swarm constellation comprises three satellites in near-polar orbits. In this study  
 142 we use EEJ intensity estimates from one spacecraft, Swarm A, which flies at an altitude  
 143 of  $\sim 440$  km with an inclination of  $87.4^\circ$ . Latitude-dependent height-integrated EEJ in-  
 144 tensity is provided by the Swarm Level 2 Product EEJ (Eastward Electric Field) (Alken  
 145 et al., 2013). The EEJ current is estimated from magnetometer measurements during  
 146 every dayside overflight of the magnetic equator (Alken, 2020). Ground-based valida-  
 147 tion is discussed by Alken et al. (2015).

148 In a manner analogous to the wind analysis, background statistics are calculated  
 149 for context, using the entire available dataset. Specifically, we use the version 0204 dataset  
 150 spanning 25 Nov 2013 to 14 Jan 2022. We first preprocess the EEJ data to remove non-  
 151 physical current distributions. These outliers are identified by computing the the total  
 152 “off-peak current” for each overflight (defined as the root-mean-square of currents pole-  
 153 ward of 5 deg quasidipole latitude). Overflights are removed if the off-peak current is larger  
 154 than 100 times the interquartile range of all the overflights (i.e., 75th percentile minus  
 155 25th percentile). This removes 25 overflights which are, by visual inspection, clear non-  
 156 physical outliers. The 45,184 remaining overflights are used in the statistics below. All  
 157 data on 15 Jan 2022 remain after this preprocessing step.

## 158 2.3 Ground-based magnetometers

159 We also use ground-based magnetometer data to support the interpretation of the  
 160 EEJ behavior on 15 Jan 2022. The intensity of the EEJ can be estimated using the hor-  
 161 izontal (H) component of the geomagnetic field observed at two stations, one being lo-  
 162 cated at the magnetic equator and the other located about the same longitude but out-  
 163 side the EEJ band (Anderson et al., 2004). The difference in H ( $\Delta H$ ) at the two stations,  
 164 after subtracting the nighttime baseline, represents the EEJ intensity. We use data from  
 165 Huancayo (HUA,  $12.0^\circ\text{S}$ ,  $75.3^\circ\text{W}$ ) and Piura (PIU,  $5.2^\circ\text{S}$ ,  $80.6^\circ\text{W}$ ) for the Peruvian sec-  
 166 tor, and Tatuoca (TTB,  $1.2^\circ\text{S}$ ,  $48.5^\circ\text{W}$ ) and Kourou (KOU,  $5.2^\circ\text{N}$ ,  $52.7^\circ\text{W}$ ) for the Brazil-



**Figure 1.** (a,b) Two consecutive orbits of ICON-MIGHTI dayside passes and Swarm equator overflights. The red and green dots indicate the location of the MIGHTI wind observation and Swarm EEJ observation at the respective time. The red and green lines indicate the observation locations just before and after this time. The dashed appearance of the red line denotes individual MIGHTI wind profile locations. The black line indicates the location of a notional radially propagating wavefront moving at 318 m/s that originated in Tonga at 04:28 UT (Wright et al., 2022). Although MIGHTI samples the wind  $\sim 35$  degrees farther north than the EEJ, the two sample the same great-circle distance from Tonga with a delay of several minutes. The yellow circles indicate the four ground-based magnetometer sites. (c,d) Daytime zonal wind profiles (positive eastward) measured by MIGHTI corresponding to the longitudes in the map above. Only daytime data are included. The small data gap near  $0^\circ$  longitude is caused by the quality control algorithm masking data when the moon is near the MIGHTI field of view. The altitude axis is in log scale to better display lower thermospheric winds. (e,f) Same as (c,d) for the meridional wind (positive northward).

167 ian sector. These stations are positioned to detect EEJ signatures in the vicinity of the  
 168 Swarm observations.

### 169 3 Results and Discussion

#### 170 3.1 Data selection and observational geometry

171 The Tonga volcanic eruption occurred around 04:15 UT on 15 Jan 2022, near local  
 172 sunset. Since the tropospheric sound speed is slower than the Earth's rotation at these  
 173 latitudes, the Lamb wave was mostly contained to the dusk and nighttime sectors for  
 174 the first several hours after the eruption. In this study we focus on thermosphere-ionosphere  
 175 signatures once the wave reaches the dayside, where ionospheric currents are strongest.  
 176 According to the parameters reported by Wright et al. (2022) (318 m/s phase speed orig-

177 inating in Tonga at 04:28 UT), the lower atmospheric Lamb wave reached the dayside  
 178 around 13 UT at low/mid-latitudes. Amores et al. (2022) provide further information  
 179 on the Lamb wave propagation, including a numerical simulation which agrees with the  
 180 timing used here. The Lamb wave entered the dayside in the American longitude sec-  
 181 tor. Serendipitously, Swarm A overflights occurred in this sector at 14:05 UT and 15:36  
 182 UT. In this study we utilize data from these two orbits and the corresponding orbits of  
 183 ICON, which samples all longitudes every orbit, albeit at different latitudes.

184 The two orbits of ICON-MIGHTI data are shown in Figure 1, an animated ver-  
 185 sion of which can be found in the Supporting Information (Movie S1). For context, we  
 186 show a reference wavefront using the Lamb wave parameters reported by Wright et al.  
 187 (2022). Given the close alignment between ICON and Swarm, these parameters are not  
 188 important for our conclusions, and similar parameters (e.g., a 310 m/s wavefront orig-  
 189 inating at 04:15 UT) do not change the interpretation. On the first orbit (panel a), Swarm  
 190 A crossed the equator and measured the EEJ at a location roughly 3000 km ahead of  
 191 the 318 m/s wavefront. At the time Swarm A measured the EEJ at the equator, MIGHTI  
 192 sampled the wind  $\sim 35$  degrees farther north but at a similar great-circle distance from  
 193 Tonga.

194 The next orbit is shown to the right (panel b). On this orbit, Swarm A measured  
 195 the EEJ at a location roughly 1500 km behind the assumed 318 m/s wavefront. At the  
 196 time of the Swarm A overflight, MIGHTI samples the wind roughly 3500 km behind the  
 197 wavefront, but reached the same great-circle distance as Swarm A 5 minutes later (15:41  
 198 UT).

199 In both orbits, the MIGHTI data (panels c, d, e, and f) show large wind fluctua-  
 200 tions, vertical shears, and coherent wave structures spanning at least 110–300 km, both  
 201 ahead of and behind the 318 m/s wavefront. Above 120 km, the horizontal wavelengths  
 202 of the wave structures are estimated by visual inspection to be 3000–5000 km, more than  
 203 three times as large as the horizontal wavelengths reported in TEC observations in the  
 204 far field. It is apparent from these observations that the thermospheric signatures of this  
 205 event are complex and likely not explained by a single wave mode. We do not comment  
 206 further on the wind features in this paper, but instead we focus on their impact on iono-  
 207 spheric currents in the next section.

208 In the meridional wind, a discontinuity is evident at 180 km where data from the  
 209 green and red channels are spliced together. This is a known artifact caused by a drift  
 210 in the zero-wind reference in the MIGHTI v04 dataset; a new calibration has been im-  
 211 plemented and will correct this bias in the next version of the MIGHTI wind dataset.  
 212 This discontinuity does not affect the qualitative conclusions of this study, which focus  
 213 on the zonal wind, nor the quantitative calculations in this study, which uses perturba-  
 214 tions from the mean, as described below.

### 215 **3.2 Comparison between MIGHTI winds and Swarm EEJ**

216 Figure 2 compares the Swarm A EEJ observations with the MIGHTI wind obser-  
 217 vations on these two orbits. On the first orbit (panel a), Swarm A observed an extremely  
 218 strong eastward EEJ (0.22 A/m). This represents the strongest EEJ observed by Swarm  
 219 A since 2017, and the 19th strongest overall (stronger than 99.96% of all observations  
 220 in the Swarm A dataset, which started in 2013). On the next orbit (panel b), Swarm A  
 221 observes an extreme westward EEJ (-0.17 A/m), often referred to as a counter-electrojet.  
 222 Except for three overflights during the 22-23 June 2015 geomagnetic storm, this repre-  
 223 sents the strongest westward EEJ in the Swarm A dataset. Swarm A data from earlier  
 224 orbits on this day do not show variations above the 90% level. Also shown are statistics  
 225 computed from all EEJ observations from the start of the mission until 14 Jan 2022. The  
 226 black line is the median, the dark gray shaded region is the interquartile range (25–75%),  
 227 and other percentile ranges are shown in lighter gray. Although Swarm B is not included

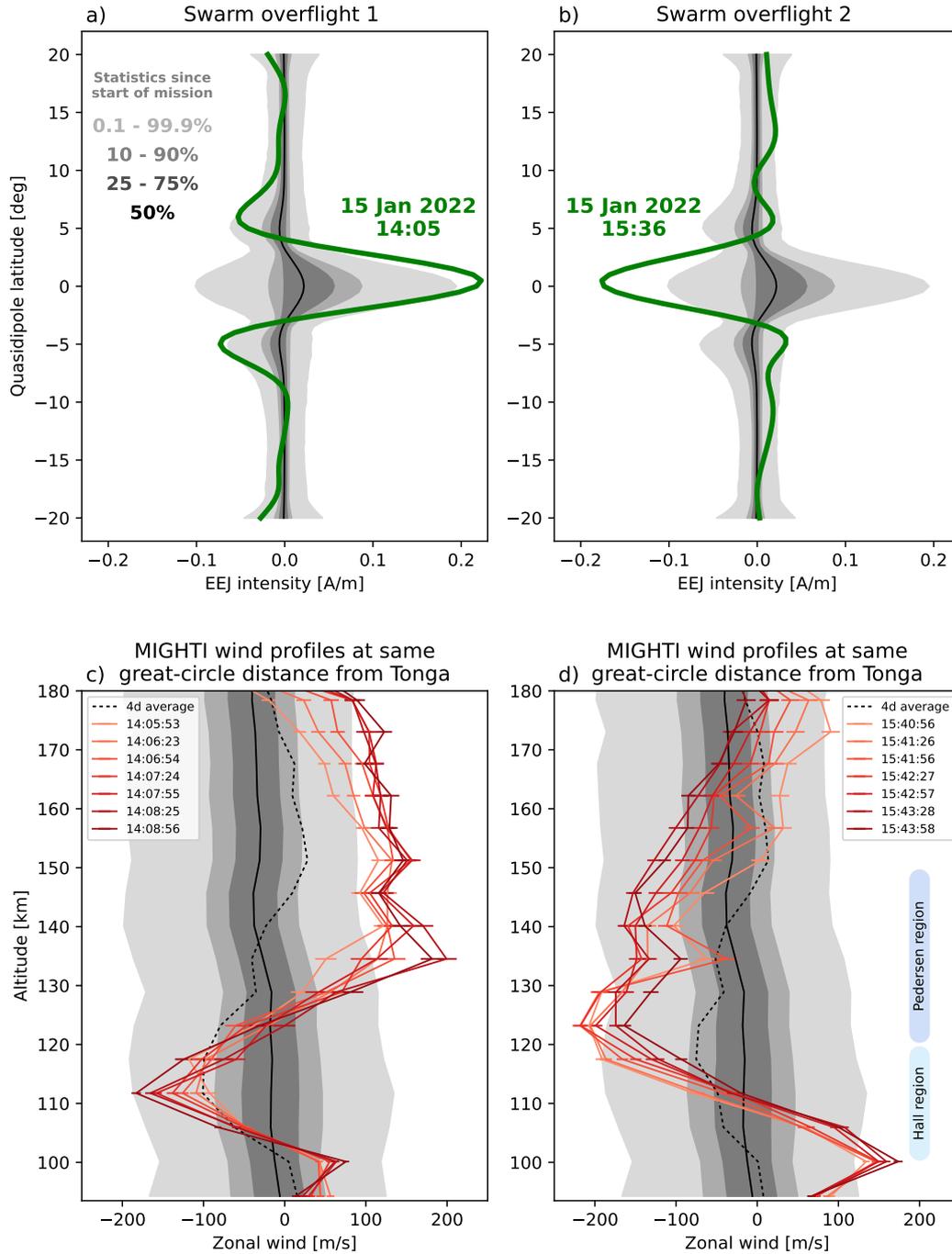
228 in this quantitative analysis, data from Swarm B also show a large positive EEJ (0.20  
 229 A/m) followed by a large negative EEJ (-0.14 A/m) on these two orbits (not shown). Swarm  
 230 C flies in a side-by-side configuration to Swarm A, and recorded similar measurements  
 231 on these two orbits (0.22 A/m and -0.17 A/m respectively, not shown).

232 The bottom of Figure 2 (panels c and d) shows the MIGHTI zonal winds correspond-  
 233 ing to the Swarm A EEJ observations, compared with background variability shown with  
 234 statistical ranges in gray, analogous to panels a and b. Although the meridional wind  
 235 fluctuations are in some cases quite significant, we focus on zonal winds because (1) merid-  
 236 ional winds are nearly parallel to the magnetic field at the equator and are not expected  
 237 to strongly influence the EEJ, and (2) the wave is propagating nearly zonally in this re-  
 238 gion. The dominant large-scale signature of the wave is therefore expected to be in the  
 239 zonal wind. A separate analysis was conducted where the zonal and meridional winds  
 240 were combined to calculate the radial wind perturbation in the direction away from Tonga.  
 241 However, this yielded identical conclusions and was more complicated to compare quan-  
 242 titatively with background statistics.

243 Zonal wind profiles (shown in red) are chosen such that their distance from Tonga  
 244 is identical to Swarm A’s distance when it crossed the magnetic equator. Insofar as the  
 245 wave can be assumed to propagate concentrically, this is a proxy for the neutral wind  
 246 fluctuations in the equatorial region during the Swarm A overflight. The various pro-  
 247 files in Figure 2 are chosen to provide exact alignment for assumed wave velocities rang-  
 248 ing from 300 to 330 m/s. The qualitative similarity of these profiles suggests that this  
 249 procedure to align the Swarm and MIGHTI observations is not significantly sensitive to  
 250 the assumed wave velocity, a consequence of the fortunate timing of the two observations.  
 251 The temporal offsets required are 0-4 minutes (overflight 1) and 5-8 minutes (overflight  
 252 2), a time scale that is not likely of importance for the large-scale waves observable by  
 253 MIGHTI. Furthermore, it is the same magnitude as the assumption of temporal persis-  
 254 tence used to produce the vector wind estimate by combining the data from the two MIGHTI  
 255 sensors (5-9 minutes) (Harding et al., 2017). A possibly non-negligible uncertainty in  
 256 this procedure is the assumption of concentric wave propagation, as the two observations  
 257 sample along different wave azimuths from Tonga, separated by 5–32 degrees.

258 These zonal wind profiles are a superposition of the volcanogenic waves and the back-  
 259 ground thermospheric state upon which they propagate. The black dashed line is an es-  
 260 timate of that background state, computed from an average of profiles at nearly the same  
 261 local time ( $\pm 1$  hr) and longitude ( $\pm 24$  deg, the amount of Earth rotation in one ICON  
 262 orbit) as the profiles shown. This average is generated using the 4 previous days (Jan  
 263 11-14), over which time the sampled latitude changes by no more than 10 deg. It is thus  
 264 an estimate of the contribution from background migrating and non-migrating tides and  
 265 planetary waves with periods  $\gtrsim 8$  days. However, there may also be contributions to the  
 266 background from short-term tidal variability, short-period planetary waves like the quasi-  
 267 two-day wave, and geomagnetic activity, which are difficult to comprehensively quan-  
 268 tify from a single observatory.

269 The wind profiles on both orbits are extreme, showing values comparable with, or  
 270 stronger than, the 0.1% and 99.9% levels. We describe the wind profile in terms of two  
 271 regions: The “Hall region,” ( $\sim 100$ – $120$  km) where the Hall conductivity is large and dom-  
 272 inant, and the “Pedersen region,” ( $\sim 120$ – $150$  km) where the Pedersen conductivity is large  
 273 and dominant. In reality the Hall and Pedersen conductivities are nonzero over larger  
 274 altitude ranges, and there is a significant overlap region in which they are both large;  
 275 however, this description is useful to connect with theoretical arguments below. On or-  
 276 bit 1, when the EEJ is strongly eastward, MIGHTI observes a westward wind in the Hall  
 277 region, which is not unusual compared to the background profile. However, there is also  
 278 a strong eastward wind in the Pedersen region which exceeds the 99.9% level. Indeed,  
 279 this represents the strongest wind observed at  $\sim 140$  km since the start of the mission.  
 280 On orbit 2, when the EEJ is strongly westward, MIGHTI observes an eastward wind,



**Figure 2.** (a,b) Swarm A equatorial electrojet (EEJ) observations during the two orbits shown in Figure 1. (c) MIGHTI wind profiles during the ICON orbit contemporaneous with the Swarm overflight in panel a. (d) Same, except one orbit later, contemporaneous with the Swarm overflight in panel b. Specific wind profiles were chosen to correspond to samples at the same great-circle distance from Tonga as the Swarm observation, for various assumed wave velocities from 300 to 330 m/s. Errorbars represent the uncertainty reported in the v04 wind product, which accounts for shot, read, and dark noise in the detector (i.e., precision). In all 4 panels, the gray shaded areas represent background statistics computed from percentiles of the entire dataset from the start of each mission until 14 Jan 2022, comprising 45,184 Swarm A overflights since 25 Nov 2013 and 1,086,830 MIGHTI profiles since 6 Dec 2019. The black solid line is the median. The black dashed line represents the background wind as computed from four previous days of data at similar local times and longitudes (see [text](#)).

281 peaking around 100 km in the lower Hall region, and a westward wind above  $\sim 110$  km,  
 282 which spans the upper Hall region and the Pedersen region. This profile is unusual relative  
 283 to the background wind and exceeds the 99.9% level at some altitudes.

284 This correspondence between the EEJ and neutral winds is consistent with the relationship  
 285 developed by Yamazaki et al. (2014) and Yamazaki et al. (2021). The early theoretical literature  
 286 on the EEJ suggested that while height-varying local winds influence the currents outside the EEJ,  
 287 they are not expected to have a significant influence on the EEJ itself, because it is dominated by  
 288 the influence of the global zonal electric field (Richmond, 1973). However, the modeling study by  
 289 Yamazaki et al. (2014) predicted that winds should have a significant role and that the EEJ should  
 290 be negatively correlated with Hall-region zonal winds and positively correlated with Pedersen-region  
 291 winds. This was observationally confirmed with the availability of concurrent MIGHTI and Swarm  
 292 observations by Yamazaki et al. (2021). The implicated mechanism is local generation of electric  
 293 fields which was not considered explicitly in the early (pre-2000) literature:  
 294 (1) in the Hall region, an eastward wind drives eastward current, which generates a westward  
 295 electric field; (2) in the Pedersen region, an eastward wind drives upward current, which  
 296 generates a downward electric field. At the footpoint of this field line, which lies in the Hall  
 297 region, the westward currents driven by this electric field will generate an eastward electric  
 298 field. Since the EEJ current flows in the Hall region, this latter case is a noteworthy example  
 299 of winds outside the EEJ perturbing currents in the EEJ.  
 300

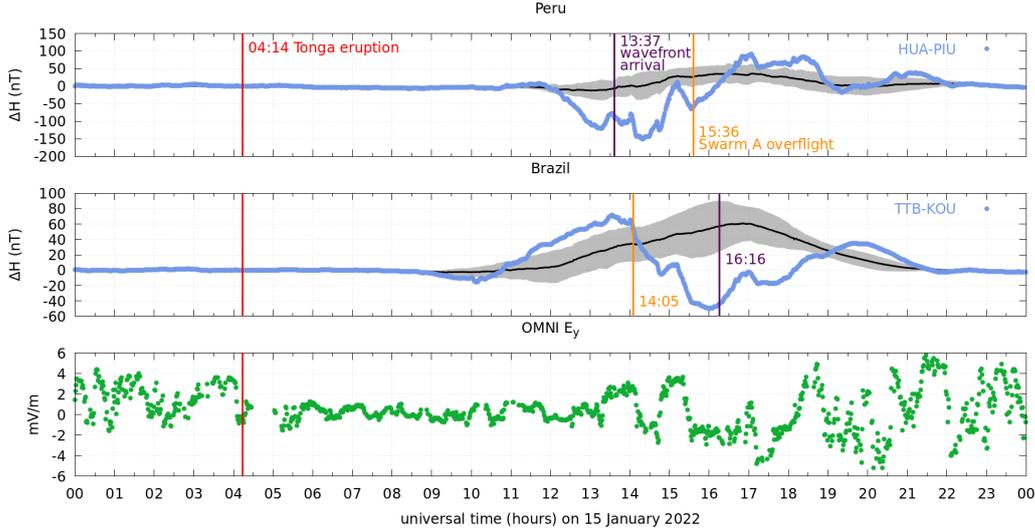
301 In orbit 1, the strong westward Hall-region wind and strong eastward Pedersen-region  
 302 wind is expected to cause a strong eastward EEJ through the Yamazaki et al. (2014) relationship.  
 303 In orbit 2, the Hall-region wind is eastward below 110 km and westward above 110 km, which is  
 304 expected to yield minimal total forcing in the Hall region. However, the Pedersen-region wind  
 305 is strongly westward, which is expected to cause a strong westward EEJ. In both cases, the Swarm  
 306 observations match the expectation. This result confirms the Yamazaki relationship holds under  
 307 extreme conditions. More interestingly, because the Hall-region effect is small in orbit 2, the  
 308 EEJ is apparently driven mostly by winds at higher altitudes, confirming the importance of  
 309 nonlocal wind driving of the EEJ. The current paths that regulate this control deserve further  
 310 inquiry, both observationally and theoretically.  
 311

### 312 3.3 Ground-based magnetometer data

313 In this section we report EEJ observations from two pairs of magnetometers located near the  
 314 Swarm overflights (see Figure 1). The observations are shown in the first two panels of Figure  
 315 3 using blue lines. The black line shows the monthly mean, and the gray shaded area represents  
 316 1 standard deviation (i.e.,  $1\sigma$ ) above and below the mean.

317 The HUA-PIU pair in Peru observes a negative  $\Delta H$  (corresponding to a westward EEJ) beginning  
 318 around 12 UT, lasting until just after 16 UT (except for one brief period of weak eastward EEJ  
 319 near 15 UT). The TTB-KOU pair in Brazil observes an eastward EEJ until  $\sim 15$  UT, followed by  
 320 a period of westward EEJ until 18 UT. Superimposed on these broad patterns are shorter, 1-hour  
 321 scale features which are discussed in the next section.  
 322

323 The broad features and relative timing seen in the magnetometer data are qualitatively  
 324 consistent with the Swarm observations. Namely, a negative disturbance is first seen over Peru,  
 325 then over Brazil 2-4 hours later, consistent with eastward propagation. The presence of 1-hour-scale  
 326 fluctuations and the lack of EEJ before sunrise makes it difficult to estimate the relative timing  
 327 with greater accuracy. On the first Swarm overflight at 14:05 UT, the nearby Brazilian pair  
 328 observes an eastward EEJ which is  $1\sigma$  or less above the climatology. On the second Swarm  
 329 overflight at 15:36 UT, the Peruvian pair observes a  $\sim 2.5\sigma$  extreme westward EEJ. Especially  
 330 for the positive EEJ on overflight 1, the fluctuations seen by the ground-based magnetometers  
 331 are not as extreme as the Swarm ob-



**Figure 3.** (top) Ground-based magnetometer EEJ intensity estimates over Peru on 15 Jan 2022 computed by subtracting PIU data (off-equator) data from HUA data (on-equator), shown in blue. The monthly mean is in black and  $\pm 1$  standard deviation range is in gray. The arrival time of a reference 318 m/s Lamb wavefront (purple line) and time of Swarm overflight (yellow line) are also shown. (middle) Same as top, but for Brazil (TTB - KOU). (bottom) Interplanetary eastward electric field from the OMNI database.

332 servations. Although the cause of this is unknown, it could be due to the ground-based  
 333 magnetometers being slightly offset from the magnetic equator. In January 2022, TTB  
 334 and HUA were  $2.2^\circ$  and  $0.8^\circ$  off the magnetic equator according to the CHAOS 7.8 model  
 335 (Finlay et al., 2020).

336 The different temporal patterns in Peru and Brazil confirm that the fluctuations  
 337 observed by Swarm A are not purely spatial but also temporal. The ground-based mag-  
 338 netometer data suggest that the most extreme EEJ activity may have been at locations  
 339 and times not sampled by Swarm A (e.g., over Brazil at 16 UT). Future work utilizing  
 340 the global network of magnetometers could help elucidate the evolution of global cur-  
 341 rents during this event.

342 The magnetometer data show disturbances before the arrival of the 318 m/s Lamb  
 343 wavefront (e.g., the negative  $\Delta H$  in Peru at 13 UT, and the positive and negative  $\Delta H$   
 344 in Brazil before 16 UT). This is consistent with the Swarm A observations ahead of the  
 345 wavefront at 14:05 UT (Figures 1a and 2a) and the MIGHTI observations on the first  
 346 orbit (Figure 1c, eastward of  $-60^\circ$  longitude). It is likely that the thermospheric response  
 347 to the eruption is not as simple as the Lamb wave observed in the lower atmosphere, due  
 348 to the effects of nonlinear evolution, dispersion, self-acceleration, and secondary wave gen-  
 349 eration, among others. Although no numerical models have yet simulated the upper at-  
 350 mospheric response to the Tonga Lamb wave, Inchin, Heale, et al. (2020) provide a dis-  
 351 cussion on these processes using a first-principles model of the thermospheric signature  
 352 of tsunamis.

### 353 3.4 Geomagnetic storm effects

354 A moderate geomagnetic storm began on 14 Jan 2022; the Tonga eruption and sub-  
 355 sequent wave propagation occurred during the recovery phase. It is thus important to  
 356 distinguish the signatures caused by the Tonga eruption from the effects of the storm.

357 The EEJ is known to be modified by electric fields penetrating from the magnetosphere  
 358 and electric fields originating from the stormtime disturbance dynamo (Yamazaki & Maute,  
 359 2017, and references therein). First, we rule out penetration electric field effects.

360 Figure 3 (bottom panel) shows the interplanetary electric field (IEF) y-component  
 361 (dawn-to-dusk electric field) from OMNI data (King & Papitashvili, 2005). The data are  
 362 taken directly from the OMNI database, except they include a 17-minute delay to ac-  
 363 count for the delay between the bowshock and the ionosphere (Manoj et al., 2008). If  
 364 the penetration electric field were the main cause of the EEJ variations, we would ex-  
 365 pect to see strong correlations between the IEF and  $\Delta H$  in both longitude sectors. Quan-  
 366 titatively, the Pearson correlations between IEF and the deviations of  $\Delta H$  from the monthly  
 367 mean, (blue lines minus solid black lines in Figure 3), between 8 and 16 hr local time is  
 368 -0.02 for Peru (13 to 22 UT) and 0.35 for Brazil (11 to 20 UT). However, the fluctua-  
 369 tions observed in IEF appear to correlate with 1-hour-scale fluctuations observed at both  
 370 ground-based sites simultaneously (e.g., positive excursions at 15 UT and 17 UT, and  
 371 possibly at 13.5 UT). After filtering  $\Delta H$  and IEF to remove their 100-minute running  
 372 mean, the correlation increased to 0.60 (Peru) and 0.61 (Brazil). Thus, it is likely that  
 373 the 1-hour-scale fluctuations are caused in part by the penetration electric field, but the  
 374 larger, longer perturbations of interest here are not. Because of this, and because of the  
 375 consistency between the EEJ signatures and the neutral wind signatures, as discussed  
 376 above, we rule out the penetration electric field as the main cause of the extreme east-  
 377 ward and westward EEJ observed by Swarm.

378 With neutral winds established as the causative mechanism, it is important to rule  
 379 out geomagnetic activity as the cause of the extreme winds seen in Figure 2(c,d). It is  
 380 well known that the EEJ can be reversed by the disturbance dynamo, a consequence of  
 381 westward Coriolis forcing of neutral winds accelerated equatorward by auroral heating  
 382 (Yamazaki & Maute, 2017). According to the modeling study by Huang et al. (2005),  
 383 disturbance winds caused by a geomagnetic storm are mainly in the westward direction  
 384 at middle and low latitudes. MIGHTI observations show both eastward and westward  
 385 wind perturbations, which are different from the predicted pattern of the disturbance  
 386 winds. Also, storm-driven wind perturbations are predicted to be much greater at F-region  
 387 heights (above 150 km) than at E-region heights (below 150 km). MIGHTI observations  
 388 show large wind perturbations below 150 km (including an eastward perturbation at 100  
 389 km exceeding 100 m/s), which does not fit the classical picture of the disturbance winds.  
 390 Furthermore, the westward disturbance wind at mid and low latitudes is stronger dur-  
 391 ing nighttime than daytime. For instance, Xiong et al. (2015) showed that the average  
 392 westward disturbance wind at 20–50° latitude is less than 50 m/s for  $K_p > 4$  at F-region  
 393 heights during daytime, while it can exceed 100 m/s during nighttime. Thus, the geo-  
 394 magnetic storm is unlikely to be the main cause of the extreme daytime winds detected  
 395 by MIGHTI.

396 The simultaneous occurrence of the Lamb wave arrival, the EEJ signal, and the wind  
 397 signal, combined with the lack of any significant wind or EEJ signals before this time,  
 398 represents strong evidence to attribute the observed fluctuations to disturbances caused  
 399 by the Tonga eruption. Nevertheless, it is possible that high-latitude heating launched  
 400 traveling atmospheric disturbances during the recovery phase, and it is likely that the  
 401 longer-term circulation changes caused by the storm have changed the background con-  
 402 ditions upon which the Tonga signal is superimposed. It will be an interesting topic for  
 403 future modeling and observational studies to elucidate the interplay of geomagnetic storm  
 404 and volcanogenic effects on the thermosphere and ionosphere during this period.

## 405 4 Conclusion

406 The 15 Jan 2022 Tonga volcanic eruption caused extreme fluctuations in the iono-  
 407 spheric wind dynamo, as observed by Swarm and ICON. On the leading part of the ini-

408 tial disturbance, Swarm A observed an extreme (0.22 A/m) eastward EEJ at the 99.96th  
 409 percentile of typical variability, while ICON-MIGHTI observed a strong westward wind  
 410 in the Hall region and a strong eastward wind in the Pedersen region, a relationship con-  
 411 sistent with recent theoretical and observational studies. Relative to typical variability,  
 412 the wind perturbations (up to 200 m/s) were at or above the 99.9th percentile, and at  
 413 some altitudes (e.g.,  $\sim 140$  km) were the strongest observed since ICON's launch. On the  
 414 trailing part of the initial disturbance, Swarm A observed the strongest westward EEJ  
 415 (-0.17 A/m) of the mission, except for the 22-23 June 2015 geomagnetic storm, while ICON-  
 416 MIGHTI observed eastward and westward winds in the Hall region and a westward wind  
 417 in the Pedersen region, also at or above the 99.9th percentile.

418 The relationship between winds and EEJ is consistent with the theoretical and ob-  
 419 servational relationships of neutral wind driving developed by Yamazaki et al. (2014) and  
 420 Yamazaki et al. (2021). The westward electrojet in this case appears to be driven mostly  
 421 by westward winds in the Pedersen region. Since the EEJ flows in the Hall region, this  
 422 is a noteworthy example of nonlocal wind driving of the EEJ. The energy and current  
 423 paths involved in this relationship would be an interesting topic for further study.

424 Initial reports on the global ionosphere-thermosphere impacts of the Tonga erup-  
 425 tion have focused on small- and meso-scale (300–1000 km wavelength) waves seen in TEC  
 426 at amplitudes of at most a few TEC units, as well as geomagnetic fluctuations 835 km  
 427 away from and soon after the eruption. The MIGHTI and Swarm observations suggest  
 428 that modifications of the ionospheric dynamo were extreme relative to background vari-  
 429 ability, even after  $\sim 10$  hours and  $\sim 10,000$  km of wave propagation. This is expected to  
 430 have caused significant and observable redistributions of ionospheric plasma. As an ex-  
 431 ample of an enormous impulse function, the Tonga eruption may be a useful test for atmosphere-  
 432 ionosphere coupled models in extreme cases, and the neutral wind and EEJ current sig-  
 433 natures reported here could be useful targets.

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## 446 Open Research

447 ICON data can be retrieved from the ICON website (<https://icon.ssl.berkeley.edu/Data>). The European Space Agency (ESA) is gratefully acknowledged for provid-  
 448 ing Swarm data, which is available from the Swarm website ([https://swarm-diss.esa.int/#swarm%2FLevel2daily%2FEntire\\_mission\\_data%2FEEF%2FTMS%2FSat\\_A](https://swarm-diss.esa.int/#swarm%2FLevel2daily%2FEntire_mission_data%2FEEF%2FTMS%2FSat_A)).  
 449 The OMNI data were obtained from the GSFC/SPDF OMNIWeb interface at <https://omniweb.gsfc.nasa.gov>. Ground-based magnetometer data from HUA, TTB, and KOU  
 450 are available from <https://intermagnet.org/data-donnee/download-eng.php>. Data  
 451 from PIU is available from the LISN network ([http://lisn.igp.gob.pe/jdata/view/magnetometer/minute/piur/?itype=magnetometer&dtype=minute&daterange=2022%2F01%2F15+-+2022%2F01%2F15&networks=on&N\\_IGP=on&N\\_LISN=on&N\\_MAGDAS=on&countries=on&C\\_Argentina=on&C\\_Brasil=on&C\\_Colombia=on&C\\_Peru=on&stations=on&S\\_areq=](http://lisn.igp.gob.pe/jdata/view/magnetometer/minute/piur/?itype=magnetometer&dtype=minute&daterange=2022%2F01%2F15+-+2022%2F01%2F15&networks=on&N_IGP=on&N_LISN=on&N_MAGDAS=on&countries=on&C_Argentina=on&C_Brasil=on&C_Colombia=on&C_Peru=on&stations=on&S_areq=)

458 on&S\_leon=on&S\_cuib=on&S\_dejp=on&S\_huan=on&S\_jica=on&S\_ancm=on&S\_huam=on&S  
 459 \_icam=on&S\_nazc=on&S\_piur=on&bt\_view=piur&S\_saol=on&S\_tara=on).

## 460 References

- 461 Aa, E., Zhang, S.-R., Erickson, P. J., Vierinen, J., Coster, A. J., Goncharenko, L. P.,  
 462 ... Rideout, W. (2022). Significant equatorial plasma bubbles and global  
 463 ionospheric disturbances after the 2022 tonga volcano eruption. *Earth and*  
 464 *Space Science Open Archive*, 17. Retrieved from [https://doi.org/10.1002/](https://doi.org/10.1002/essoar.10510637.1)  
 465 [essoar.10510637.1](https://doi.org/10.1002/essoar.10510637.1) doi: 10.1002/essoar.10510637.1
- 466 Adam, D. (2022). Tonga volcano eruption created puzzling ripples in Earth's at-  
 467 mosphere. *Nature*, 601, 497. Retrieved from [https://www.nature.com/](https://www.nature.com/articles/d41586-022-00127-1)  
 468 [articles/d41586-022-00127-1](https://www.nature.com/articles/d41586-022-00127-1) doi: 10.1038/d41586-022-00127-1
- 469 Alken, P. (2020). Estimating currents and electric fields at low latitudes from  
 470 satellite magnetic measurements. In M. W. Dunlop & H. Lühr (Eds.), *Iono-*  
 471 *spheric multi-spacecraft analysis tools: Approaches for deriving ionospheric*  
 472 *parameters* (pp. 233–254). Cham: Springer International Publishing. Re-  
 473 trieved from [https://doi.org/10.1007/978-3-030-26732-2\\_11](https://doi.org/10.1007/978-3-030-26732-2_11) doi:  
 474 10.1007/978-3-030-26732-2\_11
- 475 Alken, P., Maus, S., Chulliat, A., Vigneron, P., Sirol, O., & Hulot, G. (2015). Swarm  
 476 equatorial electric field chain: First results. *Geophysical Research Letters*,  
 477 42(3), 673–680. doi: 10.1002/2014GL062658
- 478 Alken, P., Maus, S., Vigneron, P., Sirol, O., & Hulot, G. (2013). Swarm SCARF  
 479 equatorial electric field inversion chain. *Earth, Planets and Space*, 65(11),  
 480 1309–1317. doi: 10.5047/eps.2013.09.008
- 481 Amores, A., Monserrat, S., Marcos, M., Argüeso, D., Villalonga, J., Jordà, G., &  
 482 Gomis, D. (2022). Numerical simulation of atmospheric Lamb waves generated  
 483 by the 2022 Hunga-Tonga volcanic eruption. *Geophysical Research Letters*,  
 484 1–8. doi: 10.1029/2022gl098240
- 485 Anderson, D., Anghel, A., Chau, J., & Veliz, O. (2004). Daytime vertical  $E \times B$   
 486 drift velocities inferred from ground-based magnetometer observations at low  
 487 latitudes. *Space Weather*, 2(11), n/a-n/a. doi: 10.1029/2004sw000095
- 488 Aryal, S., Evans, J. S., Correira, J., Burns, A. G., Wang, W., Solomon, S. C., ...  
 489 Jee, G. (2020). First Global-Scale Synoptic Imaging of Solar Eclipse Effects in  
 490 the Thermosphere. *Journal of Geophysical Research: Space Physics*, 125(9).  
 491 doi: 10.1029/2020JA027789
- 492 Astafyeva, E. (2019). Ionospheric Detection of Natural Hazards. *Reviews of Geo-*  
 493 *physics*, 57(4), 1265–1288. doi: 10.1029/2019RG000668
- 494 Bretherton, F. P. (1969). Lamb waves in a nearly isothermal atmosphere. *Quarterly*  
 495 *Journal of the Royal Meteorological Society*, 95(406), 754–757. Retrieved from  
 496 <https://onlinelibrary.wiley.com/doi/10.1002/qj.49709540608> doi: 10  
 497 .1002/qj.49709540608
- 498 Englert, C. R., Harlander, J. M., Brown, C. M., Marr, K. D., Miller, I. J., Stump,  
 499 J. E., ... Immel, T. J. (2017). Michelson Interferometer for Global High-  
 500 Resolution Thermospheric Imaging (MIGHTI): Instrument Design and Cali-  
 501 bration. *Space Science Reviews*, 212(1-2), 553–584. Retrieved from [http://](http://dx.doi.org/10.1007/s11214-017-0358-4)  
 502 [dx.doi.org/10.1007/s11214-017-0358-4](http://dx.doi.org/10.1007/s11214-017-0358-4)  
 503 [http://link.springer.com/](http://link.springer.com/10.1007/s11214-017-0358-4)  
 504 10.1007/s11214-017-0358-4 doi: 10.1007/s11214-017-0358-4
- 505 Finlay, C. C., Kloss, C., Olsen, N., Hammer, M. D., Tøffner-Clausen, L., Grayver,  
 506 A., & Kuvshinov, A. (2020). The CHAOS-7 geomagnetic field model and  
 507 observed changes in the South Atlantic Anomaly. *Earth, Planets and Space*,  
 72(1), 1–31.
- 508 Garrett, C. J. (1969). Atmospheric edge waves. *Quarterly Journal of the Royal Me-*  
 509 *teorological Society*, 95(406), 731–753. doi: 10.1002/qj.49709540607
- 510 Harding, B. J., Chau, J. L., He, M., Englert, C. R., Harlander, J. M., Marr, K. D.,

- 511 ... Immel, T. J. (2021). Validation of ICON-MIGHTI Thermospheric Wind  
 512 Observations: 2. Green-Line Comparisons to Specular Meteor Radars. *Jour-*  
 513 *nal of Geophysical Research: Space Physics*, 126(3), 1–12. Retrieved from  
 514 <https://onlinelibrary.wiley.com/doi/10.1029/2020JA028947> doi:  
 515 10.1029/2020JA028947
- 516 Harding, B. J., Makela, J. J., Englert, C. R., Marr, K. D., Harlander, J. M., Eng-  
 517 land, S. L., & Immel, T. J. (2017). The MIGHTI Wind Retrieval Al-  
 518 gorithm: Description and Verification. *Space Science Reviews*, 212(1-  
 519 2), 585–600. Retrieved from [http://dx.doi.org/10.1007/s11214-017-](http://dx.doi.org/10.1007/s11214-017-0359-3)  
 520 [-0359-3](http://link.springer.com/10.1007/s11214-017-0359-3) [http://link.springer](http://link.springer.com/10.1007/s11214-017-0359-3)  
 521 [com/10.1007/s11214-017-0359-3](http://link.springer.com/10.1007/s11214-017-0359-3) doi:  
 10.1007/s11214-017-0359-3
- 522 Huang, C.-M., Richmond, A., & Chen, M.-Q. (2005). Theoretical effects of geomag-  
 523 netic activity on low-latitude ionospheric electric fields. *Journal of Geophysical*  
 524 *Research: Space Physics*, 110(A5).
- 525 Immel, T. J., England, S. L., Mende, S. B., Heelis, R. A., Englert, C. R., Edelstein,  
 526 J., ... Sirk, M. M. (2018). The Ionospheric Connection Explorer Mission:  
 527 Mission Goals and Design. *Space Science Reviews*, 214(1), 13. Retrieved from  
 528 <http://dx.doi.org/10.1007/s11214-017-0449-2> <http://link.springer>  
 529 [.com/10.1007/s11214-017-0449-2](http://link.springer.com/10.1007/s11214-017-0449-2) doi: 10.1007/s11214-017-0449-2
- 530 Inchin, P. A., Heale, C. J., Snively, J. B., & Zettergren, M. D. (2020). The Dynam-  
 531 ics of Nonlinear Atmospheric Acoustic-Gravity Waves Generated by Tsunamis  
 532 Over Realistic Bathymetry. *Journal of Geophysical Research: Space Physics*,  
 533 125(12), 1–18. doi: 10.1029/2020JA028309
- 534 Inchin, P. A., Snively, J. B., Zettergren, M. D., Komjathy, A., Verkhoglyadova,  
 535 O. P., & Tulasi Ram, S. (2020). Modeling of Ionospheric Responses to Atmo-  
 536 spheric Acoustic and Gravity Waves Driven by the 2015 Nepal Mw7.8 Gorkha  
 537 Earthquake. *Journal of Geophysical Research: Space Physics*, 125(4). doi:  
 538 10.1029/2019JA027200
- 539 King, J. H., & Papitashvili, N. E. (2005). Solar wind spatial scales in and  
 540 comparisons of hourly Wind and ACE plasma and magnetic field data.  
 541 *Journal of Geophysical Research: Space Physics*, 110(A2), 1–9. doi:  
 542 10.1029/2004JA010649
- 543 Li, R., Lei, J., & Dang, T. (2021). The Solar Eclipse Effects on the Up-  
 544 per Thermosphere. *Geophysical Research Letters*, 48(15), 1–10. doi:  
 545 10.1029/2021GL094749
- 546 Lin, J.-T., Rajesh, P. K., Lin, C. C. H., Chou, M.-Y., Liu, J.-Y., Yue, J., ... Kung,  
 547 M.-M. (2022). Rapid conjugate appearance of the giant ionospheric lamb wave  
 548 in the northern hemisphere after hunga-tonga volcano eruptions. *Earth and*  
 549 *Space Science Open Archive*, 27. Retrieved from [https://doi.org/10.1002/](https://doi.org/10.1002/essoar.10510440.2)  
 550 [essoar.10510440.2](https://doi.org/10.1002/essoar.10510440.2) doi: 10.1002/essoar.10510440.2
- 551 Makela, J. J., Baughman, M., Navarro, L. A., Harding, B. J., Englert, C. R., Har-  
 552 lander, J. M., ... Immel, T. J. (2021). Validation of ICON-MIGHTI Thermo-  
 553 spheric Wind Observations: 1. Nighttime Red-Line Ground-Based Fabry-Perot  
 554 Interferometers. *Journal of Geophysical Research: Space Physics*, 126(2),  
 555 1–29. Retrieved from [https://onlinelibrary.wiley.com/doi/10.1029/](https://onlinelibrary.wiley.com/doi/10.1029/2020JA028726)  
 556 [2020JA028726](https://onlinelibrary.wiley.com/doi/10.1029/2020JA028726) doi: 10.1029/2020JA028726
- 557 Manoj, C., Maus, S., Lühr, H., & Alken, P. (2008). Penetration character-  
 558 istics of the interplanetary electric field to the daytime equatorial iono-  
 559 sphere. *Journal of Geophysical Research: Space Physics*, 113(12), 1–11. doi:  
 560 10.1029/2008JA013381
- 561 Nishida, K., Kobayashi, N., & Fukao, Y. (2014). Background Lamb waves in the  
 562 Earth's atmosphere. *Geophysical Journal International*, 196(1), 312–316. Re-  
 563 trieved from <https://academic.oup.com/gji/article/196/1/312/594635>  
 564 doi: 10.1093/gji/ggt413
- 565 Richmond, A. D. (1973). Equatorial electrojet-I. Development of a model including

- 566 winds and instabilities. *Journal of Atmospheric and Terrestrial Physics*, 35(6),  
 567 1083–1103. doi: 10.1016/0021-9169(73)90007-X
- 568 Soares, G., Yamazaki, Y., & Matzka, J. (2022). Localized geomagnetic disturbance  
 569 due to ionospheric response to the hunga tonga eruption on january 15, 2022.  
 570 *Earth and Space Science Open Archive*, 12. Retrieved from [https://doi.org/](https://doi.org/10.1002/essoar.10510482.1)  
 571 10.1002/essoar.10510482.1 doi: 10.1002/essoar.10510482.1
- 572 Themens, D. R., Watson, C., Žagar, N., Vasylyevych, S., Elvidge, S., McCaffrey, A.,  
 573 ... Jayachandran, P. (2022). Global propagation of ionospheric disturbances  
 574 associated with the 2022 Tonga Volcanic Eruption. *Geophysical Research*  
 575 *Letters*. doi: 10.1029/2022gl098158
- 576 Wright, C., Hindley, N., Alexander, M. J., Barlow, M., Hoffmann, L., Mitchell,  
 577 C., ... et al. (2022). Tonga eruption triggered waves propagating glob-  
 578 ally from surface to edge of space. *Earth and Space Science Open Archive*,  
 579 23. Retrieved from <https://doi.org/10.1002/essoar.10510674.1> doi:  
 580 10.1002/essoar.10510674.1
- 581 Xiong, C., Lühr, H., & Fejer, B. G. (2015). Global features of the disturbance winds  
 582 during storm time deduced from CHAMP observations. *Journal of Geophysical*  
 583 *Research A: Space Physics*, 120(6), 5137–5150. doi: 10.1002/2015JA021302
- 584 Yamazaki, Y., Harding, B. J., Stolle, C., & Matzka, J. (2021). Neutral Wind Profiles  
 585 During Periods of Eastward and Westward Equatorial Electrojet. *Geophysical*  
 586 *Research Letters*, 48(11), 1–10. doi: 10.1029/2021GL093567
- 587 Yamazaki, Y., & Maute, A. (2017). Sq and EEJ—A Review on the Daily Varia-  
 588 tion of the Geomagnetic Field Caused by Ionospheric Dynamo Currents. *Space*  
 589 *Science Reviews*, 206(1-4), 299–405. Retrieved from [http://dx.doi.org/10](http://dx.doi.org/10.1007/s11214-016-0282-z)  
 590 [.1007/s11214-016-0282-z](http://dx.doi.org/10.1007/s11214-016-0282-z) doi: 10.1007/s11214-016-0282-z
- 591 Yamazaki, Y., Richmond, A. D., Maute, A., Liu, H.-L., Pedatella, N., & Sassi, F.  
 592 (2014). On the day-to-day variation of the equatorial electrojet during quiet  
 593 periods. *Journal of Geophysical Research: Space Physics*, 119(8), 6966–  
 594 6980. Retrieved from <http://doi.wiley.com/10.1002/2014JA020243> doi:  
 595 10.1002/2014JA020243
- 596 Zettergren, M. D., & Snively, J. B. (2019). Latitude and Longitude Dependence of  
 597 Ionospheric TEC and Magnetic Perturbations From Infrasonic-Acoustic Waves  
 598 Generated by Strong Seismic Events. *Geophysical Research Letters*, 46(3),  
 599 1132–1140. doi: 10.1029/2018GL081569
- 600 Zhang, S.-r., Vierinen, J., Aa, E., & Goncharenko, L. P. (2022). 2022 Tonga vol-  
 601 canic eruption induced global propagation of ionospheric disturbances via  
 602 Lamb waves. *Frontiers in Astronomy and Space Sciences*, 9, 1–15. doi:  
 603 10.3389/fspas.2022.871275