

1 **Energetic proton acceleration associated with Io’s footprint tail**

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30 **Key Points:**

- 31 • Juno’s likely crossing of Io’s Main Alfvén Wing (MAW) during PJ12 reveals evidence of
32 transverse ion acceleration
- 33 • Observations suggest wave-particle interactions with ion cyclotron waves as the favored
34 acceleration mechanism; however, Alfvén acceleration was not ruled out.
- 35 • Ion conics generated in Io’s footprint tail or near the MAW are more intense and
36 energetic than observed in other auroral regions

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43 **Abstract**

44 Observations of energetic charged particles associated with Io's footprint (IFP) tail, and likely
45 within or very near the Main Alfvén Wing, during Juno's 12th perijove (PJ) crossing show
46 evidence of intense proton acceleration by wave-particle heating. Measurements made by
47 Juno/JEDI reveal proton characteristics that include pitch angle distributions concentrated along
48 the upward loss cone, broad energy distributions that span ~50 keV to 1 MeV, highly structured
49 temporal/spatial variations in the particle intensities, and energy fluxes as high as ~100 mW/m².
50 Simultaneous measurements of the plasma waves and magnetic field suggest the presence of ion
51 cyclotron waves and transverse Alfvénic fluctuations. We interpret the proton observations as
52 upgoing conics likely accelerated via resonant interactions with ion cyclotron waves. These
53 observations represent the first measurements of ion conics associated with moon-magnetosphere
54 interactions, suggesting energetic ion acceleration plays a more important role in the IFP tail
55 region than previously considered.

56 **Plain-Language Summary**

57 NASA's Juno spacecraft orbits Jupiter's polar region and makes direct measurements of the
58 fields and particles that are responsible for creating Jupiter's powerful auroras. In this article, we
59 present new observations that show intense proton acceleration occurring at altitudes near the
60 auroral emissions created by the interaction between Jupiter's moon Io and the surrounding
61 plasma and magnetic field environment. These unique observations provide clues on how
62 particles are being accelerated and will help constrain particle acceleration theories.

63

64 **1. Introduction**

65 Juno's exploration of Jupiter's polar magnetosphere (Bagenal et al., 2017, Connerney et al.,
66 2017a) has given prominence to the "far-field" region of the Io-Jupiter interaction with new *in*
67 *situ* measurements of Io's footprint (IFP) tail auroral emissions. The far-field interaction
68 specifically refers to the electromagnetic coupling between Io and Jupiter's ionosphere. Decades
69 of remote observations have established that Io generates steady auroral emissions in the radio
70 (Bigg et al., 1964, Queinsec and Zarka, 1998, Zarka 2000), infrared (Connerney et al. 1993), and
71 ultraviolet wavelengths (Clarke et al., 1996). Furthermore, more recent HST observations and
72 analyses have characterized its auroral structuring (Bonfond et al. 2008) and correlated
73 brightness changes with Io's centrifugal latitude (Gérard et al. 2006). Previous flybys of Io from
74 the Voyager and Galileo missions mapped out the local Io-plasma interaction. Plasma and
75 energetic particle (e.g., Belcher et al. 1981, Frank et al. 1996, Williams et al. 1996, Gurnett et al.
76 1996) and magnetic field (e.g., Acuña et al. 1981, Kivelson et al. 1996) perturbations were
77 observed and consistent with theories of Alfvén wing model, but also have been discussed in the
78 context of a unipolar inductor model (e.g., Goldreich and Lynden-Bell, 1969; Neubauer 1980;
79 Goertz 1980; Gurnett & Goertz 1981; Bagenal 1983, Cray and Bagenal, 1997, Saur, 2004b).
80 However, how these processes are coupled to Jupiter's ionosphere, which may be the primary
81 place where the energy is dissipated, and how that energy is transferred to accelerating charged

82 particles in the auroral region is still not well understood (e.g., Saur et al. 2004a, Kivelson et al.
83 2004, Clarke et al., 2004) due the lack of *in situ* measurements in the auroral regions during these
84 epochs.

85

86 Recent analyses of the Juno magnetic field (Gershman et al., 2019), plasma wave (Sulaiman et
87 al., 2020) and the low-energy charged particle data (Szalay et al., 2018, 2020a, 2020b) almost
88 universally depict Alfvénic acceleration as a notable, if not dominant, electron acceleration
89 mechanism associated with the IFP tail. More specifically, Gershman et al. (2019) found
90 evidence of transverse magnetic field fluctuations consistent with strong magnetohydrodynamic
91 (MHD) turbulence that can supply $\sim 3,000$ mW/m² of Alfvénic Poynting flux near Io’s Main
92 Alfvén Wing (MAW). Similarly, plasma wave observations presented by Sulaiman et al. (2020)
93 show evidence of inertial Alfvén waves, intense ion cyclotron waves and whistler-mode auroral
94 hiss radiation. Field-aligned low-energy (100 eV/Q to 100 keV/Q) electron beams with
95 broadband energy distributions further support the existence of whistler-mode hiss and the
96 imprints of stochastic particle acceleration via Alfvén waves (Szalay et al., 2018). Finally, a
97 detailed look at the low-energy (10 eV/Q to 46 keV/Q) ion population suggests there is also a
98 significant amount of proton acceleration occurring both at the high-latitudes (in similar
99 locations to the electrons) and near the Io torus “boundary” – leading Szalay et al. (2020b) to
100 hypothesize that Alfvén waves generated near Io may be an important acceleration mechanism
101 for the protons as well.

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103 In this work, we are motivated by the aforementioned studies (e.g., Gershman et al., 2019;
104 Sulaiman et al., 2020, Szalay et al., 2020a, 2020b) to present the higher-energy charged particle
105 observations with a particular focus on the proton data obtained during Juno’s 12th perijove (PJ)
106 crossing of the IFP tail in the northern hemisphere. We focus on the proton measurements
107 because the Jupiter Energetic particle Detector Instrument (JEDI) (Mauk et al., 2017a) observed
108 the most significant ion acceleration event to date, strongly suggesting that the electromagnetic
109 coupling between Io and Jupiter is responsible for energizing protons up to ~ 1 MeV away from
110 the planet. We compare these data to the magnetic field (Gershman et al., 2019) and plasma
111 wave (Sulaiman et al., 2020) data from the same PJ12 IFP tail crossing near the MAW to better
112 understand the underlying physics governing this unusually intense and unexpected event.

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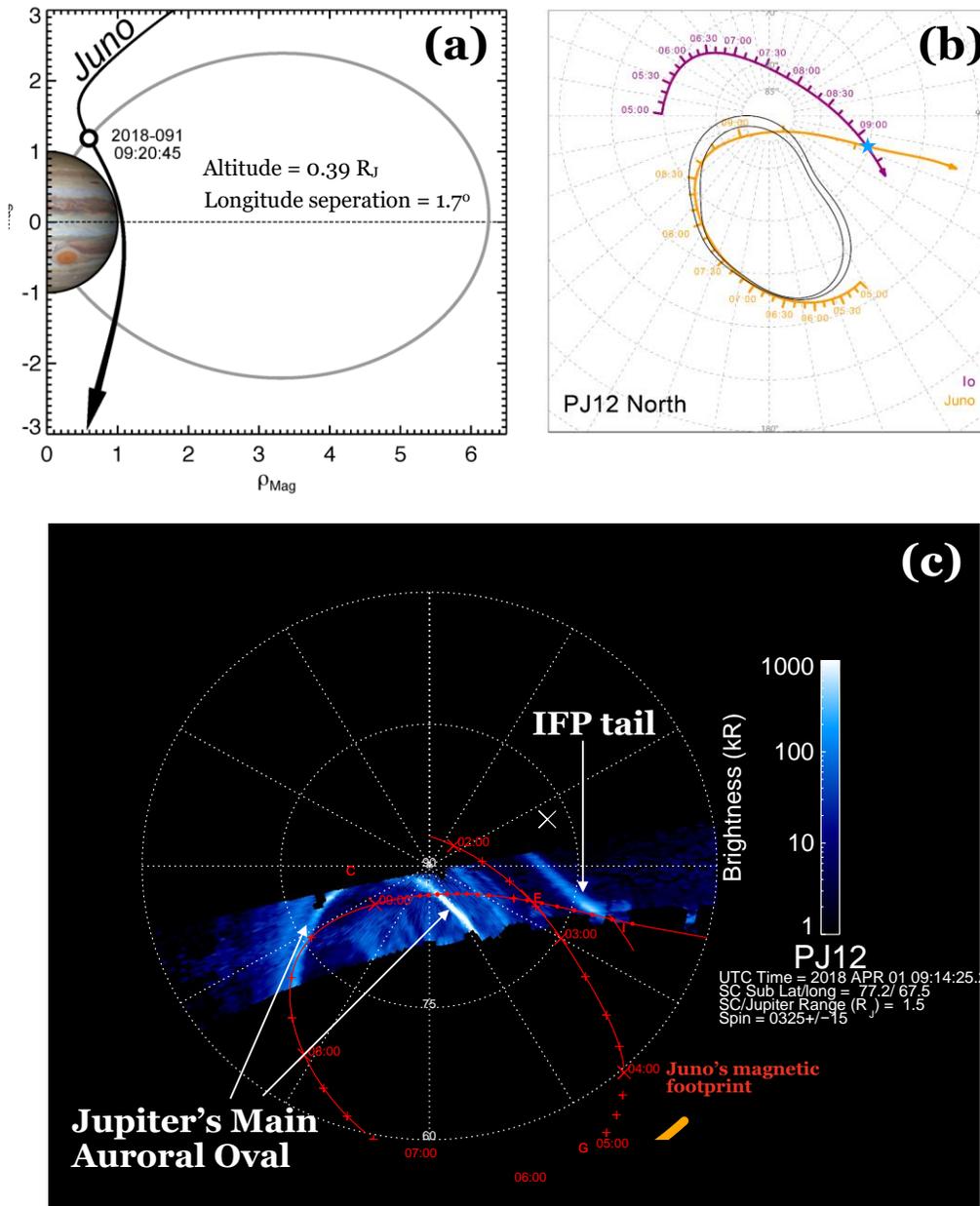
114 2. Observations

115 2.1 Juno’s crossing of the IFP tail

116 The data presented here were collected on the inbound leg of PJ12 as Juno crossed the IFP tail in
117 the northern hemisphere between $\sim 09:20:35$ to $09:20:55$ UT on 2018-091 (01 April, 2018).
118 Figure 1 is a trajectory schematic comprised of three different representations. Figure 1a
119 illustrates Juno’s intersection of a field line that maps to 5.9 Jovian radii (R_J , where $1 R_J =$

120 71,492 km), i.e., Io's orbital position, in a cylindrical magnetic dipole coordinate system. Figure
121 1b is a northern polar projection of Jupiter's auroral regions in system III coordinates that
122 encompasses the Io footprint tail (purple curve), Juno's magnetic footprint (orange curve) –
123 calculated using the JRM09 model (Connerney et al., 2018), and the statistical position of the
124 main auroral oval (black trace) derived from Hubble Space Telescope (HST) observations (e.g.,
125 Grodent et al., 2003). Juno crossed the IFP tail at an altitude of $0.39 R_J$ and with a longitudinal
126 separation of approximately 1.7 degrees from Io's Main Alfvén Wing (MAW) spot (e.g.,
127 Bonfond et al., 2009) when account for the Alfvén wave trajectory bendback between Io and
128 Jupiter's ionosphere. This remains Juno's closest approach to the MAW and potentially a direct
129 crossing (Szalay et al., 2020a). Figure 1c shows Juno ultraviolet spectrometer (UVS) (Gladstone
130 et al., 2017) observations of Jupiter's main auroral oval and the IFP tail approximately six
131 minutes prior to Juno crossing Io's footprint tail. The UVS data are presented in a system III
132 coordinate system with the red trace representing Juno's magnetic footprint.

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134

135 **Figure 1:** Io footprint tail crossing geometry. Panel a) Juno's trajectory in cylindrical magnetic coordinates with Io's
 136 M-Shell overlaid. Panel b) Magnetic footprints of Juno (orange curve), Io (purple) and the statistical location of
 137 Jupiter's main emission depicted by the bounding black curves (Bonfond et al., 2012). Panel c) Similar
 138 representation as panel b), but illustrates the ultraviolet brightness observations from Juno/UVS with Juno's
 139 magnetic footprint overlaid for reference (red curve). Juno/UVS observations occurred approximately six minutes
 140 before Juno crossed the IFP.

141

142 2.2 Brief description of Juno/JEDI

143 We focus on observations made by Juno's Jupiter Energetic particle Detector Instrument (JEDI).
 144 JEDI comprises three sensors (J90, J180 and J270) which measure the energy, angular and

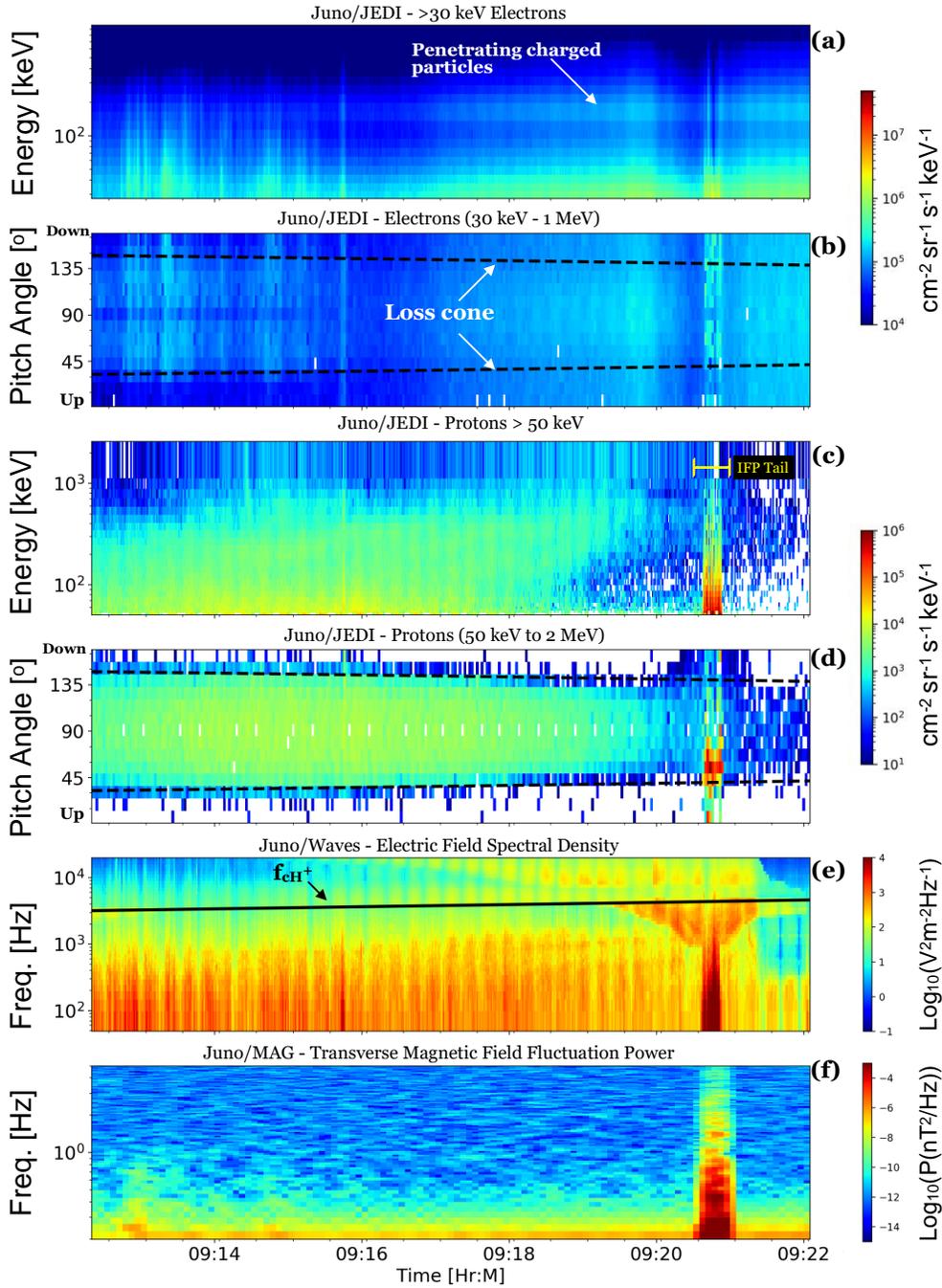
145 compositional distributions of >25 keV electrons and >10 keV ions (Mauk et al., 2017a). During
146 this event, the J90 and J270 sensors operated in a high rate mode, thus accumulating time-of-
147 flight by energy rates for 0.25 s at a cadence of 0.5 s, with no sector averaging. Pitch angle
148 distributions were obtained by combining the JEDI measurements with the measured local
149 magnetic field from Juno/MAG (Connerney et al., 2017b). The geometric loss cone size at this
150 time is 40° based on the dipole field approximation and 51° based on the JRM09 magnetic field
151 model (Connerney et al., 2018). Both methods agree well with the measured loss cone
152 distributions in the ion data. Each solid-state telescope has a full width at half maximum field-of-
153 view (FoV) that is approximately $\sim 17^\circ \times 9^\circ$ and therefore can resolve the loss cone in this
154 region. The duration of the footprint tail crossing is ~ 20 seconds, which is shorter than it takes
155 Juno to complete one revolution (Juno spins at ~ 2 revolutions per minute). This is important
156 because instantaneous look directions and pitch angle averaging between the two sensors can
157 average out fine structure in the IFP tail region. Therefore, we choose to perform all integral
158 moment calculations, i.e., characteristic energies and energy fluxes, using a 1 second sampling
159 window over a pitch angle range that contains just the upward moving protons (between 40° and
160 90°). The integral moment equations are outlined in Mauk et al. (2004) and Clark et al. (2018).

161

162 *2.3 Energetic charged particle observations*

163 Figure 2 presents an overview of the energetic charged particles (Fig. 2a – Fig. 2d) as well as the
164 plasma wave electric field spectral densities (Fig. 2e) from Sulaiman et al. (2020) and the
165 transverse magnetic field fluctuations (Fig. 2f) from Gershman et al. (2019). Plasma wave and
166 magnetic field measurements were obtained from Juno’s Waves (Kurth et al., 2017) and
167 Magnetic field (Connerney et al., 2017b) investigations, respectively. The most prominent
168 feature observed by JEDI is the dramatic proton intensity and pitch angle enhancements (Fig. 2c
169 and 2d) corresponding to the IFP tail. In Fig. 2d, protons in the IFP tail are shown to be
170 concentrated along the loss cone (horizontal dashed lines) in the upward direction. There is also
171 evidence of ions streaming upward along the local magnetic field line, but that feature only
172 persists for ~ 1 second. We do not discuss it further here. The energy-time distribution of the
173 protons (Fig. 2c) reveal broad energization ranging from ~ 50 keV to upward of ~ 1 MeV. During
174 the same time interval, the energetic electrons only show a modest response associated with the
175 IFP tail. Figures 2a and 2b show a slight enhancement in low energy (< 60 keV) electrons and a
176 slight decrease in the very energetic electron environment (> 1 MeV), which leaves a signature
177 indicated by the characteristic “penetrating charge particle” band near 160 keV (Mauk et al.,
178 2018). While ions show significant intensities in the upward loss cone, electrons mostly populate
179 the downward loss cone.

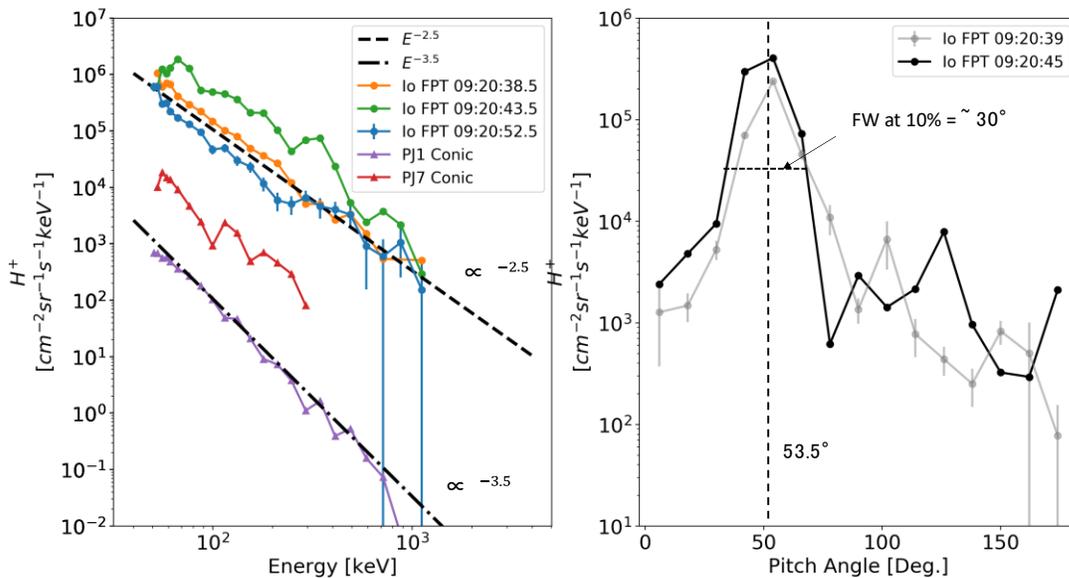
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182 **Figure 2:** Particles and Fields overview of the Juno PJ12 IFP tail crossing. Panels a-d) illustrate the Juno/JEDI
 183 observations of the energetic electrons and protons. Panels a & b) energetic electron energy-time and pitch angle-
 184 time spectrograms, respectively. Panels c & d) energetic proton energy-time and pitch angle-time spectrograms,
 185 respectively. Panel e) electric field frequency-time spectrogram and panel f) magnetic field frequency-time
 186 spectrogram. The black dashed lines in panels c & d represent the size of the loss cone in degrees using the dipole
 187 field approximation. The black solid curve in panel e represents the proton cyclotron frequency derived by
 188 Juno/MAG (Sulaiman et al., 2020).

189 Figure 3 shows proton energy spectra and pitch angle distributions for various times associated
 190 with Io's footprint tail crossing. The energy distribution of the protons resemble a power-law –
 191 monotonically decreasing intensities toward increasing particle energy (see Fig. 3 left panel).
 192 There is no clear evidence of peaked or accelerated Maxwellian-like energy distributions,
 193 representative of magnetic field-aligned electric fields (Clark et al., 2017b, Mauk et al., 2017b;
 194 Mauk et al., 2018). In Fig. 3 the energy spectra from published Juno/JEDI proton observations
 195 are compared (Clark et al., 2017a, Mauk et al., 2018). The observations made in the footprint tail
 196 suggest that the protons are more efficiently accelerated than in the other auroral regions, which
 197 can be seen by the power-law curves representing $E^{-2.5}$ and $E^{-3.5}$. Pitch angle distributions for
 198 two different times show clear peaks with centroids near 53° and a full width at the 10% level of
 199 $\sim 30^\circ$. Error bars in Figure 3 are determined by estimating the counting uncertainties associated
 200 with a Poisson distribution.



201

202 **Figure 3:** Measured energy spectra (left panel) and pitch angle distributions (right panel). For comparison purposes,
 203 energy spectra of proton conic distributions observed during PJ1 (Clark et al., 2017b) and PJ7 (Mauk et al., 2018)
 204 are also shown as well as power-law curves illustrating the different spectral slopes.

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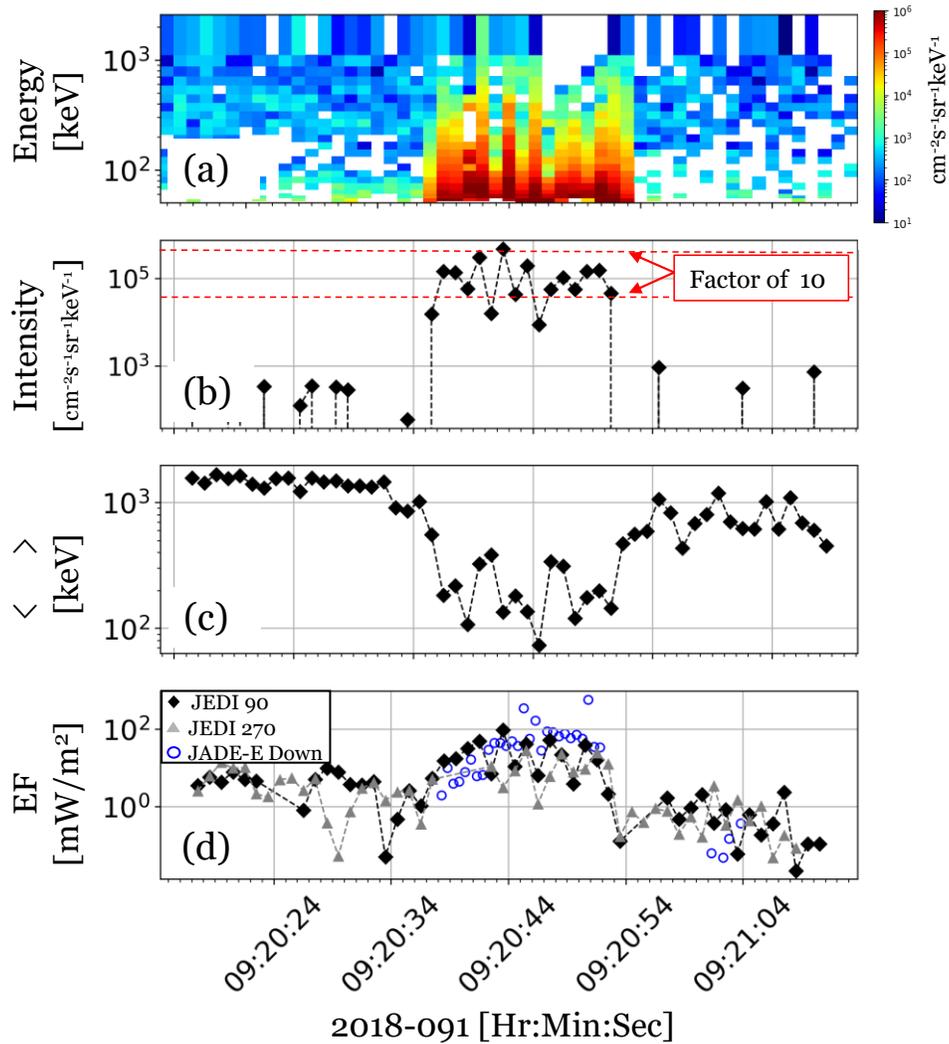
206 In Figure 4 we provide a closer inspection of the energy-time structuring and show the integral
 207 moments calculated using a one second sampling window over a pitch angle range that contains
 208 just the upward moving protons (between 40° and 90°). The energy-time spectrogram in Fig. 4a
 209 shows discrete stripes that occur somewhat regularly throughout the ~ 20 second IFP tail
 210 crossing. Similarly, in Fig. 4b the 100 keV proton intensities are chosen to highlight the
 211 variations, which fluctuate by factors of 3-10 on intervals as short as one second. Juno provides
 212 just a single point measurement and cannot disentangle the temporal/spatial ambiguity, therefore
 213 the 2-3 second variations may also be associated with ~ 50 - 100 km structures in the auroral
 214 region. It is possible that the variations are a measurement artifact due to the finite angular

215 resolution of JEDI. A crude analysis suggests that a collimated beam of particles can produce a
216 ~ 2 second variation in JEDI as a result of Juno's 12° s^{-1} rotation rate combined with the $\sim 27^\circ$
217 separation between the JEDI telescopes. This sort of temporal variation is observed in the polar
218 cap where electron beams are often narrower than JEDI can resolve (Mauk et al., 2017a;
219 Paranicas et al., 2018). In this particular event, the variation is likely not an artifact because the
220 measured width of the proton pitch angle distribution is relatively broad, i.e., $\sim 30^\circ$ (see Fig. 3),
221 compared to a single telescope FoV.

222

223 The integral moments associated with the IFP tail crossing show energetic protons characteristic
224 energies varying between $\sim 80 - 400$ keV (with a mean ~ 200 keV) (Fig. 4c) and likewise the
225 proton energy fluxes (Fig. 4d, averaged over pitch angles 40° - 90° from ~ 50 keV to 1 MeV) to
226 vary between $\sim 10 \text{ mW/m}^2$ to $\sim 100 \text{ mW/m}^2$ for the J90 sensor and $\sim 1 \text{ mW/m}^2$ to $\sim 30 \text{ mW/m}^2$ for
227 the J270 sensor. Note that instantaneous pitch angle coverage is attributed to these differences.
228 For comparison, we show the plasma electron (100 eV to 40 keV) precipitating energy fluxes
229 (Szalay et al., 2020a) measured by Juno/JADE-E (McComas et al., 2017). JADE-E energy fluxes
230 vary between $\sim 3 - 600 \text{ mW/m}^2$.

231



232

233 **Figure 4:** Panel a) JEDI J90 & 270 combined proton energy-time spectrogram filtered on pitch angles 40°
 234 $- 90^\circ$; panel b) 100 keV proton intensities; panel c) > 50 keV proton characteristic energies; panel d) J90
 235 (black diamonds), J270 (gray triangles) energetic proton energy flux vs. JADE-E (blue circles) energy
 236 fluxes of plasma electrons < 40 keV.

237

238 3. Discussion & Conclusions

239 The angular distribution of energetic protons along the upward loss cone reveal strong evidence
 240 for energetic ion conic acceleration associated with Io's footprint tail and probably the MAW.
 241 Ion conics are the result of thermal ionospheric ions heated perpendicular to the magnetic field
 242 via wave-particle interactions and then accelerated upward due to gradients in the magnetic field
 243 and/or field-aligned electric fields (e.g., Klumpar et al., 1979; Gorney et al., 1985; Chang, 1993;

244 Retterer et al., 1994; Carlson et al., 1998; Lynch et al., 2002). Wave-heating alone does not
 245 produce the most energetic ions; therefore, it is thought that electrostatic confinement via
 246 magnetic field-aligned potentials is required to trap the ions and further accelerate in the wave-
 247 heating region. This is referred to as the so-called “pressure cooker” mechanism (e.g., Gorney et
 248 al., 1985). Observations from Parker Solar Probe close to the Sun (Mitchell et al., 2020), Cassini
 249 at Saturn (Mitchell et al., 2009) and Juno at Jupiter (Clark et al., 2017b) have confirmed their
 250 existence elsewhere in the solar system, but have not been directly observed as a result of planet-
 251 moon interactions. Below we discuss possible proton acceleration mechanisms associated with
 252 the IFP tail.

253

254 The first mechanism we consider is a cyclotron resonant heating mechanism. Sulaiman et al.
 255 (2020) analyzed the Juno/Waves measurements during the PJ12 IFP tail crossing and found
 256 evidence of upward-propagating, left-hand polarized ion cyclotron waves with large spectral
 257 densities (maximum of $\sim 10^5 \text{ V}^2\text{m}^{-2}\text{Hz}^{-1}$) near and at the proton cyclotron frequency. Using the
 258 theoretical energy transfer relationship from Chang et al. (1986), Sulaiman et al. (2020)
 259 estimated the ion heating rate, denoted as dW_{\perp}/dt , to have an upper limit of $\sim 500 \text{ eV/s}$. To
 260 estimate the proton energies achievable from this heating rate we need to know the time-of-flight
 261 of the ions between their source region and the spacecraft. First, the altitude of the source region
 262 can be estimated from the measured pitch angle distributions, shown in Fig. 3, and by assuming
 263 the first adiabatic invariant is conserved as the protons are transported along the magnetic field.
 264 We also assume the protons are heated purely perpendicular to the local magnetic field, i.e., pitch
 265 angles of 90° , in the source region (see similar method outlined in Clark et al., 2017b and
 266 references therein) and neglect changes in an ion’s pitch angle as it is transported along the field
 267 line. The measured local magnetic field during the IFP tail crossing is $\sim 3 \times 10^5 \text{ nT}$ and the
 268 centroid of the proton pitch angle distributions vary between $\sim 50^\circ$ and 60° . Combining this
 269 information together and using the latest magnetic field model (JRM09; Connerney et al., 2018),
 270 we find the source location to be $\sim 11,000 \text{ km}$ or $0.16 R_J$ above Jupiter’s 1-bar oblate surface. The
 271 last piece of information required is the bulk velocity of ions in Jupiter’s ionosphere. The only
 272 published ion bulk flow measurements in this region to date is from a study of low-energy ions
 273 in Jupiter’s topside ionosphere using the JADE-Ion sensor (Valek et al., 2020). The authors
 274 performed a numerical integration of the plasma proton distributions and derived an outflow
 275 speed, v_{bulk} , of 20 km/s . Here we assume this to be the outflow speed of the protons in the
 276 region connected to the IFP tail and thus, the time-of-flight of the protons is estimated to be
 277 approximately $t \approx 900 \text{ s}$ where, $t = d/v_{bulk}$, where $d = 18,500 \text{ km}$ is the integrated length
 278 along the field between Juno at $0.33 R_J$ and the source region at $0.16 R_J$. Multiplying the
 279 proton’s time-of-flight with the heating rate derived by Sulaiman et al. (2020) suggests ion
 280 cyclotron heating may be able to produce conic energies as large as 450 keV . This number is
 281 commensurate with the characteristic energies of the proton observations in the IFP tail (see Fig.
 282 4c). Major limitations of this crude estimate include the assumption that the wave-heating is
 283 constant along the flux tube between the source region and the spacecraft and the bulk ion speed
 284 remains the same. While the wave-heating assumption appears to be reasonable in Earth’s

285 auroral region (e.g., Lynch et al., 2002), it is uncertain if the same holds true for Jupiter. Next we
286 consider the role of Alfvén waves as a possible energization mechanism.

287

288 Gershman et al. (2019) analyzed the magnetic field fluctuations in Jupiter’s polar magnetosphere
289 and found direct evidence of strong transverse perturbations associated with the IFP tail. The
290 perturbations were identified as Alfvénic (between 0.2 to 5 Hz) and the Alfvén Poynting flux
291 was calculated to be as high as $\sim 3,000$ mW/m² during the likely PJ12 crossing of Io’s MAW.
292 Sulaiman et al. (2020) used Juno/Waves data to demonstrate that Alfvénic fluctuations, first
293 observed by MAG, extend into the higher frequencies spanning a range from ~ 50 to 800 Hz.
294 Clearly, Alfvén waves are present and carry a significant source of energy in the IFP tail and
295 near the MAW. In this letter, we do not address the fundamental idea regarding the role of
296 Alfvén waves in generating ion conics, but we discuss important observational details and ideas
297 that may help future theoretical pursuits. One striking observation is the energy partitioning
298 between the Alfvénic Poynting flux and the particle energy fluxes. In Figure 4d, we show that
299 the energetic ion energy flux is comparable to the 0.1 to 40 keV electron energy flux. If the
300 energy reservoir is the same for the two populations then the observations presented here
301 suggests that energy conversion efficiencies between Alfvén waves and ions (~ 3 -5%) are
302 comparable to the lower-energy electrons except for brief moments where the electron energy
303 flux peaks as high as 580 mW/m². Numerous studies have investigated ion acceleration in
304 Earth’s aurorae and the role of Alfvén waves (e.g., Li and Temerin, 1993; Knudsen and
305 Wahlund, 1998; Chaston et al., 2004). However, in the absence of wave-particle interaction
306 models for the Io fluxtube and its tail, which consider the ion response specifically, we turn to
307 comparisons with models at larger L-shells. Saur et al. (2018) find that on L-shells between 10
308 and 40 R_J, at high latitudes, ion-Landau damping is effectively not taking place, while electron
309 Landau damping of inertial Alfvén waves is a highly effective acceleration mechanism in
310 accordance with previous modeling and existing observations of energetic electrons (e.g. Hess et
311 al. 2010, Hess et al. 2013, Bonfond et al. 2017, Saur et al. 2018, Clark et al., 2018, Szalay et al.
312 2020a). If the temporal scales of the waves become extremely small, then ion-cyclotron damping
313 becomes more prominent (e.g., Sulaiman et al. 2020). Of the two resonant mechanisms
314 discussed, i.e., Landau and cyclotron damping, non-resonant mechanisms (e.g., Lu and Li, 2007)
315 of ion acceleration through Alfvén waves have not been studied for the Jupiter system and their
316 effectiveness is thus difficult to assess without detailed studies.

317

318 The Juno/JEDI data presented in this study represent the first measurements of energetic proton
319 conics associated with Io’s footprint tail near the MAW. This discovery showcases the diversity
320 of planetary systems and interactions present where ion conics exist, e.g., Earth, Saturn and
321 Jupiter’s auroral regions and now as a result of moon-magnetospheric interactions. Our primary
322 conclusions in this study are the following:

- 323 1. Energetic proton acceleration associated with the IFP tail appears significant and perhaps the
324 most intense ion event recorded by Juno/JEDI to date.

- 325 2. The angular distributions of the protons suggest these are the ion conic distributions and are
 326 likely accelerated by ion cyclotron waves via a resonant interaction; however, Alfvénic
 327 turbulence was not ruled out and may play a role.
- 328 3. Proton acceleration associated with Io’s footprint tail is more intense than compared to the
 329 main auroral (Mauk et al., 2018) or polar cap regions (Clark et al., 2017b), thus highlighting
 330 the unique and strong electromagnetic interaction between Jupiter and Io.

331

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340

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