

# **Estimating the occurrence of geomagnetic activity using the Hilbert-Huang transform and extreme value theory**

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## **Key Points**

- A Hilbert-Huang transform is applied to the geomagnetic aa data to identify the solar cycle dependency in the data
- Extreme value theory is applied to the aa data separately in solar minimum and maximum conditions
- March 1989 (Quebec) event is shown to overall be a 1-in-25 year event. But a 1-in-130 year event during solar minimum

## Abstract

In this paper extreme value theory (EVT) has been used to estimate the return levels for geomagnetic activity based on the aa index. The aa index is the longest, continuously recorded, geomagnetic dataset (from 1868 – Present). This long, 150 year, dataset is an ideal candidate for extreme value analysis. However the data are not independent and identically distributed as required for EVT since they are impacted by the approximately 11 year solar cycle. The Hilbert-Huang Transform has been used to identify the solar cycle component in the data and the data has been split into solar maximum and minimum times. In these two regimes the generalised extreme value distribution has been fit to the datasets. These have also been combined for an estimate of the overall return times. The results suggest that the largest event in the database (March 1989) is a one in 25 year event. However, considering separate solar maximum and minimum times has a large impact on the return times. During solar minimum conditions the return time of the March 1989 event is 130 years. This suggests that the occurrence of extreme space weather events is conditionally dependent on where in the Solar Cycle we are.

## Introduction

Geomagnetic storms are disturbances in the Earth's magnetosphere. They are caused by changes in the solar wind which impact the magnetosphere. Two of the main causes for these changes in the solar wind are coronal mass ejections (CMEs) and high-speed solar wind streams (HSSs) [Schwenn, 2007]. CMEs usually have large speeds (approximately five times faster than the background solar wind), high energies and large magnetic field strengths [Riley and Love, 2017]. HSSs come from

solar coronal holes and the fast wind from these regions interacts with the slower upstream wind which create co-rotating interaction regions (CIRs) [Garton *et al.*, 2018]. These regions have increased magnetic field strength and higher particle density [Schwenn, 2007]. A lot of the largest space weather impacts are associated with geomagnetic storms: geomagnetically induced currents, radio scintillation, solar energetic particle events and enhanced fluxes of relativistic electrons [Cannon, 2013].

Indices are used to quantify the relative strength of geomagnetic events, these include Dst ([World Data Center for Geomagnetism Kyoto *et al.*, 2015b]), Kp, Ap ([Bartels, 1957]), AE, AO, AL, AU ([World Data Center for Geomagnetism Kyoto *et al.*, 2015a]), am, as, an ([Mayaud, 1980]), aa and Aa ([Mayaud, 1972]). The indices are calculated at different cadences from hourly to daily. Of these, the Dst index is used for identifying and quantifying the severity of a geomagnetic storms [e.g. Loewe and Pross, 1997]. However the various indices are, for the most part, closely related to each other.

Due to the impact of space weather events on human health and technology there is interest in estimating the return time for the most extreme events. A common reference point for extreme space weather events is the so called 'Carrington event' [Carrington, 1859], which was one of the largest space weather events in the last 200 years [Cliver and Svalgaard, 2004]. A key question is what is the return time (likelihood) of extreme space weather events. This is a difficult question to answer satisfactorily as it requires investigating the tails of probability distributions, where

there is little data. However, this can be done rigorously using extreme value theory/statistics (EVT).

EVT is mathematical rigorous and provides sensible measures of uncertainty, which can be very large when there are few data points. Therefore one of the key difficulties associated with using EVT is the need for sets of large independent large samples. For example when looking at extreme temperatures (hot or cold) in meteorology the annual maxima or minima are suitable time scales since, on the whole, yearly temperatures are independent, whilst daily temperatures are not. In the space weather domain it would be ideal to take solar cycle (11 year) minima or maxima time series. Unfortunately there is not enough recorded data to have enough data points remaining for effective analysis.

A number of authors have applied EVT to the space weather domain, including: [Elvidge and Angling, 2018; Koons, 2001; Meredith et al., 2015; Silbergleit, 1996; 1999; Siscoe, 1976; Thomson et al., 2011; Tsubouchi and Omura, 2007]. Extreme space weather events have also been estimated using techniques other than EVT such as fitting power law distributions, log-normal distributions and generalized Pareto distribution [Chapman et al., 2020; Chapman et al., 2018; Riley, 2012; Riley and Love, 2017]. The overall goal of each of the papers is to try to quantify the statistics of a particular measurable associated with an extreme space weather event.

In terms of investigating geomagnetic activity: *Silbergleit* [1996] and *Tsubouchi and Omura* [2007] used EVT to investigate extreme events in the DST index using 23 and 44 years of data respectively, *Koons* [2001] used the Ap index using 66 years of data, *Siscoe* [1976] and *Silbergleit* [1999] variants of the aa index, using 91 and 124 years of data, and *Thomson et al.* [2011] the rate of change of the magnetic field using 31 years. *Riley and Love* [2017] also estimated the probability of extreme DST events using 60 years of data. Of those, the 91 and 124 year datasets of *Siscoe* [1976] and *Silbergleit* [1999] are very useful for EVT since they capture the longest time period (8 – 11 solar cycles). However, as well as the length of the datasets, the number of data used is also crucial in reducing the uncertainty in the analysis. *Siscoe* [1976] performed EVT only using the three largest events in each solar cycle during the test period, resulting in 27 data points and *Silbergleit* [1999] used the maximum value from each solar cycle, resulting in 12 data points. Whilst these approaches break the data into suitable scale sizes, few data points remain which means there are substantial uncertainties in the results.

In this paper the aa index (1868 – 2018; 150 years) is analysed using the annual maximum values. This results in the largest temporal span of data for EVT in the space weather domain. Recent work by *Chapman et al.* [2020] has used a linear “mapping” between the top few percent of the aa index and the annual minimum Dst index value to estimate the probability of extreme Dst values. This work takes advantage of the extra data points from this aa-to-Dst mapped set.

Whilst geomagnetic storms tend to last between two and seven days, the annual maxima has been chosen to increase the likelihood of having independent,

identically distributed data, a requirement for extreme value theory. Using, for example, weekly rather than annual maxima can introduce further dependencies in the dataset as individual events may, or may not, originate from the same solar active region. This could result in the analysed data not being identically distributed. However the disadvantage of using annual maxima is that the solar variability over the ~11 year cycle remains embedded in the data. This temporal dependence is accounted for by using the Hilbert-Huang transform [*Huang and Wu, 2008*] to split the data into solar maximum and minimum times.

## **Extreme Value Theory**

EVT provides a sophisticated approach for estimating probability distribution functions, and specifically for looking at the tail of such distributions. The method avoids any starting assumption about the underlying distribution [*Coles, 2001*]. The key result from EVT is the Fisher-Tippett-Gnedenko (FTG) theorem which states that the maximum of an independent and identically distributed (iid) random variable converges to one of only three possible distributions: the Gumbel distribution [*Gumbel, 1935*], the Fréchet distribution [*Fréchet, 1927*], or the Weibull distribution [*Weibull, 1951*], which can be grouped into the generalized extreme value distribution.

Specifically, for a sequence of iid random variables  $X_1, X_2, \dots, X_n$  with common distribution function  $F$  let  $M_n = \max \{X_1, \dots, X_n\}$  and  $w = \sup \{x: F(x) < 1\}$  then

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141  $\Pr(M_n \leq x) = \Pr(X_1 \leq x, \dots, X_n \leq x) = F^n(x)$  (Equation 1)

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143 Then as  $n \rightarrow \infty$ ,  $F^n(x) \rightarrow 0$  if  $x < w$  and  $F^n(x) \rightarrow 1$  otherwise, as such  $M_n \rightarrow w$ . To

144 avoid a degenerate distribution  $F^n(x)$  is normalised. Assuming there is a non-

145 degenerate distribution  $G$  such that, for normalising constants  $a_n > 0$  and  $b_n$ :

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147  $\lim_{n \rightarrow \infty} F^n(a_n x + b_n) = G(x)$  (Equation 2)

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149 where  $G$  is the generalized extreme value (GEV) distribution defined by

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151  $G(x) = \exp \left\{ -1 \left[ 1 + \xi \left( \frac{x-\mu}{\sigma} \right) \right]^{-\frac{1}{\xi}} \right\}$  (Equation 3)

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153 defined for  $1 + \frac{\xi(x-\mu)}{\sigma} > 0$  and where  $\mu$  is the location parameter,  $\sigma > 0$  the scale

154 parameter and  $\xi$  the shape parameter [Coles, 2001]. For  $\xi < 0$  the GEV reduces to

155 the Weibull distribution, for  $\xi > 0$  the Fréchet distribution and in the limit  $\xi \rightarrow 0$ ,  $G(x)$

156 reduces to

157

158  $G(x) = \exp \left\{ -\exp \left( -\frac{x-\mu}{\sigma} \right) \right\},$  (Equation 4)

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the Gumbel distribution.

The parameters of the GEV are usually estimated using a maximum log likelihood method [Coles, 2001]. However, the requirement that the variables must be iid is usually a barrier with using raw data directly and some form of pre-processing is normally required.

## Data

The aa index is a global geomagnetic index which is based on the largest horizontal deviation of the magnetic field measured in nT. It is based on data from two nearly antipodal stations, one in the UK and another in Australia and has been continuously recorded since 1868. Over the 150 years the stations where the data have been recorded has changed. In order to maintain a constant value for the index the weighting of the different stations have varied over time (Table 1).

*Table 1: Weighting factors for the stations used to compile the aa index [International Service of Geomagnetic Indices, 2013].*

Northern hemisphere station (UK)			Southern hemisphere station (Australia)		
Time range	Station	Weighting factor	Time range	Station	Weighting factor
1868 - 1925	Greenwich	1.007	1868 - 1919	Melbourne	0.967
1926 - 1956	Abinger	0.934	1920 - 1979	Toolangi	1.033
1957 - Present	Hartland	1.059	1980 - Present	Canberra	1.084



180 The main advantage of using the aa index for EVT is that it is the longest running  
181 planetary index of geomagnetic activity. This long sample time helps in reducing the  
182 uncertainties in the EVT extrapolation, however on this time scale the impact of the  
183 solar cycles becomes apparent, which has been shown to have an impact on the  
184 results of extreme value modelling [*Riley and Love, 2017*]. Figure 1 shows the time  
185 series of the aa index in the top panel (each point is the annual maximum aa value)  
186 and the bottom panel of the figure shows the periodogram created with a Hamming  
187 windowing function on the data. The large peak in the periodogram corresponds to  
188 10.7 years and is the solar cycle contribution to the data. To perform EVT on the  
189 dataset this temporal dependence should be accounted for. In this work the Hilbert-  
190 Huang transform (HHT) is used to identify solar maximum and minimum times, which  
191 are assumed to each be iid, and EVT can be performed on each.

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193 The HHT decomposes a time series into intrinsic mode functions (IMFs) and then  
194 finds the instantaneous frequency of each IMF [*Norden E Huang and Wu, 2008*]. The  
195 first step of the HHT is to use empirical mode decomposition (EMD) [*N. E. Huang et*  
196 *al., 1998*]. Similar to the Fourier and Wavelet transform, EMD splits a signal into its  
197 components, called intrinsic mode functions (IMFs). An IMF is a function in which the  
198 number of extrema and zero-crossings differ by at most one and at each point the  
199 mean value of the envelopes, defined by the local maxima and minima, is zero. The  
200 sum of the IMFs reconstitute the original signal. Hilbert spectral analysis (HSA) can  
201 then be used by applying the Hilbert transform to each IMF to find the instantaneous  
202 frequency [*N. E. Huang et al., 1998*]. Unlike Fourier and Wavelet transforms HHT is

an algorithmic approach rather than theoretical. However the main advantage of the HHT over Fourier and Wavelet is that it is suitable for nonlinear and non-stationary data.

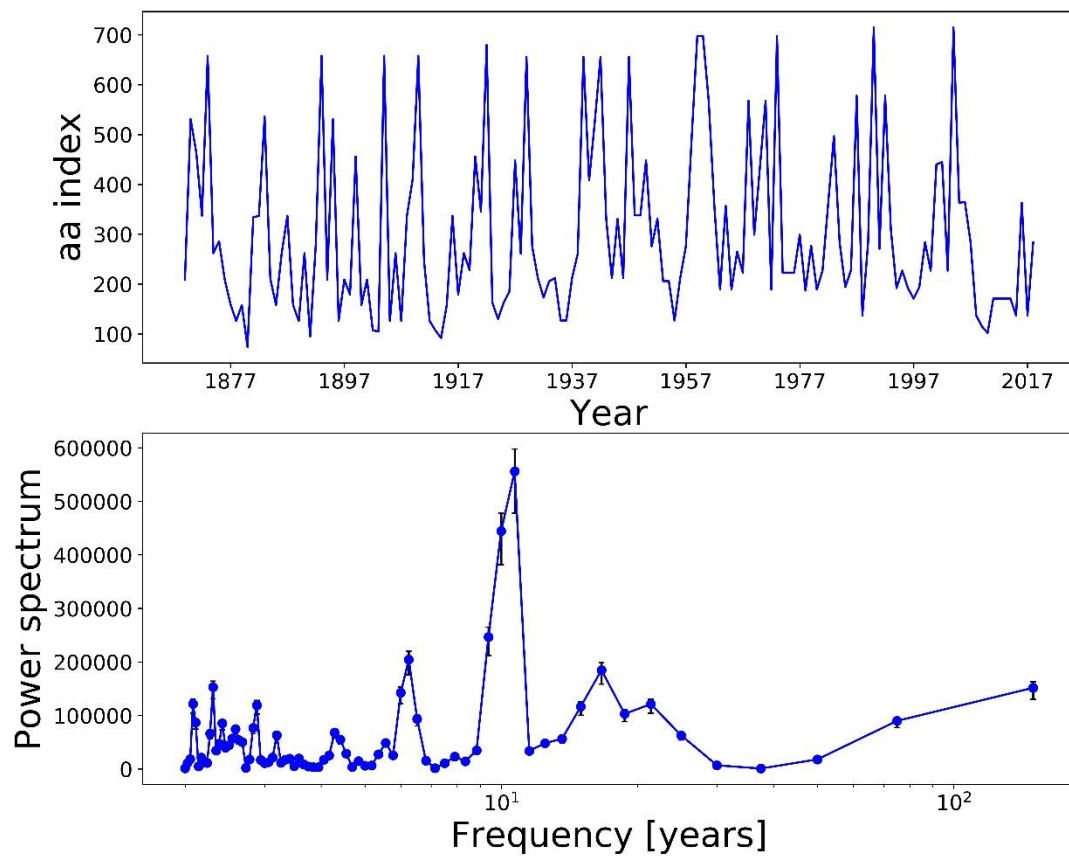


Figure 1. Top panel shows the time series of the annual maxima aa index. Bottom panel shows the periodogram of the series.

213 Applying the HHT to the aa index data (top panel of Figure 1) results in nine IMFs  
214 shown in Figure 2. The top panel of the figure shows the original aa index values (in  
215 red), then each of the IMFs are shown in blue. Using HSA the envelope of each IMF  
216 has been found (shown in green) and the value of the instantaneous time period  
217 (ITP) (one over the instantaneous frequency) is shown in the upper left of each IMF  
218 plot. A single value for the ITP is found by fitting a linear polynomial through the data.  
219 In each case the polynomial was of order zero (as expected) and the constant term  
220 is shown. From the figure it can be seen that the third IMF has an ITP which  
221 corresponds to the peak time period from Figure 1. This provides confidence that this  
222 particular IMF corresponds to the sunspot cycle component in the aa index data.

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224 This particular IMF could be used in the fitting of the GEV as part of a temporally  
225 varying location and scale parameters [Coles, 2001]. However this requires  
226 propagation of the IMF forward in time. With only simplistic ways of interpolating the  
227 IMF forward in time (it is hard to predict the next solar cycle) it makes finding the  
228 return times difficult (a key attribute for extreme value modelling).

229

230 Instead the aa variables can be made to be iid by separating the solar maximum and  
231 minimum conditions. In each of the the two cases it can be assumed that they are  
232 distributed according to the same function. Rather than using a separate dataset to  
233 try and determine when these times are, the third (solar cycle dependent) IMF can  
234 be used. The IMF is centred at zero and positive values can be used to describe  
235 solar maximum times, whilst negative values can be used as solar minimum times.  
236 Comparing the IMF estimated solar maximum/minimum times to the sunspot cycle

(as independent verification of this approach) shows that the estimated maximum overlaps with the peak of the sunspot number as would be expected (Figure 3).

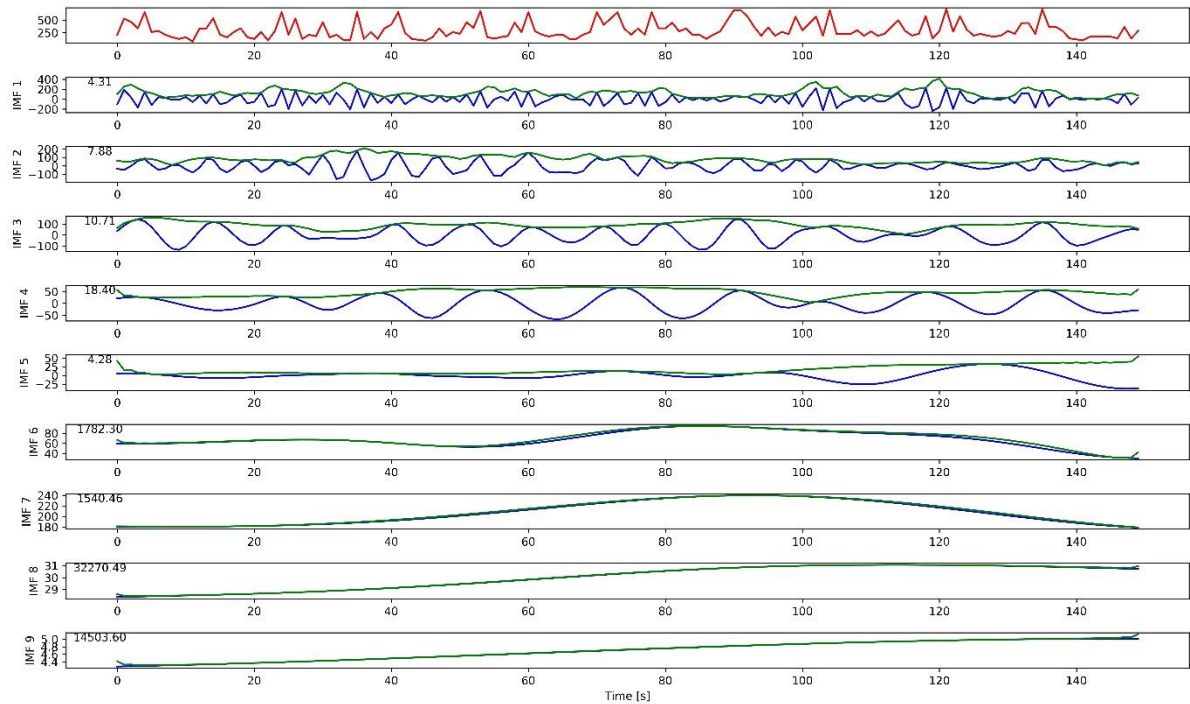


Figure 2. aa index decomposed into its 9 IMFs. The original aa index is shown in the top panel (red) and each panel below shows an IMF (blue). For each IMF the envelope has been found using the Hilbert transform (green) and the dominant instantaneous frequency value is shown in the upper left of each IMF plot.

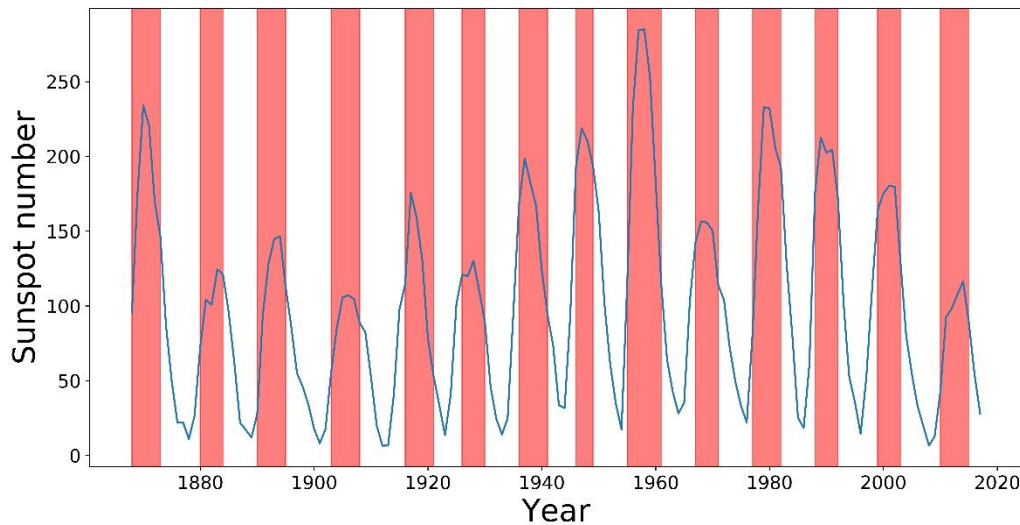


Figure 3. Annual maximum sunspot number [SILSO World Data Center, 2018]. Solar maximum times (from the third IMF from the aa index) are shaded.

## Results

Fitting the GEV to the solar maximum and minimum time series using least log likelihood results in the estimated  $\mu, \sigma, \xi$  (and standard errors) as shown in Table 2. One method for verifying the quality of the fit of the GEV distribution is by looking at the quantile plot which shows the pairs

$$\left\{ G^{-1}\left(\frac{i}{n+1}\right), x_i \right\}, \quad i \in \{1, \dots, n\} \quad (\text{Equation 5})$$

for the ordered annual aa values  $\{x_1, x_2, \dots, x_n\}$  and where  $G^{-1}(x)$  is the inverse of Equation 3, given by [Coles, 2001]:

$$G^{-1}\left(\frac{i}{n+1}\right) = \mu - \frac{\sigma}{\xi} \left( 1 - \left( -\log\left(\frac{i}{n+1}\right) \right)^{-\xi} \right). \quad (\text{Equation 6})$$

The quantile plot for solar maximum and minimum times is shown in Figure 4. The plot being roughly linear is an indication of a good agreement between the model fit and the empirical data.

Table 2. The GEV fit parameters for the solar maximum and minimum conditions, the standard error is shown in parenthesis.

	Number data points	$\mu$	$\sigma$	$\xi$
<b>Maximum</b>	78	279 (17.8)	130 (12.6)	-0.03 (0.12)
<b>Minimum</b>	72	174 (9.24)	68.8 (7.16)	0.15 (0.10)

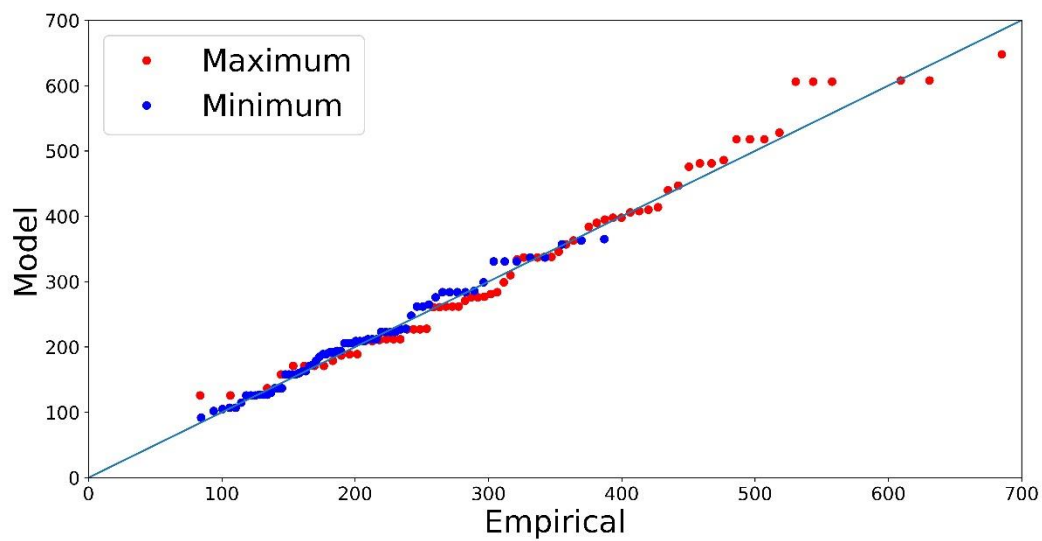


Figure 4. Quantile plot of the GEV fit for solar maximum and minimum times.

Using the fitted GEV distributions the return time for any given event can be estimated for either solar maximum or minimum conditions using the values in Table 2. However another useful return time would combine both solar maximum and minimum. The combined return time can be found by solving

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280  $G_{max}(z_p)G_{min}(z_p) = 1 - \frac{1}{p}$  (Equation 7)

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282 for  $z_p$ , the return level with return period  $1/p$  ( $z_p$  is expected to be exceeded by the  
283 annual maximum aa value with probability  $p$ ), and where  $G_{max}$ ,  $G_{min}$  are the GEV  
284 distributions defined by the parameters in Table 2 [Coles, 2001]. The return levels for  
285 10, 50, 100, 500 and 1000 years are shown in Table 3, with the standard errors  
286 shown in parenthesis. It is interesting to note the difference in return times between  
287 solar maximum and minimum times. This suggests that the probability of an extreme  
288 event is conditionally based on what part of the solar cycle we are in. This is in  
289 agreement with the findings of *Riley and Love* [2017] who determined that, assuming  
290 a power law distribution, the probability of geomagnetic storm exceeding the  
291 Carrington event (in terms of Dst) was 1.4% during solar minimum conditions and  
292 28% for solar maximum conditions.

293

294 Table 3. Return levels for 10, 50, 100, 500 and 1000 year return periods for solar  
295 maximum and minimum conditions as well combined results. The standard error is  
296 shown in parenthesis.

	10-year	50-year	100-year	500-year	1000-year
<b>Maximum</b>	611 (50)	873 (146)	989 (208)	1271 (408)	1399 (519)
<b>Minimum</b>	365 (30)	565 (89)	670 (131)	969 (287)	1127 (386)
<b>Combined</b>	631 (60)	898 (155)	1019 (216)	1321 (429)	1463 (562)

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300 These results have been compared to previous EVT work on the Ap index (an index  
301 similar to the aa index) undertaken by *Koons* [2001]. To compare the results, since  
302 the data are on different scales, they have been normalised by dividing through by  
303 the peak value of the March 1989 event [*Feynman and Hundhausen*, 1994]. This is  
304 the largest value in the aa index database (tied with the “Halloween” event of 2003)  
305 and the second largest in Ap (the largest is an event is from November 1960). *Koons*  
306 [2001] provides the fit parameters for the Gumbel distribution which was shown to  
307 have the best fit with the data. However no standard error was reported in the  
308 results. So for ease of comparison in this work the same 66 years of data was  
309 analysed and fit with the same Gumbel distribution as in *Koons* [2001] ( $\mu = 99.1409$   
310 and  $\sigma = 42.9416$ ). The normalised results for both the aa and Ap index data are  
311 shown in Table 4.

312

313 Table 4. Return levels for this work as well as previous EVT work on Ap [*Koons*,  
314 2001]. The values are normalised by dividing through by the index value of the event  
315 in March 1989.

	Normalised 10-year	Normalised 50-year	Normalised 100-year	Normalised 500-year	Normalised 1000-year
Aa (Combined)	0.88 (0.08)	1.26 (0.21)	1.43 (0.30)	1.85 (0.60)	2.05 (0.79)
Ap [ <i>Koons</i> , 2001]	0.80 (0.08)	1.09 (0.17)	1.21 (0.22)	1.43 (0.34)	1.54 (0.42)

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Comparing the EVT results between the aa and Ap index show that they are very similar for 'short' return times (10 years) with values of 0.88 and 0.80 respectively. These differences widen as the return periods get longer. This is to be expected as in this work 150 years of aa data have been used for the estimates compared to 66 years of Ap data from *Koons* [2001]. The extra data points should provide better estimates for the longer return periods, especially the 100-year return level.

## Conclusions

Extreme value theory (EVT) has been used to estimate the return levels for geomagnetic activity based on the aa index. The aa index is the longest, continuously recorded, geomagnetic dataset. This long, 150 year, dataset is an ideal candidate for extreme value analysis. Whilst the aa index is not the most commonly used space weather index its close relationship with the more commonly used Dst index [*Chapman et al., 2020*] implies that similar geoeffective impacts of extreme Dst events would be felt during extreme aa events. However, the aa data are not independent and identically distributed (iid) as required for EVT as they are impacted by the approximately 11 year solar cycle. The Hilbert-Huang Transform has been used to identify the solar cycle component in the data and the data have been split into solar maximum and minimum times. In these two regimes the variables are assumed to be iid, and the generalised extreme value (GEV) distribution has been fit to the two datasets. These have also been combined for an estimate of the return times.

The results suggest that the largest event in the database (March 1989 / October 2003) is a one in 25 year event (but with a standard error of 86 years). Whilst this may seem counter-intuitive, since there are only two events of that size (aa of 715) in the 150 year database, there are in total eight events where the aa index exceeds 650. Considering separate solar maximum and minimum times has a large impact on the return time. During solar minimum the return time of the March 1989 event is 130 years (with a standard error of 145 years). This value seems reasonable since there is one event with an aa > 650 during solar minimum in the 150 year database (August 4th 1972; [Knipp *et al.*, 2018]). However it is in contrast to the results of Riley and Love [2017] who report the probability as ~1-in-700 years during solar minimum (assuming a power law distribution). This demonstrates the uncertainty that can arise when extrapolating extreme events by using different underlying assumptions. Quantifying the impact of solar minimum is of particular importance in quantifying the likelihood associated with extreme space weather events if a period of extended solar minimum is entered. It has been estimated that there is a 15 - 20% chance of returning to Maunder Minimum-like conditions within the next 40 years [Ineson *et al.*, 2015; Lockwood, 2010].

## Acknowledgements

The aa data can be downloaded from British Geological Survey ([http://www.geomag.bgs.ac.uk/data\\_service/data/magnetic\\_indices/aaindex.html](http://www.geomag.bgs.ac.uk/data_service/data/magnetic_indices/aaindex.html)).

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