

ARTICLE TYPE

Special decompositions and linear superpositions of nonlinear systems: BKP and dispersionless BKP equations

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The existence of decomposition solutions of the well-known nonlinear BKP hierarchy is analyzed. It is shown that these decompositions provide simple and interesting relationships between classical integrable systems and the BKP hierarchy. Further, some special decomposition solutions display a rare property: they can be linearly superposed. With the emphasis on the case of the fifth BKP equation, the structure characteristic having linear superposition solutions is analyzed. Finally, we obtain similar superposed solutions in the dispersionless BKP hierarchy.

KEYWORDS:

Linear superpositions, KdV and dispersionless KdV decompositions, BKP and dBKP hierarchies, integrable systems

1 | INTRODUCTION

The linear superposition principle is encountered in many branches of physics and can be employed to solve essentially all linear problems¹. This principle states that, for linear systems, superposition provides large classes of solutions and gives rise to generalized solutions of linear problems in the form of a linear combination of the independent solutions. When linear superposition holds, a system can be decomposed into its constituent parts and the behavior of each component is independent of other components². Conversely, when linear superposition fails, such systems often exhibit behaviors reflecting the fact that individual system components are not independent of each other. Linear superposition normally operates in linear systems. Instead, the lack of linear superposition is one of the crucial characteristics of nonlinear systems. The loss of linear superposition also has implications for causation, reduction, emergence, and natural laws in nonlinear dynamics, all of which raise important issues for the application of nonlinear models to real-world problems. Generally, because of the nonlinear terms, linear superposition principle does not hold well in nonlinear theories even for nonlinear integrable systems where the nonlinear superpositions are valid. Nonetheless, it was found to apply for specific cases, such as periodic solutions^{3,4,5,6,7,8,9}. The main property that allows for the application of linear superposition principle to these specific nonlinear cases is the reduction of the nonlinear cross terms into linear ones, which then combine with other linear terms^{10,11,12}. This is possible only for certain types of solutions with this property.

Here, we will explicitly show a different method allowing for suitable linear combinations of special decomposition solutions which leads to completely different types of solutions. What is even more remarkable is that the decompositions allow us to uncover unexpected relations between various different integrable systems. To illustrate our approach, we begin with the B-type Kadomtsev-Petviashvili equation (BKP).

The BKP equation

$$u_{xt} + (u_{x4} + 15uu_{x2} + 15u^3 - 15uv - 5u_{xy})_{xx} - 5u_{yy} = 0, \quad v_x = u_y, \quad (1)$$

where $u_x = \partial_x u$, $u_{x2} = \partial_x^2 u$, $u_{x3} = \partial_x^3 u$, ..., is one of the most mystery models possessing many elegant properties^{13,14}. For simplicity only, we consider the potential form of (1) by setting $u = w_x$,

$$w_{xt} = 5w_{yy} - (w_{x5} + 15w_x w_{x3} + 15w_x^3 - 15w_x w_y - 5w_{xxy})_x \equiv K_x. \quad (2)$$

When u is y independent or $u_y \sim au_x$, the BKP equation (1) is returned to the known Sawada-Kotera (SK) equation^{15,16,17} which can be widely applied to many physical fields where the KdV equation is valid and the higher-order perturbations need to be considered¹⁸.

In the next section, the potential BKP (PBKP) hierarchy is briefly rewritten down employing the mastersymmetry method (MM)^{19,20,21} which is equivalent to the formal series symmetry approach (FSSA)^{22,23,24}. Then, we detail the steps leading to some special types of decompositions of the PBKP hierarchy in Section 3. These decompositions provide simple and interesting relationships for classical integrable systems and are fundamental for subsequent construction of linear superposition solutions. In Section 4, we show that suitable linear superpositions of some special decomposition solutions obtained in Section 3 of the PBKP hierarchy are still in fact new solutions of the same equations. Using these linear superposition formulas, we find several types of exact solutions including $m+n$ solitons, n solitons with periodic cnoidal background waves, n soliton solutions combined with soliton-cnoidal wave interaction solutions, and combination solutions of two different modified Schwarzian KdV waves, etc. of the fifth PBKP equation. In section 5, we present that similar special types of decompositions and linear superpositions still work for the dispersionless PBKP (dPBKP) hierarchy. Section 6 is devoted to conclusions.

2 | PBKP HIERARCHY VIA MASTERSYMMETRY METHOD

A symmetry, σ , of the PBKP equation (2) is defined as a solution of

$$\sigma_{xt} = \partial_x K' \sigma \equiv \partial_x \left(5\partial_x^{-1} \partial_y^2 - \partial_x^5 - 15w_{x3} \partial_x - 15w_x \partial_x^3 - 45w_x^2 \partial_x + 15w_x \partial_y + 15w_y \partial_x + 5\partial_x^2 \partial_y \right) \sigma, \quad (3)$$

which means (2) is invariant under the infinitesimal transformation $w \rightarrow w + \epsilon \sigma$ with infinitesimal parameter ϵ . If σ is not explicitly t dependent, then the symmetry definition equation (3) is equivalent to

$$K_{[.]} \sigma = 0, \quad (4)$$

where the commutate operator $K_{[.]}$ is defined as

$$K_{[.]f} \equiv K'f - f'K \equiv \lim_{\epsilon \rightarrow 0} \frac{d}{d\epsilon} [K(w + \epsilon f) - f(w + \epsilon K)] \quad (5)$$

for arbitrary f .

Conjecture 1.

$$K_{2n-1} = \frac{1}{3 \cdot 5^n n!} K_{[.]^n} y^n, \quad n = 1, 2, \dots, \infty \quad (6)$$

are all symmetries of the PBKP equation (2), that means y^n is a master symmetries for all positive integers $n \geq 1$ and $K_{[.]^{n+1}} y^n = 0$.

According to conjecture 1, the PBKP hierarchy can be written as

$$w_t = K_{2n-1} = \frac{1}{3 \cdot 5^n n!} K_{[.]^n} y^n, \quad n = 1, 2, \dots, \infty. \quad (7)$$

After finishing some tedious calculations, we have proved the conjecture 1 for $n = 1, 2, 3, 4$ and 5 with the following explicit forms,

$$w_{t_1} = K_1 = \frac{1}{15} K_{[.]} y = w_x, \quad (8)$$

$$w_{t_3} = K_3 = \frac{1}{150} K_{[.]^2} y^2 = 3w_y, \quad (9)$$

$$w_{t_5} = K_5 = \frac{1}{2250} K_{[.]^3} y^3 = 5\partial_x^{-1} w_{y2} - w_{x5} - 15w_x w_{x3} - 15w_x^3 + 5w_{yx2} + 15w_x w_y = K, \quad (10)$$

$$\begin{aligned}
w_{t_7} &= K_7 = \frac{1}{45000} K_{[.]}^4 y^4 \\
&= -w_{x7} - 21w_{x3}^2 - 21w_{x2}w_{x4} - 21w_x w_{x5} + 21\partial_x^{-1}(w_x w_{yy}) + 7\partial_x^{-2}w_{y3} + 42w_{x2}w_{xy} \\
&\quad - 42w_{x3}w_y + 42\partial_x^{-1}(w_{x4}w_y) - 63\partial_x^{-1}(w_x^2 w_{xy}) - 21w_x w_{yx2} - 63w_x w_{xx}^2 - 126w_x^2 w_{x3} + 21w_{xy2} \\
&\quad + 21w_x \partial_x^{-1}w_{y2} - 63w_x^4 + 63/2w_y^2,
\end{aligned} \tag{11}$$

$$\begin{aligned}
w_{t_9} &= K_9 = \frac{1}{1125000} K_{[.]}^5 y^5 \\
&= 9\partial_x^{-2}\{\partial_x^{-1}w_{y4} - w_{yx8} + 6w_{x2y3} - 270w_x w_{x2}^2 w_y + 3w_{x5y2} + 36w_{yx2}^2 + 9w_x w_{x3y2} + 9w_x w_{y3} \\
&\quad + 3(10w_{xy2} - 105w_{x2}w_{xy} - 111w_x w_{yx2} - f)w_{x3} - 9(39w_x w_{x2} + 5w_{xy})w_{yx3} - 18w_x w_{yx6} \\
&\quad + 27w_x w_{xy}^2 + 15w_{xy2}w_y - 3w_{x7}w_y - 9(21w_x w_{xy} + 10w_{x2}w_y - w_{y2} + 7w_{yx3})w_{x4} + 36w_{x2}w_{x2y2} \\
&\quad - 45w_{x2}w_{yx5} + 9w_{x2}\partial_x^{-1}w_{y3} - 18w_{x6}w_{xy} + 27w_{xy}w_{y2} + 9[2(5w_y - 27w_x^2)w_{xy} + 3w_x w_{y2}]w_{x2} \\
&\quad + 3(5w_y - 27w_x^2 - 20w_{x3})w_{yx4} - 45(w_x w_y + w_{yx2})w_{x5} - 63w_{x3}^2 w_y \\
&\quad + 9w_{yx2}(\partial_x^{-1}w_{y2} - 12w_x^3 + 5w_x w_y - 27w_{x2}^2)\}, \\
f_x &= 6(15w_x w_{x2} - 2w_{xy} - w_{x4})w_y + 54w_x^2 w_{xy} - \partial_x^{-1}w_{y3} - 3(w_x w_y)_y.
\end{aligned} \tag{12}$$

It should be mentioned that the ninth-order PBKP equation (12) is also only a (2+1)-dimensional extension of the seventh-order Sawada-Kotera (SK) equation. Both equations (11) and (12) will be reduced back to the seventh-order SK equation when we take $w_y = aw_x$.

Though the conjecture is difficult to prove for general n , its formal extended form can be restrictly proved by using FSSA^{22,23,24}.

Proposition 1. The PBKP equation (2) possesses formal series symmetries

$$\sigma_{2n-1} = \frac{1}{3 \cdot 5^n n!} \sum_{k=0}^{\infty} f^{(n-k)} K_{[.]}^k y^n, \quad n = 1, 2, \dots, \infty, \tag{13}$$

where f is an arbitrary function of t .

The correctness of proposition 1 has been proved for general t -independent K by one of the present authors (Lou) in several earlier works^{22,23,24,25}. Thus, conjecture 1 is equivalent to assuming that the formal series (13) is truncated up to n and the special function $f = 1$.

3 | DECOMPOSITIONS OF THE PBKP HIERARCHY

By using the symmetry constraint method to the Lax pair of the BKP equation, it is known that the BKP hierarchy can be solved by decomposing the BKP hierarchy to the KdV flows²¹. After finishing some tedious calculations, we find that for every equation of the PBKP (and then BKP) hierarchy, there are three consistent decompositions from the PBKP hierarchy to potential KdV (PKdV) flows. More different ways to relate BKP and KdV hierarchies with reductions were described in^{26,27,28}. Another point worth bringing up is that in addition to the KdV hierarchy, we have found relationships between other classical systems and the BKP hierarchy. The relationships between these various models are illustrated below.

3.1 | Decompositions of the fifth-order PBKP equation (2)

Usually, for a higher dimensional integrable system, one can formally separate some variables by using symmetry constraints on its Lax pair. For PBKP equation (2), the Lax pair possesses the form

$$\psi_y + \psi_{x3} + 3w_x \psi_x = 0, \tag{14}$$

$$\psi_t - 9\psi_{x5} - 45w_x \psi_{x3} - 45w_{x2} \psi_{x2} - 15(2w_{x3} + 3w_x^2 - w_y) \psi_x = 0. \tag{15}$$

It is not difficult to verify that ψ_x is a special symmetry of the PBKP equation, i.e., $\sigma = \psi_x$ is a solution of (3). Thus, substituting the symmetry constraint $\psi_x = w_x$, i.e.,

$$\psi = w, \tag{16}$$

into the Lax pair (14) and (15), we have

$$w_y + w_{x3} + 3w_x^2 = 0, \quad (17)$$

$$w_t - 9w_{x5} - 45(2w_{x3} + 3w_x^2)w_x = 0. \quad (18)$$

Because equation (17) is not explicitly t -dependent and (18) is not explicitly y -dependent, this kind of decomposition is called nonlinearization²⁹ and also known as variable separation approach³⁰. Since w is still a function of y and t , we call this kind of method as formally variable separation approach (FVSA)^{31,32}. The separation equations (17) and (18) are nothing but the KdV and the fifth-order KdV equations, respectively. That means if w is a solution of the KdV and fifth-order KdV decompositions (17) and (18), then it is also a solution of the PBKP equation (2).

In fact the FVSA can be applied to nonlinear systems irrelative the existence of Lax pairs³¹, since it can be conjectured that any (2+1)-dimensional nonlinear system has a solution decomposed in the form

$$w_y = F(w, w_x, w_{x2}, \dots, w_{xm}), \quad (19)$$

$$w_t = G(w, w_x, w_{x2}, \dots, w_{xn}) \quad (20)$$

with the consistent condition

$$w_{yt} - w_{ty} = [F, G] = 0 \quad (21)$$

preserved. So it is not necessary to stipulate the exact forms of the functions F and G , it is only necessary to specify that (19) and (20) satisfy the PBKP equation (2) and the consistent condition (21). Then a direct calculation shows the following decomposition theorem.

Proposition 2. If w_1, w_2, w_3, w_4, w_5 and w_6 are the solutions of the following decomposed systems

$$\begin{cases} w_{1y} = (\Phi_1 + c)w_{1x} + c_1, \quad \Phi_1 \equiv \partial_x^2 + 4w_{1x} - 2\partial_x^{-1}w_{1x2}, \\ w_{1t} = (9\Phi_1^2 + 15c\Phi_1 + 5(c^2 + 3c_1))w_{1x}, \end{cases} \quad (22)$$

$$\begin{cases} w_{2y} = \Phi_2 w_{2x} + c_1, \quad \Phi_2 \equiv \partial_x^2 + 2w_{2x} - \partial_x^{-1}w_{2x2}, \\ w_{2t} = (9\Phi_2^2 + 15c_1)w_{2x}, \end{cases} \quad (23)$$

$$\begin{cases} w_{3y} = -\frac{1}{2}(\Phi_3 - 2c)w_{3x} + c_1, \quad \Phi_3 \equiv \partial_x^2 + 4w_{3x} - 2\partial_x^{-1}w_{3x2}, \\ w_{3t} = \left(5c^2 + 15c_1 - \frac{9}{4}\Phi_3^2\right)w_{3x}, \end{cases} \quad (24)$$

$$\begin{cases} w_{4y} = w_{4x3} - \frac{3}{4}\frac{w_{4x2}^2}{W^2} + \frac{3}{2}W^4 - c_1 = \Phi_4 w_{4x} + c_1, \quad W^2 \equiv w_{4x} + c, \\ w_{4t} = 9\Phi_4^2 w_{4x} + 15c_1 w_{4x}, \quad \Phi_4 \equiv \partial_x^{-1}W\partial_x^2 W^{-1}\partial_x + W^2 + \partial_x^{-1}W^2\partial_x, \end{cases} \quad (25)$$

$$\begin{cases} w_{5y} = \frac{1}{4}(4\Phi_5^2 + 6c\Phi_5 + 3c^2)w_{5x}, \quad \Phi_5 \equiv \partial_x + \frac{1}{2}w_5 + \frac{1}{2}w_{5x}\partial_x^{-1}, \\ w_{5t} = \frac{9}{16}(16\Phi_5^4 + 40c\Phi_5^3 + 40c^2\Phi_5^2 + 20c^3\Phi_5 + 5c^4)w_{5x}, \end{cases} \quad (26)$$

$$\begin{cases} w_{6y} = cw_{6x} + c_1, \\ w_{6t} = -w_{6x5} + 5(c - 3w_{6x})w_{6x3} + 15cw_{6x}^2 - 15w_{6x}^3 + 5(c^2 + 3c_1)w_{6x}, \end{cases} \quad (27)$$

then w_1, w_2, w_3, w_4, w_5 and w_6 are all solutions of the PBKP equation (2).

Proof. Substituting (19) and (20) with $m > 3$ into the PBKP equation (2) and the decomposition consistent condition (21), and finishing some tedious calculations, one can find that there is no possible decomposition (19) and (20) with $m > 3$. Thus, we just

take $m = 3$ and then $n = 2m - 1 = 5$ in the decomposition relations (19) and (20). Substituting (19) and (20) into (2), we have

$$w_{x_6}(1 + 5F_{x_3}^2 - 5F_{x_3} + G_{x_5}) + W(x_0, x_1, \dots, x_5) = 0, \quad (28)$$

where $F = F(w, w_x, w_{x_2}, w_{x_3}) \equiv F(x_0, x_1, x_2, x_3)$, $G = G(w, w_x, w_{x_2}, w_{x_3}, w_{x_4}, w_{x_5}) \equiv G(x_0, x_1, \dots, x_5)$ and $W = W(x_0, x_1, \dots, x_5)$ is a complicated expression of x_0, x_1, \dots, x_5 . Vanishing coefficient of w_{x_6} , we have

$$G = (5F_{x_3} - 5F_{x_3}^2 - 1)x_5 + G_1(x_0, x_1, \dots, x_4), \quad (29)$$

where $G_1 = G_1(x_0, x_1, \dots, x_4)$ is a function of $\{x_0, x_1, \dots, x_4\}$. By using the relation (29), (28) is changed to

$$w_{x_5} [G_{1x_4} - 5(F_{x_3} + 2)(x_4 F_{x_3 x_3} + x_3 F_{x_2 x_3} + x_2 F_{x_1 x_3} + x_1 F_{x_0 x_3}) + 10F_{x_2} F_{x_3} + 5F_{x_2}^2] + W_1 = 0, \quad (30)$$

where $W_1 = W_1(x_0, x_1, \dots, x_4)$ is w_{x_5} independent. Eliminating the coefficient of w_{x_5} yields

$$G_1 = 5(F_{x_3} + 2) \left(\frac{1}{2} x_4 F_{x_3 x_3} + x_3 F_{x_2 x_3} + x_2 F_{x_1 x_3} + x_1 F_{x_0 x_3} \right) x_4 + 5x_4 F_{x_2} (2F_{x_3} + 1) + G_2, \quad (31)$$

with $G_2 \equiv G_2(x_0, x_1, x_2, x_3)$ being a function of $\{x_0, x_1, x_2, x_3\}$.

Similarly, substituting the decomposition (19) and (20) with $m = 3$, $n = 5$ and the results (29) and (31) into the consistent condition (21), we have

$$5w_{x_7}(1 - F_{x_3})^2(x_1 F_{x_0 x_3} + x_2 F_{x_1 x_3} + x_3 F_{x_2 x_3} + x_4 F_{x_3 x_3}) + \Gamma = 0, \quad (32)$$

where $\Gamma = \Gamma(x_0, \dots, x_6)$ is a w_{x_7} independent function of the lower-order differentiations of w with respect to x .

Vanishing the coefficient of $w_{x_7} w_{x_4}$ in (32), we get

$$F = F_1(x_0, x_1, x_2)x_3 + H(x_0, x_1, x_2). \quad (33)$$

Substituting (33) into (32) and requiring the coefficient of w_{x_7} being zero results $F_1(x_0, x_1, x_2) = a$, a constant. Thus, we have,

$$F = ax_3 + H(x_0, x_1, x_2). \quad (34)$$

Because of the simplification (34), (30) is simplified to

$$w_{x_4} [G_{2x_3} - 5(a + 2)x_3 H_{x_2 x_2} - 5(a + 2)x_2 H_{x_1 x_2} + 5x_1(3 - aH_{x_0 x_2} - 3a - 2H_{x_0 x_2}) - 10aH_{x_1} - 5H_{x_2}^2 - 5H_{x_1}] + W_2 = 0 \quad (35)$$

with $W_2 = W_2(x_0, x_1, x_2)$. Vanishing the coefficient of w_{x_4} in (35) leads to

$$G_2 = 5 \left[\frac{a+2}{2} x_3 H_{x_2 x_2} + (a+2)x_2 H_{x_1 x_2} - x_1(3 - aH_{x_0 x_2} - 3a - 2H_{x_0 x_2}) + 2aH_{x_1} + H_{x_2}^2 + H_{x_1} \right] x_3 + J = 0, \quad (36)$$

where $J = J(x_0, x_1, x_2)$. Up to now, the decomposition relation is simplified to

$$w_y = aw_{x_3} + H, \quad (37)$$

$$w_t = (5a^2 + 5a - 1)w_{x_5} + 5(2a + 1)w_{x_4} H_{x_2} + 5 \left[\frac{a+2}{2} x_3 H_{x_2 x_2} + (a+2)x_2 H_{x_1 x_2} - x_1(3 - aH_{x_0 x_2} - 3a - 2H_{x_0 x_2}) + 2aH_{x_1} + H_{x_2}^2 + H_{x_1} \right] x_3 + J = 0, \quad (38)$$

with one constant a and two undetermined functions $H = H(x_0, x_1, x_2)$ and $J = J(x_0, x_1, x_2)$.

Inserting (37) and (38) into the PBKP equation (2) and the consistent condition (21), then, vanishing the coefficients of w_{x_k} for $k \geq 3$ leaves the determining equations on $\{a, H, J\}$,

$$\begin{aligned} aH_{x_2 x_2 x_2} &= 0, (1-a)[2(a+2)(x_2 H_{x_1 x_2 x_2} + x_1 H_{x_0 x_2 x_2}) + (H_{x_2}^2 + 3H_{x_1})_{x_2}] = 0, \\ (1-a)[(a+2)(2x_1 x_2 H_{x_0 x_1 x_2} + x_1^2 H_{x_0 x_0 x_2} + x_2^2 H_{x_1 x_1 x_2}) + (1-a)(x_2 H_{x_1 x_1} + x_1 H_{x_0 x_1}) \\ &+ [(a+2)x_2 - 2x_1 H_{x_2}] H_{x_0 x_2} - 2x_2 H_{x_2} H_{x_1 x_2} - 9ax_2] = 0, \\ (1-a)(x_1 H_{x_0 x_2} + x_2 H_{x_1 x_2}) &= 0, (1-a)H_{x_2 x_2} = 0, (5a^2 - 2)\partial_{x_2}^5 H = 0, \\ (2H_{x_2} - 3ax_1 \partial_{x_0} - 3ax_2 \partial_{x_1}) H_{x_2 x_2} &= 0, (7a^2 - 4)\partial_{x_2}^4 H = 0, \end{aligned} \quad (39)$$

$$\begin{aligned}
& [(5a^2 - 2)(x_2 \partial_{x_1} + x_1 \partial_{x_0}) \partial_{x_2} + (7a^2 - 4) \partial_{x_1} - a H_{x_2} \partial_{x_2} + (6 - 9a) H_{x_2 x_2}] H_{x_2 x_2 x_2} = 0, \\
& [3(7a^2 - 4)(x_2 \partial_{x_1} + x_1 \partial_{x_0}) + 2(1 - 4a) H_{x_2}] H_{x_2 x_2 x_2} = 0, \\
& J_{x_2} - 5(2a + 1)(2x_1 x_2 \partial_{x_1 x_0} + x_1^2 \partial_{x_0}^2 + x_2^2 \partial_{x_1}^2 + x_2 \partial_{x_0}) H_{x_2} \\
& \quad - 5(a + 2)(x_2 \partial_{x_1} + x_1 \partial_{x_0}) H_{x_1} - 5(1 + 2a) H_{x_0} - 5 H_{x_2} (3x_1 + 2H_{x_1}) = 0,
\end{aligned} \tag{40}$$

$$\begin{aligned}
& 5(a + 1)(x_1^3 \partial_{x_0}^3 + x_2^3 \partial_{x_1}^3 + 3x_2^2 x_1 \partial_{x_1}^2 \partial_{x_0} + 3x_2 x_1^2 \partial_{x_1} \partial_{x_0}^2 + 3x_2 x_1 \partial_{x_0}^2 + 3x_2^2 \partial_{x_1 x_0}) H \\
& \quad + 5 H_{x_2} (x_1^2 \partial_{x_0}^2 + x_2^2 \partial_{x_1}^2 + 2x_2 x_1 \partial_{x_1 x_0} + x_2 \partial_{x_0}) H + 5(x_1 H_{x_1} + 3x_1^2 + H) H_{x_0} \\
& \quad + 5x_2 H_{x_1} (H_{x_1} + 3x_1) - x_1 J_{x_0} - x_2 J_{x_1} - 15x_2 (3x_1^2 - H) = 0, \\
& -3(x_1^2 \partial_{x_0}^2 + x_2^2 \partial_{x_1}^2 + 2x_1 x_2 \partial_{x_1 x_0} + x_2 \partial_{x_0}) \partial_{x_2}^3 H + (3x_1 \partial_{x_0} + 3x_2 \partial_{x_1} - 2H_{x_2}) H_{x_1 x_2 x_2} + 6H_{x_1 x_1 x_2} \\
& \quad - 10H_{x_2 x_2} H_{x_1 x_2} + 2(3x_1 - H_{x_1} + H_{x_2}^2) \partial_{x_2}^3 H = 0,
\end{aligned} \tag{41}$$

$$\begin{aligned}
& (a - 1)^2 (9 - H_{x_1 x_1} + x_1 H_{x_0 x_1 x_2} + x_2 H_{x_1 x_1 x_2}) + (1 - a)(2H_{x_2} H_{x_1 x_2} + 3H_{x_1 x_1} + 3H_{x_0 x_2} - 9) \\
& \quad + 9x_1 H_{x_0 x_1 x_2} + 9x_2 H_{x_1 x_1 x_2} - 6x_1 H_{x_0 x_2} H_{x_2 x_2} - 6H_{x_1 x_2} (x_2 H_{x_2})_{x_2} \\
& \quad + H_{x_2 x_2} (9x_1 - 3H_{x_1} + H_{x_2}^2) = 0,
\end{aligned} \tag{42}$$

$$\begin{aligned}
& -(a^2 + 2)(x_1^3 \partial_{x_0}^3 + 3x_2^2 x_1 \partial_{x_0}^2 \partial_{x_1} + 3x_1 x_2^2 \partial_{x_0} \partial_{x_1}^2 + x_2^3 \partial_{x_1}^3 + 3x_1 x_2 \partial_{x_0}^2 + 3x_2^2 \partial_{x_0 x_1}) H_{x_2} \\
& \quad + (2a + 1)[(x_1^2 H_{x_0 x_0} + 2x_1 x_2 H_{x_0 x_1} + x_2^2 H_{x_1 x_1} + x_2 H_{x_0}) H_{x_2 x_2} + (x_1^2 H_{x_0 x_0 x_2} + 2x_1 x_2 H_{x_0 x_1 x_2} \\
& \quad + x_2^2 H_{x_1 x_1 x_2} + x_2 H_{x_0 x_2}) H_{x_2} + x_1 H_{x_0} H_{x_1 x_2}] + (1 - a)[x_1 (a - 1) H_{x_0 x_0} - (a + 2)(x_1^2 H_{x_0 x_0 x_1} \\
& \quad + 2x_1 x_2 H_{x_0 x_1 x_1} + x_2^2 H_{x_1 x_1 x_1}) + 2x_2 H_{x_2} H_{x_1 x_1} + 2x_1 H_{x_2} H_{x_0 x_1} + 2x_1 H_{x_1} H_{x_0 x_2} - 3x_2 H_{x_0 x_1}] \\
& \quad + 3(2 - a)(x_1 H_{x_0 x_2} + x_2 H_{x_1 x_2})^2 - x_2 (H_{x_2}^2 + 9x_1 - 3H_{x_1}) H_{x_1 x_2} - x_1 H_{x_0 x_2} H_{x_2}^2 - 15a x_2 H_{x_2} \\
& \quad + 2a H H_{x_0 x_2} - 9x_1^2 H_{x_0 x_2} + 6x_2 H_{x_2} + H H_{x_0 x_2} = 0,
\end{aligned} \tag{43}$$

$$\begin{aligned}
& 4(a - 5a^2 - 5)x_2 H_{x_0 x_1 x_2} + (7 - a)(4x_1 x_2 H_{x_2 x_2} H_{x_0 x_1 x_2} + 2x_1^2 H_{x_2 x_2} H_{x_0 x_0 x_2} + 2x_2^2 H_{x_2 x_2} H_{x_1 x_1 x_2} \\
& \quad + 2x_2 H_{x_0 x_2} H_{x_2 x_2}) + (4 - a)(x_1 H_{x_1} H_{x_0 x_2 x_2} + 4x_2 H_{x_2 x_2} H_{x_1 x_1} + 4x_1 H_{x_2 x_2} H_{x_0 x_1} + 6x_2 H_{x_1 x_2}^2 \\
& \quad + 6x_1 H_{x_0 x_2} H_{x_1 x_2}) + (1 - a)[2H_{x_2} H_{x_1 x_1} - 2(a + 2)H_{x_0 x_1} - 6(a + 1)(x_1 \partial_{x_0} \partial_{x_1}^2 - x_2 \partial_{x_1}^3) H] \\
& \quad - 6(a^2 + 2)(x_2^2 H_{x_1 x_1 x_1 x_2} + 2x_1 x_2 H_{x_0 x_1 x_1 x_2} + x_1^2 H_{x_0 x_0 x_1 x_2}) + (2 + a)[(a - 2)(x_2^3 H_{x_1 x_1 x_1 x_2 x_2} \\
& \quad + x_1^3 H_{x_0 x_0 x_0 x_2 x_2} + 3x_1^2 x_2 H_{x_0 x_0 x_1 x_2 x_2} + 3x_1 x_2^2 H_{x_0 x_1 x_1 x_2 x_2} + 3x_1 x_2 H_{x_0 x_0 x_2 x_2} + 3x_2^2 H_{x_0 x_1 x_2 x_2}) \\
& \quad + 2x_1 H_{x_2} H_{x_0 x_1 x_2} + 2x_2 H_{x_2} H_{x_1 x_1 x_2} + x_1 H_{x_0} H_{x_1 x_2 x_2} + (x_1^2 H_{x_0 x_0} + x_2^2 H_{x_1 x_1} + 2x_1 x_2 H_{x_0 x_1} \\
& \quad + x_2 H_{x_0}) H_{x_2 x_2 x_2}] + 2(10 - 7a)(x_2 H_{x_1 x_2} + x_1 H_{x_0 x_2})(x_1 H_{x_0 x_2 x_2} + x_2 H_{x_1 x_2 x_2}) \\
& \quad + 3a H_{x_2} (x_1^2 \partial_{x_0}^2 + 2x_1 x_2 \partial_{x_0 x_1} + x_2^2 \partial_{x_1}^2) H_{x_2 x_2} - 2x_1 (7a^2 - 2a + 4) H_{x_0 x_0 x_2} - (2x_1 H_{x_2}^2 - 3ax_2 H_{x_2} \\
& \quad - aH + 18x_1^2 - 2H) H_{x_0 x_2 x_2} - (2x_2 H_{x_2}^2 - 6x_2 H_{x_1} + 18x_1 x_2) H_{x_1 x_2 x_2} + 4(1 + 2a) H_{x_2} H_{x_0 x_2} \\
& \quad - (18x_1 - 2H_{x_2}^2 - 6H_{x_1}) H_{x_1 x_2} + [(4 + 5a) H_{x_0} - 18ax_2] H_{x_2 x_2} - 18a H_{x_2} = 0,
\end{aligned} \tag{44}$$

$$\begin{aligned}
& (1 - 5a^2 - 5a)(10x_2^4\partial_{x_0}\partial_{x_1}^3 + 10x_2x_1^3\partial_{x_0}^4 + 30x_1x_2^3\partial_{x_0}^2\partial_{x_1}^2 + 30x_1^2x_2^2\partial_{x_0}\partial_{x_1}^3 + x_2^5\partial_{x_1}^5 + x_1^5\partial_{x_0}^5 \\
& + 5x_1x_2^4\partial_{x_0}\partial_{x_1}^4 + 5x_2x_1^4\partial_{x_0}^4\partial_{x_1} + 10x_2^2x_1^3\partial_{x_0}^3\partial_{x_1}^2 + 10x_1^2x_2^2\partial_{x_0}^2\partial_{x_1}^3)H \\
& - 5(2a + 1)(2x_1x_2H_{x_0x_1x_2} + x_1^2H_{x_0x_0x_2} + x_2^2H_{x_1x_1x_2})(x_1^2H_{x_0x_0} + x_2^2H_{x_1x_1} + 2x_2x_1H_{x_0x_1} + x_2H_{x_0}) \\
& - 5x_1^2(ax_2H_{x_1x_2} + ax_1H_{x_0x_2} + H_{x_2}^2 + 2x_2H_{x_1x_2} + 2x_1H_{x_0x_2})(x_1\partial_{x_0}^3 + 3x_2\partial_{x_0}^2\partial_{x_1})H \\
& - 5x_2^2[(a + 2)(x_2H_{x_1x_2} + x_1H_{x_0x_2}) + (2a + 1)H_{x_1} + H_{x_2}^2 - 3(1 - a)x_1](x_2\partial_{x_1}^3 + 3x_1\partial_{x_1}^2\partial_{x_0})H \\
& + 5(2a + 1)\{-H_{x_2}[6x_2^3\partial_{x_0}\partial_{x_1}^2 + x_1^4\partial_{x_0}^4 + x_2^4\partial_{x_1}^4 + 6x_1^2x_2^2\partial_{x_0}^3 + 12x_1x_2^2\partial_{x_0}^2\partial_{x_1} + 4x_1x_2^2\partial_{x_0}\partial_{x_1}^3 \\
& + 6x_1^2x_2^2\partial_{x_0}^2\partial_{x_1}^2 + 4x_1^3x_2\partial_{x_0}^3\partial_{x_1}] - H_{x_1}[x_1^3\partial_{x_0}^3 + 3x_2x_1^2\partial_{x_0}^2\partial_{x_1} + 3x_2x_1\partial_{x_0}^2 + 3x_2^2\partial_{x_0}\partial_{x_1}] - x_2^3H_{x_0x_2}\partial_{x_1}^2 \\
& - H_{x_0}[x_1^2\partial_{x_0}^2 + x_2^2\partial_{x_0}\partial_{x_2} + x_2\partial_{x_0}]\}H - 5H_{x_0x_0}[(a + 2)x_1^3H_{x_0x_1} + x_1x_2(a + 2)(3x_2H_{x_1x_2} \\
& + x_1H_{x_1x_1}) + x_2x_1^2(5a + 7)H_{x_0x_2} + x_1H_{x_2}(3x_2H_{x_2} + 2x_1H_{x_1}) + 3(2ax_2^2 + x_1^3 + x_2^2)H_{x_2} \\
& - 9(1 - a)x_2x_1^2] - 15[x_1(5a^2x_2^2 + ax_1^3 + 5ax_2^2 - x_1^3 - x_2^2)\partial_{x_0}^3 + x_2(5a^2x_2^2 + 3ax_1^3 + 5ax_2^2 - 3x_1^3 \\
& - x_2^2)\partial_{x_0}^2\partial_{x_1}]H - 10(a + 2)x_2x_1^2H_{x_0x_1}^2 - 5H_{x_0x_1}[3(a + 2)(x_2^3H_{x_1x_2} + x_1x_2^2H_{x_1x_1}) \\
& + x_1x_2^2(7a + 8)H_{x_0x_2} + (5a + 4)x_1x_2H_{x_0} + x_2H_{x_2}(3x_2H_{x_2} + 4x_1H_{x_1} + 6x_1^2) - 9(1 - a)x_1x_2^2] \\
& - 5(a + 2)x_2^3H_{x_1x_1}^2 + a(3x_2x_1^2\partial_{x_0}^2\partial_{x_1} + 3x_2^2x_1\partial_{x_0}\partial_{x_1}^2 + x_2^3\partial_{x_1}^3 + x_1^3\partial_{x_0}^3)J - 5x_2^2[(2H_{x_1} + 3x_1)H_{x_2} \\
& + 3(a + 1)H_{x_0}]H_{x_1x_1} + x_1(x_1H_{x_2} + 3ax_2)J_{x_0x_0} + x_2(2x_1H_{x_2} + 3ax_2)J_{x_0x_1} + x_2^2H_{x_2}J_{x_1x_1} \\
& - (10x_2H_{x_2}H_{x_1} + 15x_1x_2H_{x_2} + x_1J_{x_1} - J)H_{x_0} + J_{x_0}(x_2H_{x_2} + x_1H_{x_1} - H) = 0, \tag{45}
\end{aligned}$$

$$\begin{aligned}
& -(a + 2)(x_1^2H_{x_0x_0} + x_2^2H_{x_1x_1} + 2x_2x_1H_{x_0x_1} + x_2H_{x_0})(x_2\partial_{x_1} + x_1\partial_{x_0})H_{x_2x_2} - (a + 8)(x_2H_{x_1x_2} \\
& + x_1H_{x_0x_2})(2x_1x_2\partial_{x_0x_1} + x_1^2\partial_{x_0}^2 + x_2^2\partial_{x_1}^2 + x_2\partial_{x_0})H_{x_2} - [x_1^2H_{x_2}^2 + (1 - a)x_1^2H_{x_1} \\
& + 3ax_1x_2H_{x_2} + x_1(a + 2)H - 3(2a^2 + 1)x_2^2 - 9x_1^3]H_{x_0x_0x_2} - [(a + 2)x_1^2H_{x_0} + (4 - a)x_1x_2H_{x_1} \\
& + 2x_1x_2H_{x_2}^2 + 3ax_2^2H_{x_2} + x_2(aH - 18x_1^2 + 2H)]H_{x_0x_1x_2} + (2a^2 + 1)[6x_1^2x_2\partial_{x_0}^3 + 6x_2^3\partial_{x_0}\partial_{x_1}^2 \\
& + 12x_1x_2^2\partial_{x_0}^2\partial_{x_1} + x_1^4\partial_{x_0}^4 + 4x_1x_2^3\partial_{x_0}\partial_{x_1}^3 + 6x_1^2x_2^2\partial_{x_0}^2\partial_{x_1}^2 + 4x_2x_1^3\partial_{x_0}^3\partial_{x_1} + x_2^4\partial_{x_1}^4]H_{x_2} \\
& + (1 - a)[(a + 1)(6x_1x_2^2\partial_{x_1}^3\partial_{x_0} + 6x_1^2x_2\partial_{x_0}^2\partial_{x_1}^2 + 2x_1^3\partial_{x_0}^3\partial_{x_1} + 2x_2^3\partial_{x_1}^4)H \\
& - H_{x_2}(x_2^3\partial_{x_1}^3 + 2x_1x_2\partial_{x_0}\partial_{x_1}^2 + x_1^2\partial_{x_0}^2\partial_{x_1} + 2x_1\partial_{x_0}^2 + 3x_2\partial_{x_0x_1})H - x_1H_{x_0}H_{x_1x_1}] \\
& - (10 - 7a)x_1H_{x_0x_2}(x_2H_{x_1x_1} + x_1H_{x_0x_1}) + 2x_1^2H_{x_2}H_{x_0x_2}^2 + [9(2a + 1)x_1x_2 + 4x_2x_1H_{x_2}H_{x_1x_2} \\
& - (5a + 4)x_1H_{x_0} + (2x_1H_{x_2} + 3x_2)H_{x_1} - x_2H_{x_2}^2 - 2H_{x_2}H]H_{x_0x_2} - 3ax_1x_2^2H_{x_2}H_{x_0x_1x_1x_2} \\
& + (a + 2)x_2^2(4 - 3x_1H_{x_2x_2} - a)H_{x_0x_1x_1} - [x_2(4x_1H_{x_2} + 3(a + 2)x_2)H_{x_2x_2} + (16 - 7a)x_1x_2H_{x_1x_2} \\
& + (a + 2)x_1H_{x_1} - x_1H_{x_2}^2 - (aH + 9x_1^2 - H)]H_{x_0x_1} + x_2H_{x_2}^2(6 - x_2H_{x_1x_1x_2} + H_{x_1x_1}) \\
& + H_{x_2}(2x_2^2H_{x_1x_2}\partial_{x_1} - ax_1^3\partial_{x_0}^3 - ax_2^3\partial_{x_1}^3 - 3ax_2x_1^2\partial_{x_0}^2\partial_{x_1} - 2x_1H_{x_0}\partial_{x_1})H_{x_2} - H_{x_2x_2}[2(x_1^2H_{x_0x_0} \\
& + x_2^2H_{x_1x_1} + x_2H_{x_0})H_{x_2} + (a + 2)(x_1^3\partial_{x_0}^3 + x_2^3\partial_{x_1}^3 + 3x_2x_1^2\partial_{x_0}^2\partial_{x_1} + 3x_2x_1\partial_{x_0}^2)H] \\
& + (5a^2 + 2a + 2)x_1^2H_{x_0x_0x_0} + 2(2a^2 + 2a + 5)x_1x_2H_{x_0x_0x_1} - [(a + 2)x_1x_2H_{x_0} + 3x_2^2H_{x_1} \\
& - 9x_2^2x_1]H_{x_1x_1x_2} + [(5a^2 + 2a + 2)x_2 - 3x_1^2H_{x_1x_2}]H_{x_0x_0} + [(7a - 13)x_2^2H_{x_1x_2} - 3x_2H_{x_1} \\
& + 9x_1x_2]H_{x_1x_1} - 54ax_2x_1 + 3[(5a + 1)x_1 + 2aH_{x_1}]H_{x_0} + 18ax_2H_{x_1} - [(4a + 5)x_2H_{x_0} \\
& - 18ax_2^2]H_{x_1x_2} - 3x_1x_2(a + 2)H_{x_0x_0}H_{x_2x_2} - 3aJ_{x_0}/5 = 0. \tag{46}
\end{aligned}$$

From the determining equation (39), we know that the determining equations (39)-(46) should be solved in three separated cases for $a = 0$, $a = 1$ and $a \neq 0, 1$, respectively. Now, it is easy to finish the final work by solving (39)-(46).

Case 1. $a = 0$. In this simple case, the final solutions of H and J read

$$H = cx_1 + c_1 = cw_x + c_1, \quad J = 15cw_x^2 - 15w_x^3 + (5c^2 + 15c_1)w_x. \tag{47}$$

Case 2. $a = 1$. In this case, one can find four different solutions

$$H = \frac{3}{2}x_1^2 - c_1, \quad J = \frac{45}{2}(x_2^2 + x_1^3) + 15c_1x_1, \quad x_0 = w, \quad x_1 = w_x, \quad x_2 = w_{x_2}, \tag{48}$$

$$H = cw_x + 3w_x^2 + c_1, \quad J = 90w_x^3 + 45cw_x^2 + 5c^2w_x + 15c_1w_x + 45w_{x2}^2, \quad (49)$$

$$H = c_1 - \frac{3w_{x2}^2}{4(c+w_x)} + \frac{3}{2}w_x^2, \quad J = 15c_1w_x - \frac{315w_{x2}^4}{16(c+w_x)^3} - \frac{45(w_x-2c)}{4(c+w_x)}w_{x2}^2 + \frac{45}{2}w_x^3, \quad (50)$$

$$H = \frac{3}{4}[2(w+c)w_{x2} + 2w_x^2 + (c+w)^2w_x], \quad (51)$$

$$J = \frac{45}{16}\{16w_{x2}^2 + [40(w+c)w_x + 4(w+c)^3]w_{x2} + 12w_x^3 + 12(w+c)^2w_x^2 + (w+c)^4w_x\}.$$

Case 3. $a \neq 0, 1$. In this case, we find that a should be fixed as $a = -\frac{1}{2}$ and the functions H and J are also fixed as

$$H = -\frac{3}{2}w_x^2 + cw_x + c_1, \quad J = -\frac{45}{4}w_{xx}^2 - \frac{45}{2}w_x^3 + 5c^2w_x + 15c_1w_x. \quad (52)$$

c and c_1 in all cases are arbitrary constants. By substituting these solutions (47)–(52) into the decomposition relations (37) and (38), proposition 2 is proved. \square

Proposition 2 provides simple and interesting relationships between classical integrable systems and PBKP equation (2). The first three decompositions (22)–(24) are all the KdV decompositions. The first equation of the fourth decomposition (25) is an integrable model which is a special form of the Svinolupov-Sokolov (SS) equation proposed in³³. The decomposition (26) is related to the so-called Sharma-Tasso-Olver (STO) equation (the third-order equation of the Burgers hierarchy)^{34,35}. The sixth decomposition is a natural SK decomposition. The unexpected linear superposition property of these decompositions continues to surprise us in Sec. 4.

3.2 | Decompositions of the seventh-order PBKP equation (11)

For the seventh-order PBKP equation, decompositions can be found in the similar way. We just list the results in the following proposition.

Proposition 3. The functions w_i , $i = 1, 2, \dots, 6$ satisfying the following decomposition systems

$$\begin{cases} w_{1y} = (\Phi_1 + c)w_{1x} + c_1, \\ w_{1t} = [27\Phi_1^3 + 63c\Phi_1^2 + (42c^2 + 63c_1)\Phi_1 + 7c(c^2 + 9c_1)]w_{1x}, \end{cases} \quad (53)$$

$$\begin{cases} w_{2y} = \Phi_2 w_{2x} + c_1, \\ w_{2t} = (27\Phi_2^3 + 63c_1\Phi_2)w_{2x}, \end{cases} \quad (54)$$

$$\begin{cases} w_{3y} = -\frac{1}{2}(\Phi_3 - 2c)w_{3x} + c_1, \\ w_{3t} = \left[\frac{27}{8}\Phi_3^3 - \frac{63}{4}c\Phi_3^2 + \frac{7}{2}(3c^2 - 9c_1)\Phi_3 + 7c(c^2 + 9c_1) \right] w_{3x}, \end{cases} \quad (55)$$

$$\begin{cases} w_{4y} = \Phi_4 w_{4x} + c_1, \\ w_{4t} = 27\Phi_4^3 w_{4x} + 63c_1\Phi_4 w_{4x}, \end{cases} \quad (56)$$

$$\begin{cases} w_{5y} = \frac{1}{4}(4\Phi_5^2 + 6c\Phi_5 + 3c^2)w_{5x}, \\ w_{5t} = \frac{27}{64}(64\Phi_5^6 + 224c\Phi_5^5 + 336c^2\Phi_5^4 + 280c^3\Phi_5^3 + 140c^4\Phi_5^2 + 42c^5\Phi_5 + 7c^6)w_{5x}, \end{cases} \quad (57)$$

$$\begin{cases} w_{6y} = cw_{6x} + c_1, \\ w_{6t} = 84c_1cw_{6x} + 7c^3w_{6x} + 21(3w_{6x}^2 + w_{6x3})c^2 - 21(w_{6x}^3 + w_{6x}w_{6x3} - w_{6x2}^2)c \\ \quad - 63w_{6x}(w_{6x}^3 + 2w_{6x}w_{6x3} + w_{6x2}^2) - 21(w_{6x}w_{6x5} + w_{6x2}w_{6x4} + w_{6x3}^2) - w_{6x7} \end{cases} \quad (58)$$

are all solutions of the seventh-order PBKP equation (11) with $t_7 = t$.

To prove proposition 3, one can directly substitute the decompositions to the seventh-order PBKP equation (11) and finish some integrations by parts. Here, we omit the details on the proof procedure. Same as in the fifth-order BKP equation, decompositions make it possible to find interrelations between classical integrable systems and the seventh-order PBKP equation. The first three decompositions are the usual potential KdV and seventh-order potential KdV reductions. The fourth decomposition is related to the special SS equation (25) and its seventh-order flow. The fifth decomposition is related to the third-order Burgers (STO) and seventh-order Burgers equations in potential forms. The sixth decomposition is simply the seventh-order potential SK reduction of the seventh-order PBKP equation.

3.3 | Decompositions of the ninth-order PBKP equation (12)

In the same way, we directly write down the decomposition theorem for the ninth-order PBKP equation (12) without detailed verifications.

Proposition 4. The functions w_i , $i = 1, \dots, 6$ provide solutions of the ninth-order PBKP equation (12) with $t_9 = t$ if they satisfy

$$\begin{cases} w_{1y} = (\Phi_1 + c)w_{1x} + c_1, & \mu \equiv \frac{9}{2}(2c^4 + 54c_1c^2 + 45c_1^2), \\ w_{1t} = [81\Phi_1^4 + 243c\Phi_1^3 + 243(c^2 + c_1)\Phi_1^2 + 18c(5c^2 + 27c_1)\Phi_1 + \mu]w_{1x}, \end{cases} \quad (59)$$

$$\begin{cases} w_{2y} = \Phi_2 w_{2x} + c_1, \\ w_{2t} = \left(81\Phi_2^4 + 243c_1\Phi_2^2 + \frac{405}{2}c_1^2\right)w_{2x} - c, \end{cases} \quad (60)$$

$$\begin{cases} w_{3y} = -\frac{1}{2}(\Phi_3 - 2c)w_{3x} + c_1, \\ w_{3t} = \left[\frac{81}{16}\Phi_3^4 - \frac{81(2c^2 + 3c_1)}{4}\Phi_3^2 + \frac{9c(8c^2 - 9c_1)}{2}\Phi_3 + \mu\right]w_{3x}, \end{cases} \quad (61)$$

$$\begin{cases} w_{4y} = \Phi_4 w_{4x} + c_1, \\ w_{4t} = 81\Phi_4^4 w_{4x} + 243c_1\Phi_4^2 w_{4x} + \frac{405c_1^2}{2}w_{4x}, \end{cases} \quad (62)$$

$$\begin{cases} w_{5y} = \frac{1}{4}(4\Phi_5^2 + 6c\Phi_5 + 3c^2)w_{5x}, \\ w_{5t} = \frac{81}{256}(256\Phi_6^8 + 1152c\Phi_5^7 + 2304c^2\Phi_5^6 + 2688c^3\Phi_5^5 + 2016c^4\Phi_5^4 + 1008c^5\Phi_5^3 \\ + 336c^6\Phi_5^2 + 72c^7\Phi_5 + 9c^8)w_{5x}, \end{cases} \quad (63)$$

$$\begin{cases} w_{6y} = cw_{6x} + c_1, \\ w_{6t} = -9(63w_{6x}^4 + 126w_{6x}^2 w_{6x3} + 63w_{6x} w_{6x2}^2 + 21w_{6x} w_{6x5} + 21w_{6x2} w_{6x4} + 21w_{6x3}^2 + w_{6x7})c \\ + 54(3w_{6x}^2 + w_{6x3})c^3 + 27(8w_{6x}^3 + 8w_{6x} w_{6x3} + 8c_1 w_{6x} + 7w_{6x2}^2 + w_{6x5})c^2 + 135(3w_{6x}^2 \\ + w_{6x3})c_1 c + \frac{9}{2}w_{6x}(2c^4 + 45c_1^2) - 27(15w_{6x}^3 + 15w_{6x} w_{6x3} + w_{6x5})c_1. \end{cases} \quad (64)$$

Similar to the decompositions of the fifth- and seventh-order PBKP equations, the first three decompositions (59)–(61) are the potential KdV decompositions, the fourth decomposition (62) is the special SS decomposition, the fifth decomposition (63) is the higher-order Burgers decomposition and the sixth decomposition (64) is the SK reduction. By observing the decomposition propositions 2, 3 and 4, we conclude that classical integrable systems solve the PBKP hierarchy after simple decompositions.

Conjecture 2. Each order equation in PBKP hierarchy (7) possesses six types of $\{y, t\}$ decomposed solutions.

The first three types of KdV decompositions can be obtained from the nonlinearization procedure of the Lax pairs of the PBKP hierarchy, the fourth decomposition is the special SS decomposition and the fifth type of the STO decomposition can be concluded from hints of the Bäcklund transformations of the bilinear PKP equation. The sixth type of SK reduction is a direct conclusion because the PBKP hierarchy is just a (2+1)-dimensional extension of the (1+1)-dimensional SK hierarchy.

4 | SPECIAL LINEAR SUPERPOSITIONS OF THE PBKP HIERARCHY

It is well known that for the physical systems frequently characterized by nonlinear differential equations, there is no linear superposition theorem. However, there exist some types of nonlinear superposition properties if the solutions are linked by

some special requirements such as the Bäcklund/Darboux transformations^{36,37,38,39,40,41,42}. For the BKP equation (alias (2+1)-dimensional Sawada-Kotera equation), the nonlinear superpositions have been studied by Li and Hu⁴³ in bilinear forms. In this section, we investigate the possible linear superpositions of the PBKP hierarchy for the possible decomposition solutions obtained in the last section.

4.1 | Special linear superpositions of the fifth-order PBKP equation (2)

For the fifth-order PBKP equation (2), if w_i and w_j are the decompositions of the PBKP equation listed in the last section, then requiring the linear combination $a_1 w_i + a_2 w_j$ is also a solution of the PBKP equation (2) yields the following special linear superposition theorem.

Proposition 5. Suppose that w_1, w_2, w_3, w_4, w_5 and w_6 are solutions of the PBKP equation (2) with the conditions

$$\begin{aligned} w_{1y} &= (\Phi_1 + c)w_{1x}, \quad w_{1t} = (9\Phi_1^2 + 15c\Phi_1 + 5c^2)w_{1x}, \quad \Phi_1 \equiv \partial_x^2 + 4w_{1x} - 2\partial_x^{-1}w_{1x2}, \\ w_{2y} &= (\Phi_2 - c)w_{2x}, \quad w_{2t} = (9\Phi_2^2 - 15c\Phi_2 + 5c^2)w_{2x}, \quad \Phi_2 \equiv \partial_x^2 + 4w_{2x} - 2\partial_x^{-1}w_{2x2}, \\ w_{3y} &= \Phi_3 w_{3x}, \quad w_{3t} = 9\Phi_3^2 w_{3x}, \quad \Phi_3 \equiv \partial_x^2 + 4w_{3x} - 2\partial_x^{-1}w_{3x2}, \\ w_{4y} &= \Phi_4 w_{4x}, \quad w_{4t} = 9\Phi_4^2 w_{4x}, \quad \Phi_4 \equiv \partial_x^2 + 2w_{4x} - \partial_x^{-1}w_{4x2}, \\ w_{5y} &= \Phi_5 w_{5x} - \frac{c^2}{6}, \quad w_{5t} = (9\Phi_5^2 - \frac{5c^2}{2})w_{5x}, \quad \Phi_5 \equiv \partial_x^2 + 2w_{5x} - \partial_x^{-1}w_{5x2}, \\ w_{6y} &= \Phi_6 w_{6x} - \frac{c^2}{6}, \quad w_{6t} = (9\Phi_6^2 - \frac{5c^2}{2})w_{6x}, \quad \Phi_6 \equiv \partial_x^2 + 2w_{6x} - \partial_x^{-1}w_{6x2}, \end{aligned} \quad (65)$$

then the linear superpositions

$$w_7 = w_1 + w_2, \quad w_8 = w_3 + \frac{1}{2}w_4, \quad w_9 = \frac{1}{2}(w_5 + w_6) \quad (66)$$

are at the same time the solutions of the PBKP equation (2).

Proof. By substituting the linear superposition relations (66) with the decomposition results (65) into the PBKP equation (2), one can directly prove the proposition. \square

In fact, one can directly substitute the general superposition assumption $w = f(w_i, w_j)$ with w_i and w_j being any one of the decompositions given in the proposition 2 to prove that the results (66) are the only possible superpositions if there is no further relations among w_i and w_j .

It is noted that though there are six possible decompositions for the fifth-order PBKP equation (2) as shown in the proposition 2, only the first two types of decomposition solutions (22) and (23) can be linearly combined to construct new solutions of the PBKP equation (2). The third type of KdV decomposition (24), the special SS decomposition (25), the STO decomposition (26) and the third SK reduction (27) can not be used to find new solutions via superposition assumption $w = f(w_i, w_j)$ if there is no further relations among seed solutions w_i and w_j . It should be emphasized that w_i and w_j , $i, j = 1, 2, \dots, 6$, appeared in superpositions (66) are independent solutions of the decompositions. There is no further relations. Here, we list some special linear superposition structures of the fifth BKP equation.

Example 1. $n + m$ solitons.

$$w_1 = 2(\ln f_n^+)_x, \quad w_2 = 2(\ln f_m^-)_x, \quad w_3 = w_1|_{c=0}, \quad w_4 = w_6|_{c=0}, \quad w_5 = 4(\ln g_n^+)_x + \frac{cx}{3}, \quad w_6 = w_5|_{g_n^+ \rightarrow g_m^-}, \quad (67)$$

$$f_n^\pm = \sum_{\mu=0,1} \exp\left(\sum_{i=1}^n \mu_i \xi_i^\pm + \sum_{1 \leq i < j \leq n} \mu_i \mu_j \theta_{ij}^\pm\right), \quad g_n^\pm = \sum_{\mu=0,1} \exp\left(\sum_{i=1}^n \mu_i \eta_i^\pm + \sum_{1 \leq i < j \leq n} \mu_i \mu_j \theta_{ij}^\pm\right), \quad (68)$$

where the summation of μ should be done for all permutations of $\mu_i = 0, 1$, $i = 1, 2, \dots, n$,

$$\xi_{i\pm} = k_{i\pm}x + k_{i\pm}(k_{i\pm} \pm c)y + (9k_{i\pm}^5 \pm 15ck_{i\pm}^3 + 5c^2k_{i\pm})t + \xi_{i0}^\pm, \quad (69)$$

$$\eta_{i\pm} = k_{i\pm}x + k_{i\pm}(k_{i\pm} + c)y + (9k_{i\pm}^5 + 15ck_{i\pm}^3 + 5c^2k_{i\pm})t + \eta_{i0}^\pm, \quad \exp(\theta_{ij}^\pm) = \frac{(k_{i\pm} - k_{j\pm})^2}{(k_{i\pm} + k_{j\pm})^2},$$

and $k_{i\pm}$, c , ξ_{i0}^\pm and η_{i0}^\pm are arbitrary constants. All the related solutions of the BKP equation (1) with $u = w_x$ and $w = w_7, w_8, w_9$ expressed in the proposition 5 with (67), (68) and (69) denote some types of $n + m$ solitons solutions. To illustrate the linear superposition structure more clearly, let us look at some figures. The following figures are intended to illuminate two-soliton,

three-soliton and their linear superposition of the BKP equation (1) respectively with condition $\{c = 2, k_1 = -1, k_2 = 1.5, k_3 = -2, \xi_{10} = 2, \xi_{20} = 10, \xi_{30} = 10\}$ at time $t = 0$.

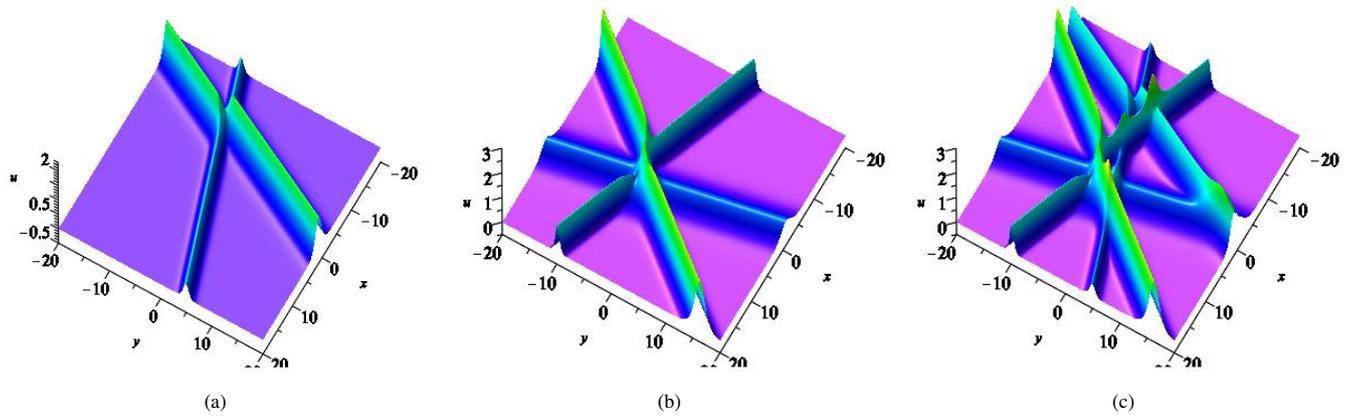


FIGURE 1 Two-soliton (a), three-soliton (b) and their linear superposition (c) of BKP equation located at $t = 0$, respectively.

Example 2. n solitons with periodic cnoidal background waves. If one of w_1 and w_2 in example 1 is replaced by a periodic solution, say,

$$w_{2x} = 2k^2 m^2 \text{cn}^2(kx + [4k^3(2m^2 - 1) - ck]y + [5c^2k - 60k^3(2m^2 - 1)c + 72k^5(7m^4 - 7m^2 + 2)]t + \xi_0, m), \quad (70)$$

where $\text{cn}(\xi, m)$ is a Jacobi cn function with the modula m , then $w_7 = w_1 + w_2$ becomes an n -soliton solution on a periodic cnoidal background wave. Figure 2 is drawn for the case of three-soliton moving on a periodic cnoidal wave by taking simply $\{c = 1.2, k_1 = -1, k_2 = 1.5, k_3 = -2, \xi_0 = \xi_{10} = 2, \xi_{20} = \xi_{30} = 10, m = 1.2, k = 0.5\}$ at $t = 0$.

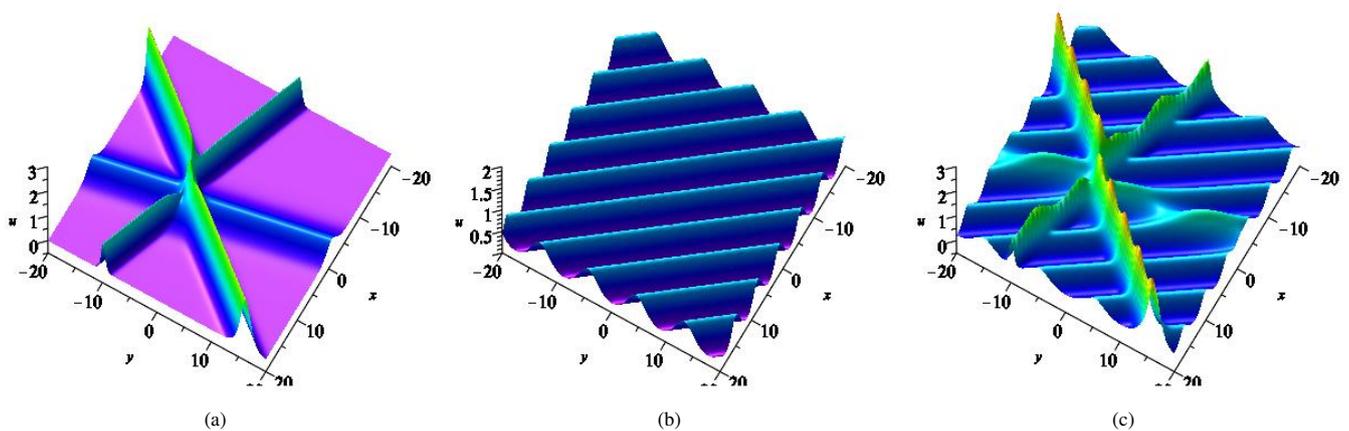


FIGURE 2 Structures of three-soliton (a), periodic cnoidal wave (b) and their linear superposition (c) of BKP equation at $t = 0$, respectively.

Example 3. n solitons solution combined with soliton-cnoidal wave interaction solution. If w_3 is given by (67) and w_4 is fixed as

$$w_{4x} = \frac{8}{3}s_x^2 - \frac{4s_{x3}}{3s_x} + \frac{s_{x2}^2}{3s_x^2} + 4s_{x2} \tanh(s) - 4s_x^2 \tanh^2(s), \quad (71)$$

$$s = mkx + l_0y + \omega_0t + \xi_0 - \operatorname{arctanh}(msn(kx + ly + \omega t + \xi_1, m^2)) \quad (72)$$

with arbitrary constants $k, m, \xi_0, \xi_1, \lambda$ and $l_0 = mk\lambda - \frac{mk^3}{2}(1 + 3m^2)(m^2 + 3)$, $l = k\lambda + \frac{k^3}{2}[(m^2 - 1)^2 - 16m^2]$, $\omega = 15l\lambda - \frac{15}{2}k\lambda^2 + \frac{9k^5}{8}(3m^8 + 20m^6 + 722m^4 + 20m^2 + 3)$ and $\omega_0 = 15l_0\lambda - \frac{15}{2}mk\lambda^2 - \frac{9mk^5}{8}(5m^8 - 180m^6 - 428m^4 - 180m^2 + 5)$, then $w = w_3 + \frac{1}{2}w_4$ expresses an n -soliton solution combined with a soliton-cnoidal wave interaction solution. In figure 3, we draw the behavior of a two-soliton, a soliton-cnoidal wave and their linear combination of BKP equation (1) at $y = 0$ with arbitrary parameters fixed as $\{c = 0, k_1 = 0.9, k_2 = 0.65, \xi_{10} = 12, \xi_{20} = 0, k = 0.3, m = 0.8, \xi_0 = -10, \xi_1 = 10, \lambda = 0.3\}$.

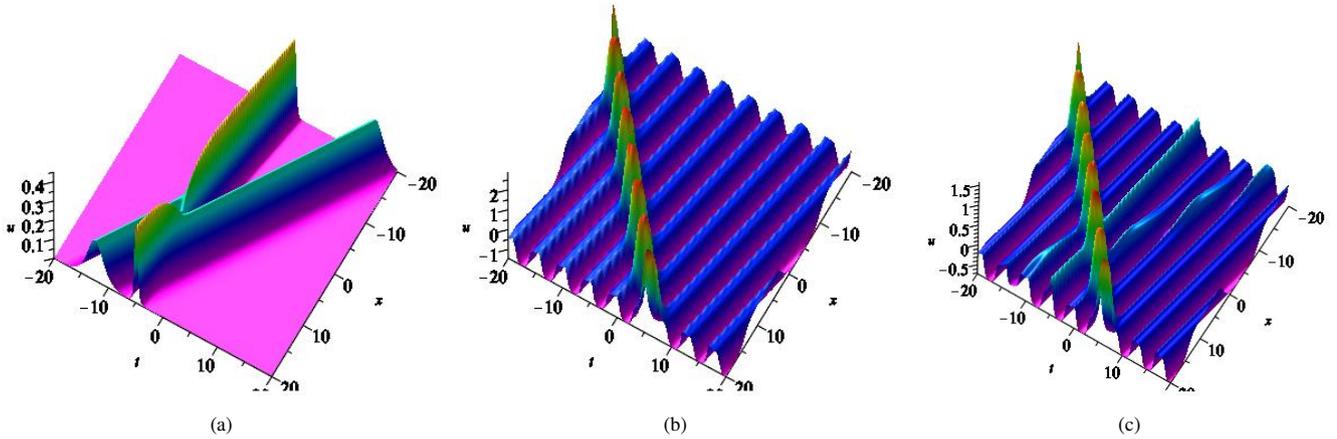


FIGURE 3 Two-soliton (a), soliton-cnoidal interaction wave (b) and their linear combination (c) of BKP equation at $y = 0$, respectively.

Example 4. Combination solutions of two different modified Schwarzian KdV waves. If the solutions of (65) for w_5 and w_6 are rewritten in the forms

$$w_{5x} = 2s_{1x}^2 - \frac{s_{1x2}^2}{s_{1x}^2} - S_1 + \frac{\lambda_1}{3} + 4s_{1x2} \tanh(s_1) - 4s_{1x}^2 \tanh^2(s_1), \quad (73)$$

$$w_{6x} = 2s_{2x}^2 - \frac{s_{2x2}^2}{s_{2x}^2} - S_2 + \frac{\lambda_2}{3} + 4s_{2x2} \tanh(s_2) - 4s_{2x}^2 \tanh^2(s_2), \quad (74)$$

where s_1 and s_2 are any solutions of the modified Schwarzian KdV systems

$$Y_i = S_i - 2s_{ix}^2 + \lambda_i, \quad Y_i \equiv \frac{s_{iy}}{s_{ix}}, \quad T_i \equiv \frac{s_{it}}{s_{ix}}, \quad S_i \equiv \frac{s_{ix3}}{s_{ix}} - \frac{3s_{ix2}^2}{2s_{ix}^2}, \quad i = 1, 2, \quad (75)$$

$$T_i = 9S_{ix2} + \frac{27}{2}S_i^2 + 15\lambda_i S_i + \frac{5c^2}{2} + \frac{15\lambda_i^2}{2} + 6s_{ix}^2(9s_{ix}^2 - 15S_i - 5\lambda_i) - 90s_{ixx}^2 \quad (76)$$

with arbitrary constants λ_i , then $w = \frac{1}{2}(w_5 + w_6)$ is a combination solution of two modified Schwarzian KdV waves. The modified Schwarzian KdV systems (75) and (76) are related to the Schwarzian KdV systems (the systems of (75) and (76) contain only the Möbius transformation invariant terms related to the Möbius invariants $T_i, Y_i, S_i, S_{ix2}, \lambda_i$ and c) by the transformation $s_i \rightarrow \operatorname{arctanh}(s_i^{-1})$. The modified Schwarzian KdV systems (75) and (76) possess various interaction solutions including the

following soliton-cnoidal interaction solutions^{44,45}

$$s_1 = m_1 k_1 x + p_1 y + \omega_1 t + \xi_{10} - \operatorname{arctanh}[m_1 \operatorname{sn}(k_1 x + q_1 y + \Omega_1 t + \eta_{10}, m_1^2)], \quad (77)$$

$$s_2 = k_0(x + p_2 y + \omega_2 t) + \xi_{20} - \operatorname{arctanh}[m_2 \operatorname{sn}(k_2 x + q_2 y + \Omega_2 t + \eta_{20}, m_2^2)] \quad (78)$$

with arbitrary constants $k_i, m_i, \lambda_i, \xi_{i0}, \eta_{i0}, i = 1, 2$, and the other constants $\{k_0, p_1, q_i, \omega_i, \Omega_i\}$ being given by

$$p_1 = m_1 k_1 \lambda_1 - \frac{m_1 k_1^3}{2}(m_1^2 + 3)(3m_1^2 + 1),$$

$$q_1 = k_1 \lambda_1 + \frac{k_1^3}{2}[(m_1^2 - 1)^2 - 16m_1^2],$$

$$\omega_1 = 15p_1 \lambda_1 + \frac{5}{2}m_1 k_1 (c^2 - 3\lambda_1^2) - \frac{9}{8}m_1 k_1^5 [5m_1^8 - 180m_1^2(1 + m_1^4) - 418m_1^4],$$

$$\Omega_1 = 15q_1 \lambda_1 + \frac{5k_1}{2}(c^2 - 3\lambda_1^2) + \frac{9k_1^5}{8}(3m_1^8 + 20m_1^6 + 722m_1^4 + 20m_1^2 + 3),$$

$$k_0 = -\frac{k_2}{2}(m_2^2 + 1), \quad p_2 = \lambda_2 - \frac{k_2^2}{2}(1 + 14m_2^2 + m_2^4),$$

$$q_2 = k_2 \lambda_2 - \frac{k_2^3}{2}(5 + 6m_2^2 + 5m_2^4),$$

$$\omega_2 = 15p_2 \lambda_2 + \frac{5}{2}(c^2 - 2\lambda_2^2) + \frac{9k_2^4}{8}(3m_2^8 + 212m_2^2(1 + m_2^4) + 3),$$

$$\Omega_2 = 15q_2 \lambda_2 - \frac{5}{2}k_2(c^2 - 3\lambda_2^2) + \frac{9k_2^5}{8}(43(1 + m_2^8) + 180m_2^2(1 + m_2^4) + 322m_2^4).$$

Figure 4 shows a linear superposition structure of two soliton-cnoidal interaction waves for the field u with the conditions $k_1 = 0.3, k_2 = 0.25, m_1 = 1.2, m_2 = 0.6, \xi_{10} = \xi_{20} = 0, \eta_{10} = \eta_{20} = 10, \lambda_1 = \lambda_2 = 0$ at $y = 0$. Obviously, figures (c)

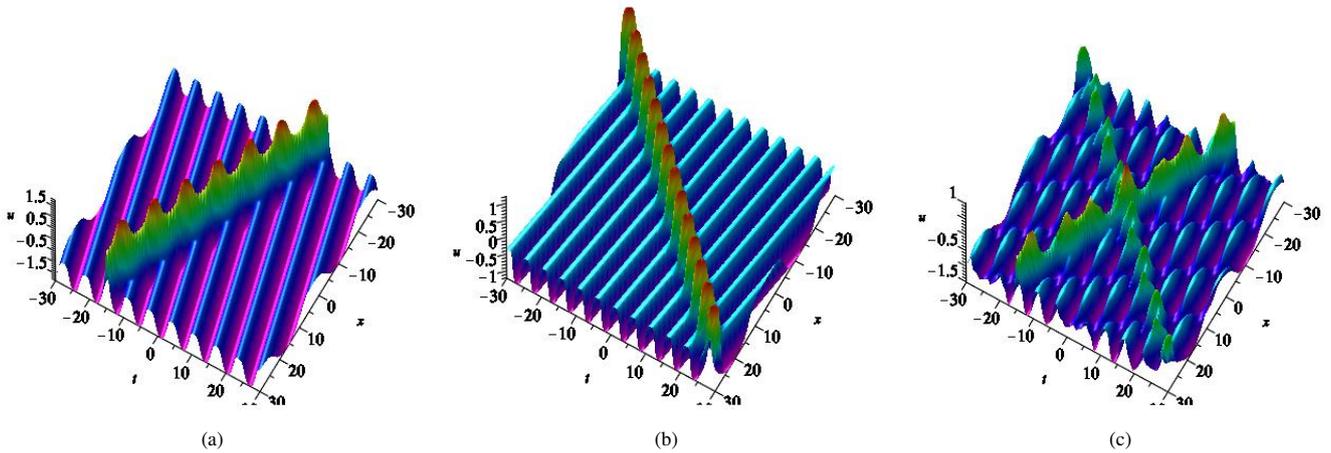


FIGURE 4 Two soliton-cnoidal interaction waves (a) and (b), their linear superposition (c) of BKP equation at $y = 0$.

in Figures 1 -4 clearly display the linear superpositions of (a) and (b), correspondingly. Interestingly, two waves can pass through each other without altering each other. Sometimes, this leads to some truly weird behaviors. It also means that waves can constructively or destructively interfere under the effect of linear superposition.

4.2 | Special linear superpositions of the seventh-order PBKP equation (11)

Similar to the fifth-order PBKP equation, though there are six decomposition solutions (53)–(58), only the first two decompositions (53) and (54) can be used to construct new solutions of the seventh-order PBKP equation (11) via linear superpositions. The result is summarized in the following proposition.

Proposition 6. If w_1, w_2, w_3, w_4, w_5 and w_6 are solutions of the seventh-order PBKP equation (11) with conditions

$$\begin{aligned}
 w_{1y} &= (\Phi_1 + c)w_{1x}, \quad w_{1t} = (27\Phi_1^3 + 63c\Phi_1^2 + 42c^2\Phi_1 + 7c^3)w_{1x} - c, \\
 w_{2y} &= (\Phi_2 - c)w_{2x}, \quad w_{2t} = (27\Phi_2^3 - 63c\Phi_2^2 + 42c^2\Phi_2 - 7c^3)w_{2x} - c_1, \\
 w_{3y} &= \Phi_3 w_{3x}, \quad w_{3t} = 27\Phi_3^3 w_{3x} - c, \\
 w_{4y} &= \Phi_4 w_{4x}, \quad w_{4t} = 27\Phi_4^3 w_{4x} - c_1, \\
 w_{5y} &= \Phi_5 w_{5x} + c_1, \quad w_{5t} = 9(3\Phi_5^3 + 7c_1\Phi_5)w_{5x} - c, \\
 w_{6y} &= \Phi_6 w_{6x} + c_1, \quad w_{6t} = 9(3\Phi_6^3 + 7c_1\Phi_6)w_{6x} - c_1,
 \end{aligned} \tag{79}$$

then the special combinations corresponding to

$$w_7 = w_1 + w_2, \quad w_8 = w_3 + \frac{1}{2}w_4, \quad w_9 = \frac{1}{2}(w_5 + w_6) \tag{80}$$

are also solutions of the seventh-order PBKP equation (11).

Substituting (80) with (79) into (11), one can directly prove proposition 6. Therefore, in the seventh PBKP case, the principle of linear superposition (80) holds for the special types of decomposition solutions given by (79).

4.3 | Special linear superpositions of the ninth-order PBKP equation (12)

For the ninth-order PBKP equation (12), we have an analogous result for the linear superpositions of decomposition solutions.

Proposition 7. Let w_1, w_2, w_3, w_4, w_5 and w_6 be solutions of the ninth-order PBKP equation (12) with the decompositions

$$\begin{aligned}
 w_{1y} &= (\Phi_1 + c)w_{1x}, \quad w_{1t} = (81\Phi_1^4 + 243c\Phi_1^3 + 243c^2\Phi_1^2 + 90c^3\Phi_1 + 9c^4)w_{1x} - c, \\
 w_{2y} &= (\Phi_2 - c)w_{2x}, \quad w_{2t} = (81\Phi_2^4 - 243c\Phi_2^3 + 243c^2\Phi_2^2 - 90c^3\Phi_2 + 9c^4)w_{2x} - c_1, \\
 w_{3y} &= \Phi_3 w_{3x}, \quad w_{3t} = 81\Phi_3^4 w_{3x} - c, \\
 w_{4y} &= \Phi_4 w_{4x}, \quad w_{4t} = 81\Phi_4^4 w_{4x} - c_1, \\
 w_{5y} &= \Phi_5 w_{5x} + c_1, \quad w_{5t} = \left(81\Phi_5^4 + 243c_1\Phi_5^2 + \frac{405}{2}c_1^2\right)w_{5x} - c, \\
 w_{6y} &= \Phi_6 w_{6x} + c_1, \quad w_{6t} = \left(81\Phi_6^4 + 243c_1\Phi_6^2 + \frac{405}{2}c_1^2\right)w_{6x} - c_1,
 \end{aligned} \tag{81}$$

then $w_7 = w_1 + w_2$, $w_8 = w_3 + \frac{1}{2}w_4$ and $w_9 = \frac{1}{2}(w_5 + w_6)$ are still solutions of the ninth-order PBKP equation (12).

From the detailed examples of the fifth-, seventh-, and ninth-order PBKP equations, similar linear superposition conjecture can be found for the whole PBKP hierarchy (7).

Conjecture 3. For every equation of the PBKP hierarchy, there exist three possible special types of linear superposition solutions.

It should be mentioned that though the solutions w_i , $i = 1, \dots, 6$ shown in the propositions 5, 6 and 7 are the potential KdV decomposition solutions, their special linear superposition solutions w_7, w_8 and w_9 are not the decomposition solutions.

In the next section, we study the possible decompositions and linear superpositions for the dispersionless PBKP hierarchy.

5 | DECOMPOSITIONS AND SPECIAL LINEAR SUPERPOSITIONS OF THE DISPERSIONLESS PBKP HIERARCHY

5.1 | Mastersymmetries and dPBKP hierarchy

The dPBKP equation possesses the form

$$w_{xt} = 15(w_x w_y - w_x^3)_x + 5w_{y2} = \partial_x (15w_x w_y - 15w_x^3 + 5\partial_x^{-1} w_{yy}) \equiv \partial_x V. \quad (82)$$

Using the mastersymmetry approach similar to that described in conjecture 1 for the PBKP equation, one can make the following conjecture for the dPBKP equation (82).

Conjecture 4. y^n for arbitrary positive integers are mastersymmetries with $V_{[1]}^{n+1} y^n = 0$.

If the conjecture is correct, then we can write the dPBKP hierarchy in the form

$$w_{t_{2n-1}} = V_{2n-1} = \frac{1}{3n!5^n} V_{[1]}^n y^n. \quad (83)$$

We have check that the conjecture 4 is correct for $n = 1, 2, \dots, 6$. The corresponding equations of the dPBKP hierarchy (83) possess the forms

$$w_{t_1} = \frac{1}{15} V_{[1]} y = w_x, \quad (84)$$

$$w_{t_3} = \frac{1}{150} V_{[1]}^2 y^2 = 3w_y, \quad (85)$$

$$w_{t_5} = \frac{1}{2250} V_{[1]}^3 y^3 = V_5 = V, \quad (86)$$

$$w_{t_7} = \frac{1}{45000} V_{[1]}^4 y^4 = V_7 \quad (87)$$

$$= 21\partial_x^{-1}(w_x w_{y2}) + 7\partial_x^{-2} w_{y3} - 63w_x^2 w_y + 126\partial_x^{-1}(w_x w_{x2} w_y) - 63w_x^4 + 21w_x \partial_x^{-1}(w_{y2}) + \frac{63}{2} w_y^2,$$

$$w_{t_9} = \frac{1}{1125000} V_{[1]}^5 y^5 = V_9 \quad (88)$$

$$= \partial_x^{-1} \left\{ \frac{27}{2} w_{x2} [15w_y^2 - 108w_x^2 w_y + 36\partial_x^{-1}(w_{x2} w_x w_y) + 6\partial_x^{-1}(w_x w_{y2}) + 2\partial_x^{-2} w_{y3}] \right. \\ \left. - 81w_{xy} (12w_x^3 - 5w_x w_y - \partial_x^{-1} w_{y2}) + 135w_y w_{y2} + 54w_x \partial_x^{-1} w_{y3} + 27\partial_x^{-1}(w_x w_{y2})_y \right. \\ \left. + 162\partial_x^{-1}(w_{x2} w_x w_{y2} - w_x w_{xy}^2) + 9\partial_x^{-2} w_{y4} \right\},$$

$$w_{t_{11}} = \frac{1}{33750000} V_{[1]}^6 y^6 = V_{11} \quad (89)$$

$$= \partial_x^{-2} \left\{ w_{x3} [7128w_x^5 + 33\partial_x^{-3} w_{y4} - 594\partial_x^{-2}(w_x w_{xy}^2) + 99\partial_x^{-2}(w_{xy} w_{y2}) + 99\partial_x^{-2}(w_x w_{y3}) \right. \\ \left. + 594\partial_x^{-1}(w_x w_{x2} w_{y2}) + 198\partial_x^{-1}(w_y w_{y2}) - 99\partial_x^{-1}(w_{y3} w) + 3564\partial_x^{-1}(w_x^2 w_{x2} w_y) - 297\partial_x^{-1}(w_x^2 w_{y2}) \right. \\ \left. - 9504w_y w_x^3 - 1485w_x^2 \partial_x^{-1} w_{y2} + 495w_y \partial_x^{-1} w_{y2} + 99w \partial_x^{-1} w_{y3}] + w_{yx2} [2079/2w_y^2 + 99\partial_x^{-2} w_{y3} \right. \\ \left. + 1782\partial_x^{-1}(w_x w_{x2} w_y) + 297\partial_x^{-1}(w_x w_{y2}) - 3267w_x^4 - 2970w_x^2 w_y + 495w_x \partial_x^{-1} w_{y2}] \right. \\ \left. + w_{xy2} (1386w_x w_y - 1980w_x^3 + 165\partial_x^{-1} w_{y2}) + (198w_x^2 + 231w_y) w_{y3} + w_{x2}^2 (35640w_x^4 - 24948w_x^2 w_y \right. \\ \left. - 2970w_x \partial_x^{-1} w_{y2}) + w_{x2} [w_{xy} (990\partial_x^{-1} w_{y2} - 22572w_x^3 - 4158w_x w_y) + w_{y2} (2079w_y - 5346w_x^2) \right. \\ \left. + 297w_x \partial_x^{-1} w_{y3} + 297\partial_x^{-1}(w_{xy} w_{y2}) + 297\partial_x^{-1}(w_x w_{y3}) - 1782\partial_x^{-1}(w_x w_{xy}^2) + 1782\partial_x^{-1}(w_x w_{x2} w_{y2}) \right. \\ \left. + 99\partial_x^{-2} w_{y4}] + w_{xy}^2 (-5346w_x^2 + 2079w_y) + w_{xy} (1782w_x w_{y2} + 297\partial_x^{-1} w_{y3}) + 99w_x \partial_x^{-1} w_{y4} \right. \\ \left. + 33\partial_x^{-1} w_x w_{y4} + 11\partial_x^{-2} w_{y5} + \frac{495}{2} w_y^2 - 198\partial_x^{-1} w_{xy}^3 + 594\partial_x^{-1}(w_{y2} w_{x2} w_{xy}) + 594\partial_x^{-1}(w_{y2} w_x w_{yx2}) \right. \\ \left. + 198\partial_x^{-1}(w_x w_{x2} w_{y3}) + 66\partial_x^{-1}(w_{xy} w_{y3}) \right\}.$$

5.2 | Decompositions of dPBKP hierarchy

Proposition 8. If w_1 , w_2 , w_3 and w_4 are solutions of the following dispersionless KdV systems,

$$\begin{cases} w_{1y} = a_2 w_{1x}^2 + a_1 w_{1x} + c_1, \\ w_{1t_5} = \frac{5}{3}(a_2 + 3)(4a_2 - 3)w_{1x}^3 + 5a_1(2a_2 + 3)w_{1x}^2 + 5(a_1^2 + 3c_1)w_{1x}, \end{cases} \quad (90)$$

$$\begin{cases} w_{2y} = a_2 w_{2x}^2 + a_1 w_{2x} + c_1, \\ w_{2t_7} = \frac{7}{2}(a_2 - 1)(4a_2 + 3)(a_2 + 6)w_{2x}^4 + 7a_1(4a_2^2 + 19a_2 - 3)w_{2x}^3 + bw_{2x}^2 + 7a_1(a_1^2 + 9c_1)w_{2x}, \end{cases} \quad (91)$$

$$\begin{cases} w_{3y} = a_2 w_{3x}^2 + a_1 w_{3x} + c_1, \\ w_{3t_9} = \frac{9}{10}a_2(a_2 + 9)(8a_2 - 9)(4a_2 + 9)w_{3x}^5 + c_4 w_{3x}^4 + c_3 w_{3x}^3 + c_2 w_{3x}^2 + \frac{9}{2}c_1 w_{3x}, \end{cases} \quad (92)$$

$$\begin{cases} w_{4y} = a_2 w_{4x}^2 + a_1 w_{4x} + c_1, \\ w_{4t_{11}} = \frac{11}{30}(d_6 w_{4x}^6 + 3d_5 w_{4x}^5 + 15d_4 w_{4x}^4 + 10d_3 w_{4x}^3 + 15d_2 w_{4x}^2 + 15d_1 w_{4x}), \end{cases} \quad (93)$$

with $b = 21(a_1^2 a_2 + 3c_1 a_2 + 3a_1^2)$, $c_4 = 9a_1(2a_2 + 3)(4a_2^2 + 30a_2 - 21)$, $c_3 = 9(8a_1^2 a_2^2 + 20c_1 a_2^2 + 58a_1^2 a_2 + 45c_1 a_2 + 24a_1^2 - 45c_1)$, $c_2 = 9a_1(4a_1^2 a_2 + 39c_1 a_2 + 18a_1^2 + 45c_1)$, $c_1 = 2a_1^4 + 54c_1 a_1^2 + 45c_1^2$, $d_6 = (a_2 + 12)(5a_2 - 6)(4a_2 + 15)(8a_2 - 3)(a_2 + 1)$, $d_5 = a_1(160a_2^4 + 2236a_2^3 + 4551a_2^2 - 3357a_2 - 2160)$, $d_4 = (40a_1^2 a_2^3 + 84c_1 a_2^3 + 490a_1^2 a_2^2 + 483c_1 a_2^2 + 780a_1^2 a_2 - 189c_1 a_2 - 270a_1^2 - 378c_1)$, $d_3 = a_1(40a_1^2 a_2^2 + 252c_1 a_2^2 + 390a_1^2 a_2 + 1197c_1 a_2 + 360a_1^2 - 189c_1)$, $d_2 = 10a_1^4 a_2 + 144c_1 a_1^2 a_2 + 60a_1^4 + 189c_1^2 a_2 + 378c_1 a_1^2$ and $d_1 = a_1(2a_1^4 + 66c_1 a_1^2 + 189c_1^2)$, then w_1 , w_2 , w_3 and w_4 are solutions of the fifth-, seventh-, ninth- and eleventh-order dPBKP equations (86), (87), (88) and (89), respectively.

In fact, for the dPBKP hierarchy there are more general decomposition solutions than those shown in proposition 8. Here is a general decomposition conjecture:

Conjecture 5. The dPBKP hierarchy (83) possesses a $\{y, t_{2n-1}\}$ decomposed solution w_n for the n th-order equation

$$\begin{cases} w_{ny} = F(w_{nx}), \\ w_{nt_{2n-1}} = G(w_{nx}) \end{cases} \quad (94)$$

with $F(w_{nx})$ being an arbitrary function of w_{nx} and $G(w_{nx})$ being a function determined by $F(w_{nx})$. Especially, if

$$F(w_{nx}) = \sum_{k=0}^m a_k w_{nx}^k, \quad (95)$$

then

$$G(w_{nx}) = \sum_{k=0}^{(m-1)(n-1)+1} b_k w_{nx}^k \quad (96)$$

with b_k being determined by a_k .

For the fifth-order dPBKP equation (86), $n = 3$, it is easy to find that

$$G(x_3) = -15x_3^3 + 15x_3 F + 5\partial_{x_3}^{-1} F_{x_3}^2, \quad F = F(x_3), \quad x_3 = w_{3x}.$$

For the seventh-order dPBKP equation (87), $n = 4$, we have

$$G(x_4) = -\frac{21}{2}[6x_4^4 - 3F^2 + \partial_{x_4}^{-1}(6x_4^2 F_{x_4} - 4x_4 F_{x_4}^2)], \quad F = F(x_4), \quad x_4 = w_{4x}.$$

Proposition 8 is just for the special cases of conjecture 5 related to $m = 2$ of (95) and (96).

5.3 | Linear superpositions of dPBKP hierarchy

Similar to the PBKP hierarchy (7), for some special decomposition solutions, the dPBKP hierarchy also admits solutions in terms of their sum. For the fifth-, seventh-, ninth- and eleventh-order dPBKP equations (86), (87), (88) and (89), some special linear superposition theorems are summarised in the following propositions 9, 10, 11 and 12.

Proposition 9. Assume w_1, w_2, w_3, w_4, w_5 and w_6 are solutions of the fifth-order dPBKP equation (86) with the decomposition conditions

$$w_{1y} = 3w_{1x}^2 + cw_{1x}, w_{1t_5} = 90w_{1x}^3 + 45cw_{1x}^2 + 5c^2w_{1x}, \quad (97)$$

$$w_{2y} = 3w_{2x}^2 - cw_{2x}, w_{2t_5} = 90w_{2x}^3 - 45cw_{2x}^2 + 5c^2w_{2x}, \quad (98)$$

$$w_{3y} = 3w_{3x}^2, w_{3t_5} = 90w_{3x}^3, \quad (99)$$

$$w_{4y} = \frac{3}{2}w_{4x}^2, w_{4t_5} = \frac{45}{2}w_{4x}^3, \quad (100)$$

$$w_{5y} = \frac{3}{2}w_{5x}^2 + c_1, w_{5t_5} = \frac{45}{2}w_{5x}^3 + 15c_1w_{5x}, \quad (101)$$

$$w_{6y} = \frac{3}{2}w_{6x}^2 + c_1, w_{6t_5} = \frac{45}{2}w_{6x}^3 + 15c_1w_{6x}, \quad (102)$$

then w_7, w_8 and w_9 with the linear superposition properties

$$w_7 = w_1 + w_2, w_8 = w_3 + \frac{1}{2}w_4, w_9 = \frac{1}{2}(w_5 + w_6) \quad (103)$$

are also solutions of the fifth-order dPBKP equation (86).

Proposition 10. If the solutions $w_i, i = 1, 2, \dots, 12$, of the seventh-order dPBKP equation (87) are satisfied by the decomposition properties

$$w_{1y} = 3w_{1x}^2 + cw_{1x}, w_{1t_7} = 7(135w_{1x}^4 + 90cw_{1x}^3 + 18c^2w_{1x}^2 + c^3w_{1x}), \quad (104)$$

$$w_{2y} = 3w_{2x}^2 - cw_{2x}, w_{2t_7} = 7(135w_{2x}^4 - 90cw_{2x}^3 + 18c^2w_{2x}^2 - c^3w_{2x}), \quad (105)$$

$$w_{3y} = 3w_{3x}^2, w_{3t_7} = 945w_{3x}^4, \quad (106)$$

$$w_{4y} = \frac{3}{2}w_{4x}^2, w_{4t_7} = \frac{945}{8}w_{4x}^4, \quad (107)$$

$$w_{5y} = \frac{3}{2}w_{5x}^2 + c_1, w_{5t_7} = \frac{945}{8}w_{5x}^4 + \frac{189}{2}c_1w_{5x}^2, \quad (108)$$

$$w_{6y} = \frac{3}{2}w_{6x}^2 + c_1, w_{6t_7} = \frac{945}{8}w_{6x}^4 + \frac{189}{2}c_1w_{6x}^2, \quad (109)$$

$$w_{7y} = 3w_{7x}^2, w_{7t_7} = 945w_{7x}^4, \quad (110)$$

$$w_{8y} = -\frac{2}{3}w_{8x}^2, w_{8t_7} = -\frac{280}{27}w_{8x}^4, \quad (111)$$

$$w_{9y} = -\frac{2}{3}w_{9x}^2, w_{9t_7} = -\frac{280}{27}w_{9x}^4, \quad (112)$$

$$w_{10y} = -\frac{2}{3}w_{10x}^2, w_{10t_7} = -\frac{280}{27}w_{10x}^4, \quad (113)$$

$$w_{11y} = -\frac{2}{3}w_{11x}^2, w_{11t_7} = -\frac{280}{27}w_{11x}^4, \quad (114)$$

$$w_{12y} = \frac{3}{2}w_{12x}^2, w_{12t_7} = \frac{945}{8}w_{12x}^4, \quad (115)$$

then $w_i, i = 13, \dots, 18$, with the linear superposition behaviors

$$\begin{aligned} w_{13} &= w_1 + w_2, w_{14} = w_3 + \frac{1}{2}w_4, w_{15} = \frac{1}{2}(w_5 + w_6), \\ w_{16} &= w_7 - \frac{2}{9}w_8, w_{17} = -\frac{2}{9}(w_9 + w_{10}), w_{18} = -\frac{2}{9}w_{11} + \frac{1}{2}w_{12} \end{aligned} \quad (116)$$

are also solutions of the seventh-order dPBKP equation (87).

Proposition 11. The solutions w_i , $i = 1, 2, \dots, 12$, with the decomposition relations

$$w_{1y} = 3w_{1x}^2 + cw_{1x}, w_{1t_9} = 9w_{1x}(6w_{1x} + c)(3w_{1x} + c)(63w_{1x}^2 + 21cw_{1x} + c^2), \quad (117)$$

$$w_{2y} = 3w_{2x}^2 - cw_{2x}, w_{2t_9} = 9w_{2x}(6w_{2x} - c)(3w_{2x} - c)(63w_{2x}^2 - 21cw_{2x} + c^2), \quad (118)$$

$$w_{3y} = 3w_{3x}^2, w_{3t_9} = 10206w_{3x}^5, \quad (119)$$

$$w_{4y} = \frac{3}{2}w_{4x}^2, w_{4t_9} = \frac{5103}{8}w_{4x}^5, \quad (120)$$

$$w_{5y} = \frac{3}{2}w_{5x}^2 + c_1, w_{5t_9} = \frac{5103}{8}w_{5x}^5 + \frac{1215}{2}c_1w_{5x}^3 + \frac{405}{2}c_1^2w_{5x}, \quad (121)$$

$$w_{6y} = \frac{3}{2}w_{6x}^2 + c_1, w_{6t_9} = \frac{5103}{8}w_{6x}^5 + \frac{1215}{2}c_1w_{6x}^3 + \frac{405}{2}c_1^2w_{6x}, \quad (122)$$

$$w_{7y} = 3w_{7x}^2, w_{7t_9} = 10206w_{7x}^5, \quad (123)$$

$$w_{8y} = -\frac{3}{2}w_{8x}^2, w_{8t_9} = \frac{5103}{8}w_{8x}^5, \quad (124)$$

$$w_{9y} = -\frac{3}{2}w_{9x}^2 + c_1, w_{9t_9} = \frac{5103}{8}w_{9x}^5 - \frac{1215}{2}c_1w_{9x}^3 + \frac{405}{2}c_1^2w_{9x}, \quad (125)$$

$$w_{10y} = -\frac{3}{2}w_{10x}^2 + c_1, w_{10t_9} = \frac{5103}{8}w_{10x}^5 - \frac{1215}{2}c_1w_{10x}^3 + \frac{405}{2}c_1^2w_{10x}, \quad (126)$$

$$w_{11y} = -\frac{3}{2}w_{11x}^2 + c_1, w_{11t_9} = \frac{5103}{8}w_{11x}^5 - \frac{1215}{2}c_1w_{11x}^3 + \frac{405}{2}c_1^2w_{11x}, \quad (127)$$

$$w_{12y} = \frac{3}{2}w_{12x}^2 + c_1, w_{12t_9} = \frac{5103}{8}w_{12x}^5 + \frac{1215}{2}c_1w_{12x}^3 + \frac{405}{2}c_1^2w_{12x} \quad (128)$$

of the ninth-order dPBKP equation (88) give rise to the linear superposition solutions

$$\begin{aligned} w_{13} &= w_1 + w_2, w_{14} = w_3 + \frac{1}{2}w_4, w_{15} = \frac{1}{2}(w_5 + w_6), \\ w_{16} &= w_7 - \frac{1}{2}w_8, w_{17} = -\frac{1}{2}(w_9 + w_{10}), w_{18} = -\frac{1}{2}w_{11} + \frac{1}{2}w_{12} \end{aligned} \quad (129)$$

of the same equation.

Proposition 12. For any solution w_i , $i = 1, 2, \dots, 13$, satisfying decomposition conditions

$$w_{1y} = 3w_{1x}^2 + cw_{1x}, w_{1t_{11}} = 11w_{1x}(c + 3w_{1x})(3402w_{1x}^4 + 2268cw_{1x}^3 + 504c^2w_{1x}^2 + 42c^3w_{1x} + c^4), \quad (130)$$

$$w_{2y} = 3w_{2x}^2 - cw_{2x}, w_{2t_{11}} = 11w_{2x}(3w_{2x} - c)(3402w_{2x}^4 - 2268cw_{2x}^3 + 504c^2w_{2x}^2 - 42c^3w_{2x} + c^4), \quad (131)$$

$$w_{3y} = 3w_{3x}^2, w_{3t_{11}} = 112266w_{3x}^6, \quad (132)$$

$$w_{4y} = \frac{3}{2}w_{4x}^2 + c_1, w_{4t_{11}} = \frac{6237}{16}w_{4x}^2(9w_{4x}^4 + 10c_1w_{4x}^2 + 4c_1^2), \quad (133)$$

$$w_{5y} = \frac{3}{2}w_{5x}^2 + c_1, w_{5t_{11}} = \frac{6237}{16}w_{5x}^2(9w_{5x}^4 + 10c_1w_{5x}^2 + 4c_1^2), \quad (134)$$

$$w_{6y} = -\frac{3}{2}w_{6x}^2 + c_1, w_{6t_{11}} = -\frac{6237}{16}w_{6x}^2(9w_{6x}^4 - 10c_1w_{6x}^2 + 4c_1^2), \quad (135)$$

$$w_{7y} = -\frac{3}{2}w_{7x}^2 + c_1, w_{7t_{11}} = -\frac{6237}{16}w_{7x}^2(9w_{7x}^4 - 10c_1w_{7x}^2 + 4c_1^2), \quad (136)$$

$$w_{8y} = \frac{3}{2}w_{8x}^2, w_{8t_{11}} = \frac{56113}{16}w_{8x}^6, \quad (137)$$

$$w_{9y} = -\frac{3}{2}w_{9x}^2, w_{9t_{11}} = -\frac{56113}{16}w_{9x}^6, \quad (138)$$

$$w_{10y} = 6c_+w_{10x}^2, w_{10t_{11}} = \frac{4032(503275411c_+ - 30935290)}{519921875}w_{10x}^6, c_{\pm} = \frac{-109 \pm \sqrt{44881}}{1605}, \quad (139)$$

$$w_{11y} = 6c_+w_{11x}^2, w_{11t_{11}} = \frac{4032(503275411c_+ - 30935290)}{519921875}w_{11x}^6, \quad (140)$$

$$w_{12y} = 6c_-w_{12x}^2, w_{12t_{11}} = \frac{4032(503275411c_- - 30935290)}{519921875}w_{12x}^6, \quad (141)$$

$$w_{13y} = 6c_-w_{13x}^2, w_{13t_{11}} = \frac{4032(503275411c_- - 30935290)}{519921875}w_{13x}^6, \quad (142)$$

there exist the linear superposition solutions

$$\begin{aligned}
 w_{14} &= w_1 + w_2, \quad w_{15} = \frac{1}{2}(w_4 + w_5), \quad w_{16} = -\frac{1}{2}(w_6 + w_7), \quad w_{17} = \frac{1}{2}(w_4 - w_6), \\
 w_{18} &= w_3 + \frac{1}{2}w_8, \quad w_{19} = w_3 - \frac{1}{2}w_9, \quad w_{20} = w_3 + 2c_+w_{10}, \quad w_{21} = w_3 + 2c_-w_{12}, \\
 w_{22} &= \frac{1}{2}w_8 + 2c_+w_{10}, \quad w_{23} = \frac{1}{2}w_8 + 2c_-w_{12}, \quad w_{24} = 2c_+w_{10} - \frac{1}{2}w_9, \quad w_{25} = 2c_-w_{12} - \frac{1}{2}w_9, \\
 w_{26} &= 2c_+w_{10} + 2c_+w_{11}, \quad w_{27} = 2c_+w_{10} + 2c_-w_{12}, \quad w_{28} = 2c_-w_{12} + 2c_-w_{13}
 \end{aligned} \tag{143}$$

of the eleventh-order dPBKP equation (89).

Conjecture 6. As the orders of the dPBKP equations increase, the possible types of linear superposition solutions increase.

Same as in the PBKP hierarchy cases, the special linear superposition solutions of the dPBKP equations are not decomposition solutions.

6 | CONCLUSIONS

We analyze the special decompositions and some linear superpositions of the BKP hierarchy and the dispersionless BKP hierarchy from their potential forms which are constructed in terms of the mastersymmetry approach. For the PBKP hierarchy, we conjecture that every order PBKP equation admits six different types of decomposition solutions. Particularly, we prove this conjecture for three lower-order PBKP equations. These decompositions provide simple and surprising relationships between several classic integrable systems and the PBKP hierarchy. The interrelations between different integrable systems enable us to discover unexpected connections of various mathematical and physical problems, ultimately solve them. Another interesting result about these special decomposition solutions is their suitable linear superpositions can also yield new solutions of the PBKP hierarchy, although the linear superposition theorem does not apply to nonlinear systems in general. This is equivalent to say that a possible linear combination of two special types of decomposition solutions holds again a solution to the same equation. The obtained linear superposition formulas allow us to construct many new exact solutions including $m+n$ solitons, n solitons with periodic cnoidal background waves, n soliton solutions combined with soliton-cnoidal wave interaction solutions, and combination solutions of two different modified Schwarzian KdV waves, etc. of the PBKP hierarchy. We mention that only two of the decomposition solutions can be linearly combined to construct three new possible solutions, although there are six possible decompositions for all equations in the PBKP hierarchy. The correctness of this has been verified through the fifth-, seventh-, and ninth-order PBKP equations.

Furthermore, we propose conjectures about the existence of decompositions and linear superposition solutions for each equation in the dPBKP hierarchy and, particularly, verify these conjectures for the fifth-, seventh-, ninth-, and eleventh-order dPBKP equations. Different from that which is common to all the members of the PBKP hierarchy containing the same possible linear superpositions, the possible ways of linear superposition increase rapidly as the order of the dPBKP equation increases.

We also emphasize that the significance of our findings about linear superposition properties is not restricted to the PBKP hierarchy or the dispersionless PBKP hierarchy. Our findings confirm the existence of some possible special linear superposition solutions in nonlinear systems and add the richness of exact solutions. Particularly, the existence of such linear superposition solutions in nonlinear systems provides us a totally new insight into the physical nature.

Naturally, generalizing the linear superposition principle, we will further consider whether possible linear superpositions of three or more special types of decomposition solutions exist in the PBKP hierarchy or dPBKP hierarchy. And is it possible to find linear superposition properties in other wonderful properties such as Hirota's bilinear representations, Schwarzian forms, etc. of the integrable systems? There is much remaining to be explored about how linear superposition principle can challenge and enrich our understanding of nonlinear systems. These, and other issues concerning linear superpositions, merit investigation.

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Conflict of interest

This work does not have any conflicts of interest.

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