

1           **Antarctic vortex dehydration in 2023 as a substantial**  
2           **removal pathway for Hunga Tonga-Hunga Ha'apai**  
3           **water vapour**

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18           **Key Points:**

- 19           • Antarctic dehydration is a major removal pathway of stratospheric H<sub>2</sub>O injected  
20           from Hunga Tonga-Hunga Ha'apai (HTHH) eruption  
21           • HTHH H<sub>2</sub>O caused small (up to 10 DU) additional chemical ozone depletion in  
22           2023 Antarctic spring  
23           • Model indicates e-folding timescale of 4 years for removal of HTHH H<sub>2</sub>O from strato-  
24           sphere

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## 25 Abstract

26 The January 2022 eruption of Hunga Tonga-Hunga Ha’apai (HTHH) injected a huge  
 27 amount ( $\sim 150$  Tg) of water vapour ( $\text{H}_2\text{O}$ ) into the stratosphere, along with small amount  
 28 of  $\text{SO}_2$ . An off-line 3-D chemical transport model (CTM) successfully reproduces the spread  
 29 of the injected  $\text{H}_2\text{O}$  through October 2023 as observed by the Microwave Limb Sounder  
 30 (MLS). Dehydration in the 2023 Antarctic polar vortex caused the first substantial ( $\sim 20$   
 31 Tg) removal of HTHH  $\text{H}_2\text{O}$  from the stratosphere. The CTM indicates that this pro-  
 32 cess will dominate removal of HTHH  $\text{H}_2\text{O}$  for the coming years, giving an overall e-folding  
 33 timescale of 4 years; around 25 Tg of the injected  $\text{H}_2\text{O}$  is predicted to still remain in the  
 34 stratosphere by 2030. Following relatively low Antarctic column ozone in midwinter 2023  
 35 due to transport effects, additional springtime depletion due to  $\text{H}_2\text{O}$ -related chemistry  
 36 was small and maximised at the vortex edge (10 DU in column).

## 37 Plain Language Summary

38 Around 150 Tg (150 million tons) of water vapour was injected into the stratosphere  
 39 during the eruption of Hunga Tonga-Hunga Ha’apai. Water vapour is a greenhouse gas  
 40 and this increase is expected to have a warming effect in the troposphere, as well caus-  
 41 ing perturbations in stratospheric chemistry and aerosols. We use an atmospheric model  
 42 to study the residence time of this excess water vapour and its impact on the recent Antarc-  
 43 tic ozone hole. The model performance is evaluated by comparison with satellite mea-  
 44 surements. Wintertime dehydration in the Antarctic stratosphere in 2023 is found to be  
 45 an important mechanism for removal of the volcanic water from the stratosphere. How-  
 46 ever, the overall removal rate is predicted to be slow; around 25 Tg (17%) is still present  
 47 in 2030. The direct impact of the excess water vapour on ozone via chemical processes  
 48 in the Antarctic ozone hole in 2023 is small.

## 49 1 Introduction

50 The eruption on 15<sup>th</sup> January 2022 of the submarine Hunga Tonga - Hunga Ha’apai  
 51 (HTHH) volcano ( $20.54^\circ\text{S}$ ,  $175.38^\circ\text{W}$ ) is recognized as the most explosive in the last 30  
 52 years, with emissions reaching up to  $\sim 55$  km (Carr et al., 2022; Taha et al., 2022). It  
 53 was unusual due to the huge amount of water vapour ( $\text{H}_2\text{O}$ ) injected very high into the  
 54 stratosphere, along with only small quantities of sulfur dioxide ( $\text{SO}_2$ ), thereby challeng-  
 55 ing many preconceptions about the atmospheric impacts of volcanic eruptions. This ex-  
 56 ceptional event is a global experiment allowing us to study, for the first time, a water-  
 57 rich volcanic eruption. Microwave Limb Sounder (MLS) satellite measurements indicate  
 58 that around 150 Tg of  $\text{H}_2\text{O}$  was injected, increasing the stratospheric burden by around  
 59 10% (Millán et al., 2022; Xu et al., 2022; Khaykin et al., 2022), while the  $\text{SO}_2$  injection  
 60 was only 0.5 Tg. This is expected to generate a very different climate forcing to other  
 61 satellite-observed  $\text{SO}_2$ -rich volcanic eruptions, possibly leading to a net warming of the  
 62 global surface temperature due to the dominant radiative effect of  $\text{H}_2\text{O}$  perturbations  
 63 (Sellitto et al., 2022; Jenkins et al., 2023).

64 The slow spreading of the injected  $\text{H}_2\text{O}$  throughout the stratosphere via the Brewer-  
 65 Dobson circulation (BDC) (Coy et al., 2022; Manney et al., 2023) is also expected to af-  
 66 fect stratospheric chemistry and dynamics. Rapid ozone depletion was observed in the  
 67 initial plume (Evan et al., 2023), along with the rapid formation of a dense aerosol layer  
 68 as a result of the water vapour injection (Asher et al., 2023; Zhu et al., 2022). In addi-  
 69 tion, evident processing of chlorine and depletion of nitrogen was observed in the south-  
 70 ern tropical stratosphere immediately after the HTHH eruption, which then spread pole-  
 71 ward over the following months (Santee et al., 2023). The aerosol layer was transported  
 72 polewards at lower altitudes than the  $\text{H}_2\text{O}$  enhancement. The excess  $\text{H}_2\text{O}$  caused a strong  
 73 cooling in the SH mid-latitude stratosphere shortly after the eruption (Schoeberl et al.,

2022; Vömel et al., 2022), which in turn strengthened the mid-latitude jet and slowed down the BDC (Coy et al., 2022). When the HTHH H<sub>2</sub>O reaches high latitudes with the descent of the BDC, it can affect gas-phase and heterogeneous processes related to polar ozone loss. Determining the timing and longevity of the excess H<sub>2</sub>O is thus critically important for assessing the impact on stratospheric ozone recovery and near-term climate change.

Water vapour can affect processes that drive stratospheric ozone in many ways. One important example is the formation of polar stratospheric clouds (PSCs), which initiate ozone-depleting heterogeneous chemistry. Sedimentation of ice PSCs irreversibly changes the H<sub>2</sub>O amount in the polar vortex, and affects the ozone-depletion processes via dehydration and denitrification (e.g., Fahey et al., 2001; Kelly et al., 1989; Feng et al., 2011; Tabazadeh et al., 2000). Dehydration in Antarctic winter has long been observed (e.g., Kelly et al., 1989; Vömel et al., 1995; Rosenlof et al., 1997; Tomikawa et al., 2015), but its representation by models can vary when applying PSC schemes with different complexity. If we use the HTHH water transport and its dehydration at polar regions as metrics to test a model’s stratospheric transport and PSC processes, we can then predict the longevity of the excess H<sub>2</sub>O by calculating its annual removal amount.

In this paper we use an off-line 3-D chemical transport model (CTM) to simulate the spatio-temporal evolution of the injected H<sub>2</sub>O with the results showing good agreement with MLS measurements in terms of plume spread and removal of HTHH H<sub>2</sub>O from the stratosphere. We estimate the longevity of the excess H<sub>2</sub>O and the amount that may remain in the stratosphere over the coming decade. We also diagnose the direct chemical impact of the increased H<sub>2</sub>O on stratospheric ozone through gas-phase and heterogeneous chemistry (e.g. PSCs and aerosols). The impacts are simulated with specified realistic post-eruption meteorology and hence do not account explicitly for dynamical feedbacks. Our CTM setup nonetheless provides useful constraints on changes seen in more complex coupled radiative-dynamical-chemical models (Wang et al., 2023).

## 2 Model and Observations

The TOMCAT/SLIMCAT CTM (Chipperfield, 1999, 2006) was run at a horizontal resolution of  $2.8^\circ \times 2.8^\circ$  and 32 levels from the surface to about 60 km forced with European Centre for Medium-Range Weather Forecasts (ECMWF) ERA-5 meteorology (Hersbach et al., 2020). The model uses a detailed gas-phase stratospheric chemistry scheme, and a simplified PSC scheme for the simulation of heterogeneous chemistry based on the assumption of thermodynamic equilibrium between PSC particles, including liquid aerosol, solid nitric acid trihydrate, and/or solid ice particles (Groß et al., 2018), and the gas phase (e.g., Feng et al., 2011, 2021). In this scheme, ice particles with assumed radius of 10  $\mu\text{m}$  sediment with a fall velocity of 1500 m/day. A control simulation (**Control**) without treatment of HTHH was integrated from 1980 to October 2023. Output from run **Control** for January 1st 2022 was used to initialise a run (**HT**) until October 31st 2023 with the injection of 150 Tg of H<sub>2</sub>O into the low-mid stratosphere at southern subtropical latitudes. We experimented with the timing of the model H<sub>2</sub>O injection between January 15<sup>th</sup> and April 1<sup>st</sup>. A later injection date, when the plume is already well spread longitudinally and latitudinally (i.e. April 1st,  $0^\circ$ - $360^\circ\text{E}$ ,  $2^\circ\text{S}$ - $28^\circ\text{S}$ ), overcomes inconsistencies between the initial plume dynamics and the coarse resolution CTM. The model used here employs a climatological distribution of H<sub>2</sub>O in the troposphere, so that any excess H<sub>2</sub>O transported to this region is removed from the model. The runs used background fields for sulfuric acid aerosols with no enhancement due to HTHH. The magnitude and impact of the HTHH SO<sub>2</sub> is uncertain (Wang et al., 2023) and in this study we focus on H<sub>2</sub>O alone.

To test the possible future evolution of the HTHH H<sub>2</sub>O three further model runs were performed. These were integrated from January 1<sup>st</sup> 2023 until 2030 using repeat-

ing ERA-5 meteorology for 2022. Run **Con\_2022** was essentially an extension of run **Control**; run **HT\_2022** was an extension of **HT**; run **HT\_2022ns** was the same as run **HT\_2022** but had sedimentation of PSC particles turned off. The experiments are summarized in Supplementary Table S1.

The modelled results are compared to satellite measurements of H<sub>2</sub>O from MLS (Waters et al., 2006) and total column ozone from the Ozone Monitoring Instrument (OMI) (Levelt et al., 2006) on the NASA Aura satellite, and the Infrared Atmospheric Sounding Interferometer (IASI) (Siddans et al., 2018) on MetOp-B satellite.

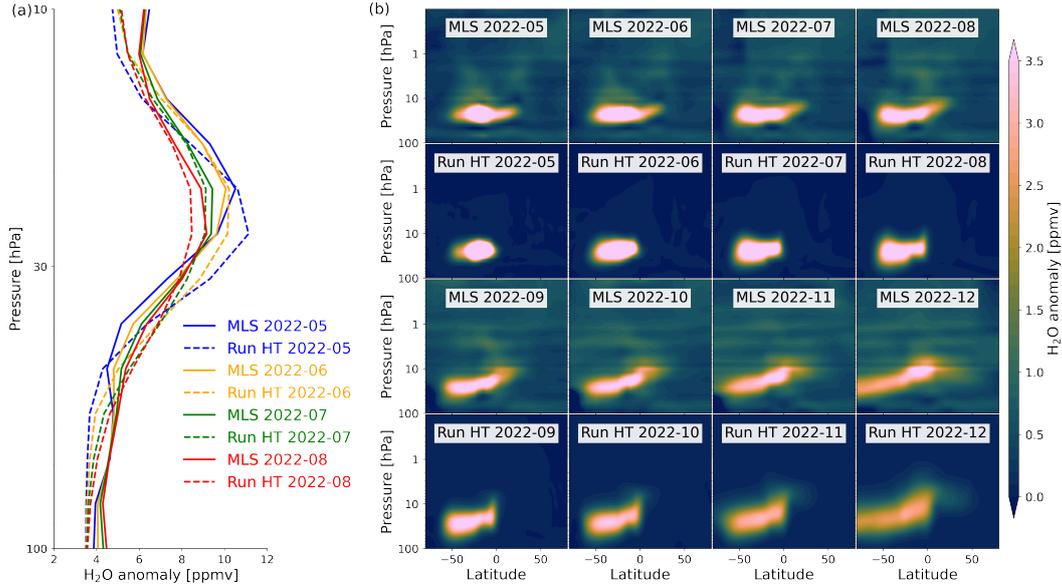
MLS H<sub>2</sub>O anomalies are calculated as deviations from the climatology of 2005-2021. The stratospheric total mass of H<sub>2</sub>O is estimated as the global (80°S-80°N) sum of the stratospheric column over each 5° latitude band from MLS volume mixing ratio measurements on pressure levels. For this study, we use MLS version 4 (v4) and version 5 (v5) products for the H<sub>2</sub>O mass, but use v4 for analysis of the vertical structure of the plume in view of the poor fits of v5 to H<sub>2</sub>O signals in regions with extremely enhanced humidity (Millán et al., 2022).

### 3 Evaluation of the model post-eruption stratospheric transport

We first assess the model performance for the H<sub>2</sub>O transport after eruption. Figure 1a shows the H<sub>2</sub>O anomaly profiles after the HTHH injection in model run **HT** (dashed line), compared with MLS measurements (solid line). The two are in good agreement, both showing the positive water vapour anomaly of 8–11 ppmv peaking between 30 hPa and 10 hPa from April to September. While the injected total mass in the model is consistent with MLS, the simulation has slightly larger peak anomalies and smaller horizontal extent after injection. The simulated plume spread is in very good agreement with the observations (Figure 1b and Supplementary Movie S1), in particular regarding the characteristics and behaviour of the excess H<sub>2</sub>O at the mixing barriers in the stratosphere, including the polar vortex edge, the extratropical tropopause, and the tropical pipe. Around 4-6 months after the eruption, the excess H<sub>2</sub>O moves into the Southern Hemisphere (SH) mid-latitudes within the shallow branch of the BDC, i.e. via the tropical pipe (Plumb, 1996). However, it does not intrude into the 2022 Antarctic polar vortex due to the polar jet. Only after the breakdown of the Antarctic polar vortex in November 2022 did the H<sub>2</sub>O reach the pole (see also Manney et al. (2023)). The circulation associated with the easterly phase of the Quasi-Biennial Oscillation (QBO) confine the excess H<sub>2</sub>O to the SH until the transition to westerlies at the end of 2022. When the SH moves into austral winter, H<sub>2</sub>O enters the deep branch of the BDC, ascending from the tropics and descending into the high latitudes in the SH. The model reproduces well the timing of the HTHH-injected H<sub>2</sub>O penetrating the polar vortex, and the altitudes of the H<sub>2</sub>O plume. This indicates the model has a good representation of both the poleward horizontal H<sub>2</sub>O transport by the shallow branch of the BDC and the ascent of the water-enriched air to high levels by the deep branch of the BDC.

### 4 Dehydration in 2023 Antarctic winter

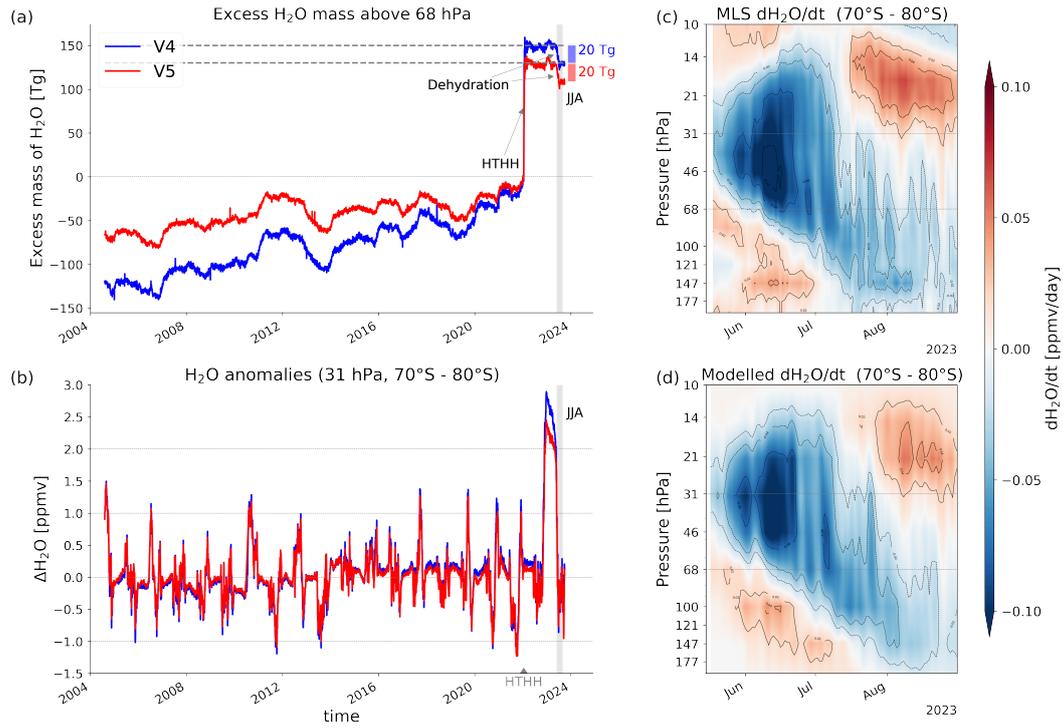
Figure 2a shows the time evolution series of total excess stratospheric H<sub>2</sub>O mass above 68 hPa observed by the MLS. After the enhancement of the stratospheric H<sub>2</sub>O mass by ~150 Tg after HTHH injection in January 2022, the amount of excess H<sub>2</sub>O remained steady until a sudden drop of ~20 Tg from June to July 2023. This strong dehydration is also seen in the time series of the Antarctic H<sub>2</sub>O mixing ratio at 31 hPa (Figure 2b). First, the excess H<sub>2</sub>O rose to its highest and unprecedented level at the end of May 2023 in this 20-year record, when the HTHH injected H<sub>2</sub>O entered the Antarctic stratosphere via the deep branch of the BDC. Then, a striking drop in H<sub>2</sub>O occurred within just three months from June to August 2023. The H<sub>2</sub>O anomaly fell from 2.5 ppmv to close to zero in August 2023. The amplitude of this stratospheric dehydration is also unprecedented



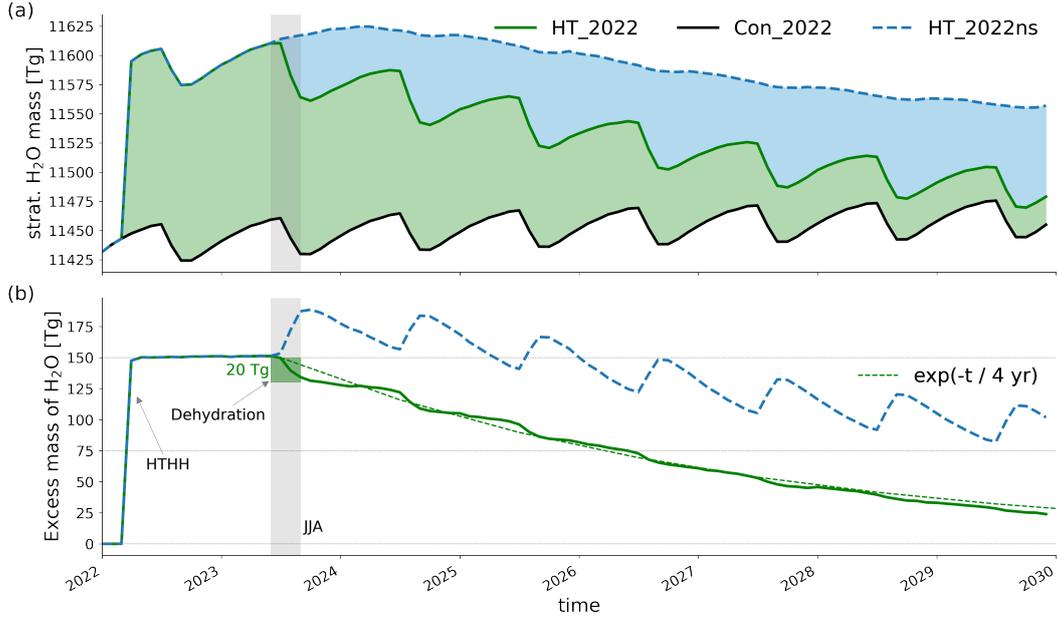
**Figure 1.** Water vapour ( $\text{H}_2\text{O}$ ) profile and evolutions after the HTHH eruption. (a) Zonal mean  $\text{H}_2\text{O}$  anomaly profiles (ppmv) between  $40^\circ\text{S}$ - $20^\circ\text{S}$  from May to September 2022 from MLS v4 observations (solid lines) and run **HT** (dashed lines). (b) Zonal mean latitude-pressure cross sections of  $\text{H}_2\text{O}$  anomalies observed by MLS v4 and simulated by model run **HT** from May to December 2022.

175 in the observational record since 2004. The fact that the  $\text{H}_2\text{O}$  anomaly returns to essentially  
 176 zero in August, like the previous years, shows that the vortex  $\text{H}_2\text{O}$  is controlled  
 177 by the local polar temperatures (so that water vapour remaining in the gas-phase is deter-  
 178 mined by its saturation vapour pressure) and is not affected by additional volcanic  
 179 injection; i.e. ice PSC processes effectively cancel out the in-vortex impact of the addi-  
 180 tional HTHH water vapour by the end of Austral winter. Note that the model has a good  
 181 representation of  $\text{H}_2\text{O}$  inside the Antarctic polar vortex core in winter/spring 2023 (Fig-  
 182 ure S1).

183 Figures 2c and 2d show the daily tendencies of the  $\text{H}_2\text{O}$  mixing ratio inside the Antarc-  
 184 tic vortex observed by MLS v4 and simulated by the model. The dramatic dehydrated  
 185 areas above 68 hPa are clearly seen in the lower stratosphere in June 2023 with the largest  
 186 rate of  $\text{H}_2\text{O}$  decrease of  $-0.19$  ppmv/day. This led to a fast transition of the lower strato-  
 187 sphere from anomalously wet conditions to dry in the Antarctic vortex core. The region  
 188 of strong dehydration is in good agreement with the vertical domain of PSCs that are  
 189 usually observed between 15 and 25 km. The strong dehydration is accompanied by an  
 190 increase in  $\text{H}_2\text{O}$  mixing ratio below, indicating enriched  $\text{H}_2\text{O}$  below the dehydrated re-  
 191 gion. This dehydration and rehydration below are visible until August. Importantly, we  
 192 see a clear descent of the rehydration over time, causing water originating from higher  
 193 altitudes to accumulate initially in the polar lowermost stratosphere. The dehydrated  
 194 and rehydrated air carries on descending throughout winter forced by the BDC with ul-  
 195 timately  $\text{H}_2\text{O}$  being irreversibly removed from the stratosphere. Based on the good agree-  
 196 ment between the model simulation and MLS data, the observed dehydration and rehy-  
 197 dration below can only be linked to the sedimentation of the ice PSC particles, evap-  
 198 oration at lower levels and descent of this rehydrated air with the BDC. The important  
 199 consequence is that a substantial amount of HTHH  $\text{H}_2\text{O}$  is removed from the stratosphere  
 200 during the austral winter.



**Figure 2.** MLS observed and TOMCAT simulated dehydration of the Antarctic polar vortex. (a) Time series of observed excess H<sub>2</sub>O mass (Tg) above 68 hPa from MLS v4 and v5. (b) As panel (a) but for MLS v4 mean SH polar cap (80°S-70°S) H<sub>2</sub>O mixing ratio (ppmv) at 31 hPa. The grey bar marks the months of June, July and August 2023. (c) MLS v4 observed daily tendencies of mean SH polar cap H<sub>2</sub>O mixing ratio (ppmv/day) with a 30-day smoothing from June to August in 2023. (d) Same as (c) but for TOMCAT simulated (run **HT**) daily tendencies.

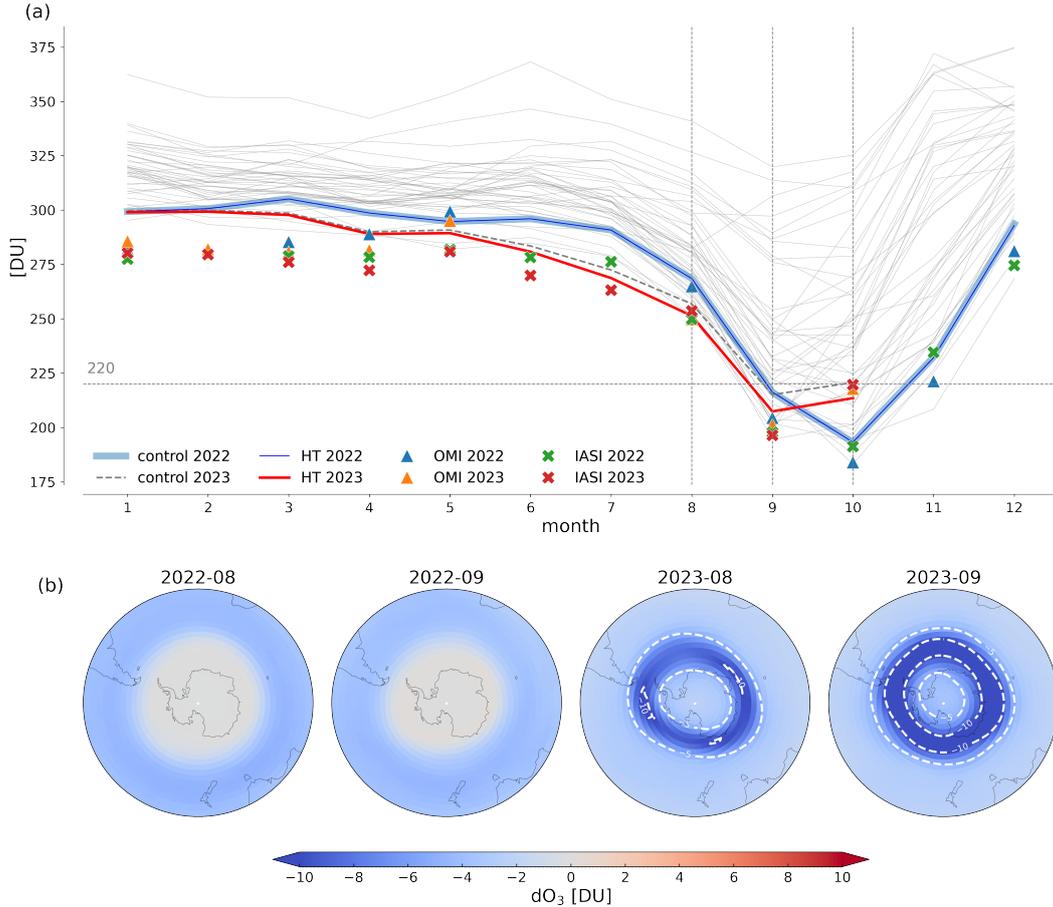


**Figure 3.** TOMCAT projections of the decay rate of the excess H<sub>2</sub>O. (a) Time series of simulated total abundance of H<sub>2</sub>O above 100 hPa (Tg) for runs **Con\_2022**, **HT\_2022**, and **HT\_2022ns**. Grey shading indicates the Austral winter (JJA) 2023. (b) Differences (Tg) in panel (a) with respect to control run **Con\_2022** for runs **HT\_2022** and **HT\_2022ns**. The dotted curve illustrates exponential decay with timescale of 4 years starting in July 2023. A horizontal dotted line indicates 75 Tg, 50% of the initial injection.

## 201 5 Long-term decay of HTHH water vapour

202 To clarify the mechanism for the removal of the HTHH injected H<sub>2</sub>O, we compare  
 203 simulations **HT\_2022**, **HT\_2022ns** and **Con\_2022**. Figure 3a shows the total burden  
 204 of gas phase H<sub>2</sub>O above 100 hPa projected by TOMCAT simulation through 2030 for  
 205 these runs. The annual cycle for run **Con\_2022**, with a decrease in austral winter and  
 206 increase afterward, is related to the annual cycle of its sinks and sources, the dehydra-  
 207 tion inside the polar vortex in June to August, and the tape recorder signal with enhanced  
 208 H<sub>2</sub>O from the tropical troposphere to the stratosphere that has maximum enhancement  
 209 around September. A similar annual cycle exists in run **HT\_2022**. The differences be-  
 210 tween the two runs is because the dehydration is stronger for run **HT\_2022**, causing a  
 211 decreasing year-to-year difference between them. In contrast, the run **HT\_2022ns** does  
 212 not show this strong dehydration and annual cycle, with a slower decay throughout the  
 213 following years. Here, the H<sub>2</sub>O decay is caused by the stratosphere-to-troposphere trans-  
 214 port where stratospheric air with HTHH-injected H<sub>2</sub>O slowly descends into the tropo-  
 215 sphere at high latitudes via the deep branch of the BDC and is replaced by air from the  
 216 troposphere with H<sub>2</sub>O values determined by tropopause temperatures.

217 Figure 3b shows the differences in the total burden of H<sub>2</sub>O mass between the HTHH  
 218 perturbed runs and the control run. The excess H<sub>2</sub>O, once injected into the stratosphere,  
 219 is removed in an almost step-like fashion during austral winter starting in 2023. The amount  
 220 of H<sub>2</sub>O removal, around 20 Tg in 2023, is in very good agreement with the MLS observed  
 221 value, supporting the TOMCAT representation of the HTHH H<sub>2</sub>O transport and removal.  
 222 The modelled e-folding lifetime for the removal of the excess H<sub>2</sub>O is around 4 years (half-  
 223 life of  $\sim 2.8$  yrs) from the point at which removal starts (over a year after the eruption),



**Figure 4.** Stratospheric ozone changes after HTHH. (a) Time series of the mean total column ozone (DU) at SH high-latitudes (65°S-90°S), comparing run **Control** in 2022 (bold blue), 2023 (dashed grey) and years 1980 to 2021 (grey) with run **HT** in 2022 (thin navy) and 2023 (solid red), and with OMI (triangle), and IASI (cross) satellite observations in 2022 and 2023. (b) Modelled monthly mean difference in column ozone (DU) between runs **HT** and **Control** for August and September in 2022 and 2023.

224 so that  $\sim 25$  Tg remains in the stratosphere at 2030. The longevity of the HTHH injected  
 225 H<sub>2</sub>O thus exceeds 7 years. Comparing the simulated H<sub>2</sub>O mass in the stratosphere be-  
 226 tween runs **HT\_2022** and **HT\_2022ns**, it can be seen that PSC sedimentation plays  
 227 a key role in the removal of the HTHH H<sub>2</sub>O. Without PSC sedimentation, the decline  
 228 of H<sub>2</sub>O mass is much slower so that only around 50 Tg H<sub>2</sub>O is removed by 2030, account-  
 229 ing for  $\sim 38\%$  of the total removal in run **HT\_2022** by that time. The dehydration due  
 230 to the sedimentation of ice PSC particles thus accounts for more than 60% of the modelled  
 231 total removal over this period, serving as a main removal pathway of the HTHH  
 232 H<sub>2</sub>O. It is worth noting that PSC sedimentation is also important for the removal of back-  
 233 ground (non-volcanic) H<sub>2</sub>O. Without PSC sedimentation (run **HT\_2022ns**) the strato-  
 234 spheric total H<sub>2</sub>O burden would increase slightly during austral winter under the effect  
 235 of the positive anomalies in H<sub>2</sub>O at the tropical tropopause (Gilford & Solomon, 2017).  
 236 In the model PSC sedimentation removes  $\sim 30$  Tg of background H<sub>2</sub>O each austral win-  
 237 ter.

## 6 Ozone depletion due to the increased H<sub>2</sub>O

We now quantify the impact of the additional H<sub>2</sub>O on Antarctic ozone through October 2023 and compare with previous years (Figure 4). In 2022, there is negligible modelled chemical impact on the total column ozone at high SH latitudes (65°S-90°S), consistent with the failure of HTHH H<sub>2</sub>O to penetrate into the Antarctic polar vortex. In contrast, in 2023 it reached the pole before the vortex formation and thus caused a direct impact. Interestingly, in mid-winter 2023, before the onset of substantial chemical ozone depletion, modelled column ozone is smaller than other modelled years since 1980. This apparent earlier start of the ozone hole was observed by IASI. However, the small difference (10 DU) between model runs **HT** and **Control** shows that this early onset was dynamically driven rather than a chemical impact of the enhanced H<sub>2</sub>O. The HTHH injection has been found to be linked to a stable and colder-than-normal vortex, with a slowdown of the BDC (Wang et al., 2023); our CTM may be capturing this effect via the specified ERA-5 meteorology but here we cannot quantify it. An unusual transport of ozone-poor air from the upper stratosphere to the lower stratosphere, indicated by the increase of age-of-air in the stratosphere (Figure S2), can lead to an anomalous decrease of the ozone at lower altitudes.

Meanwhile, an earlier formation of the PSCs can lead to more extensive heterogeneous processing and ozone depletion. However, the strong dehydration due to the sedimentation of ice PSC particles (see above) limits the impact of the additional H<sub>2</sub>O in the core of the polar vortex. From June to September, an additional depletion of ozone up to 10 DU (around 4% of the background) due to the injected H<sub>2</sub>O occurs at the vortex edge, a region of available sunlight and where ozone loss is not saturated. While the modelled mean column ozone in 2023 is outside the range of previous years in June and July, by September and October it is no longer an outlier. Hence, possible early indications of record low springtime ozone did not occur.

Model runs **Con\_2022** and **HT\_2022** can be used to estimate the longer term impact of the HTHH H<sub>2</sub>O over the next 5 years (Figure S3). The largest column depletion occurs at the edge of the Antarctic vortex in 2023 at 10 DU (see also Figure 4). The impact on the Antarctic ozone hole then decreases in subsequent years as the HTHH H<sub>2</sub>O decays. Other large impacts on column ozone occur in the SH midlatitudes in 2022 and, to a lesser extent, the Arctic winter/spring. Interestingly, the impact on the Arctic maximises in winter 2024/25 due to the spread of the HTHH H<sub>2</sub>O, but note that these runs use repeating 2022 meteorology. In reality Arctic ozone loss is dominated by meteorological variability.

## 7 Summary and Discussion

Near-term projections of the HTHH climate impacts depend strongly on the estimation of the transport and longevity of the injected water vapour. Here we show that the Antarctic ice PSC sedimentation is likely a major removal pathway from the stratosphere for the HTHH-injected H<sub>2</sub>O. This PSC sedimentation is partly responsible for a small estimated impact on chemical loss in the 2023 Antarctic vortex. Nevertheless, there are many other potential chemical, microphysical and radiative impacts of H<sub>2</sub>O (and initial SO<sub>2</sub>) in the stratosphere. The projected long residence time of the HTHH H<sub>2</sub>O means that we can expect it to influence the atmosphere for many years. Many more modelling and observational studies are needed to quantify these impacts further.

## 8 Open Research

The v4 MLS water vapour data is available at [https://acdisc.gesdisc.eosdis.nasa.gov/data/Aura\\_MLS\\_Level12/ML2H20.004/](https://acdisc.gesdisc.eosdis.nasa.gov/data/Aura_MLS_Level12/ML2H20.004/); the v5 data is available at [https://acdisc.gesdisc.eosdis.nasa.gov/data/Aura\\_MLS\\_Level12/ML2H20.005/](https://acdisc.gesdisc.eosdis.nasa.gov/data/Aura_MLS_Level12/ML2H20.005/). IASI ozone

287 data can be downloaded from the Seris portal <http://iasi.aeris-data.fr/03/>. OMI  
 288 total column ozone product is available at [https://disc.gsfc.nasa.gov/datasets/  
 289 OMD0A03e\\_003/summary](https://disc.gsfc.nasa.gov/datasets/OMDOA03e_003/summary). TOMCAT model data is available at [https://doi.org/10.5281/  
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