

Verification of He's frequency formulation by Duffing-Harmonic Oscillator

Z.L. Tao¹

¹ Nanjing Univ Informat Sci & Technol

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Abstract

He's frequency formulation suggests the fast identification of the amplitude-frequency relationship of a nonlinear oscillator. This paper applies the formulation to the Duffing-harmonic oscillator with great success. A modification of He's frequency formulation and the variational iteration method are also used for comparison.

A modification of He's frequency-amplitude formulation

He's original frequency formulation was derived from an ancient Chinese mathematical algorithm [10, 13], recently there was a hot discussion on He's original frequency formulation [14-22]. Here we check Ren-Hu's modification [14].

Rewrite the problem (1) in this form

$$\ddot{x}(1+x^2)+x^3=0, x(0)=A, \dot{x}(0)=0. \quad (10)$$

In view of He's frequency-amplitude formulation [7-9], we use a trial solution

$$x_1 = A \cos \omega_1 t. \quad (11)$$

Taking equation (11) into the first equation of problem (10) results in the residual

$$\begin{aligned} R_1 &= -\omega_1^2 A \cos \omega_1 t (1 + A^2 \cos^2 \omega_1 t) + A^3 \omega_1 t \\ &= \left[\frac{3A^3}{4} (1 - \omega_1^2) - \omega_1^2 A \right] \cos \omega_1 t + \frac{A^3}{4} (1 - \omega_1^2) \cos 3\omega_1 t. \end{aligned} \quad (12)$$

Introducing defined in [14]

$$\tilde{R}_1 = \frac{4}{T} \int_0^{\frac{T}{4}} R_1(t) dt, \quad (13)$$

with T the period of the oscillator.

Submitting equation (12) into equation (13), we obtain

$$\begin{aligned} \tilde{R}_1 &= \frac{4}{T} \int_0^{\frac{T}{4}} \left\{ \left[\frac{3A^3}{4} (1 - \omega_1^2) - \omega_1^2 A \right] \cos \omega_1 t + \frac{A^3}{4} (1 - \omega_1^2) \cos 3\omega_1 t \right\} dt \\ &= \frac{2}{\pi} \left[\frac{3A^3}{4} (1 - \omega_1^2) - \omega_1^2 A \right] - \frac{(1 - \omega_1^2) A^3}{6\pi} \end{aligned} \quad (14)$$

Let $x_2 = A \cos \omega_2 t$ with $\omega_1 \neq \omega_2$, by the similar operation as the above, we have

$$\tilde{R}_2 = \frac{2}{\pi} \left[\frac{3A^3}{4} (1 - \omega_2^2) - \omega_2^2 A \right] - \frac{(1 - \omega_2^2) A^3}{6\pi}. \quad (15)$$

Using He's frequency-amplitude formulation [10, 13], we have

$$\omega^2 = \frac{\tilde{R}_2\omega_1^2 - \tilde{R}_1\omega_2^2}{\tilde{R}_2 - \tilde{R}_1} = \frac{\left\{ \frac{2}{\pi} \left[\frac{3A^3}{4} (1 - \omega_2^2) - \omega_2^2 A \right] - \frac{(1 - \omega_2^2)A^3}{6\pi} \right\} \omega - \left\{ \frac{2}{\pi} \left[\frac{3A^3}{4} (1 - \omega_1^2) - \omega_1^2 A \right] - \frac{(1 - \omega_1^2)A^3}{6\pi} \right\} \omega}{\frac{2}{\pi} \left[\frac{3A^3}{4} (1 - \omega_2^2) - \omega_2^2 A \right] - \frac{(1 - \omega_2^2)A^3}{6\pi} - \left\{ \frac{2}{\pi} \left[\frac{3A^3}{4} (1 - \omega_1^2) - \omega_1^2 A \right] - \frac{(1 - \omega_1^2)A^3}{6\pi} \right\}} = \frac{18A^2}{(9\pi - 2)A^2 + 24} \quad (16)$$

When , Eq. (16) becomes

$$\omega^2 = \frac{18}{(9\pi - 2)} \quad (17)$$

while the exact one is $\omega = 1$. When $A \ll 1$, we have

$$\omega^2 = \frac{18A^2}{24} = \frac{3A^2}{4} \quad (18)$$

This agrees with that by the homotopy perturbation method [10-12]. So Ren-Hu's modification is valid for $A \ll 1$.

Variational iteration method

Reconsider the problem (1), according to the variational iteration method with Laplace transform; we have the iteration equality [23]

$$\begin{aligned} L[x_{n+1}] &= L[x_n] - L \left[\int_0^t \frac{1}{\omega} \sin \omega(t - \xi) \delta \xi \right] \\ &= L[x_n] - \frac{1}{\omega} L[\sin \omega \tau] \bullet L[\ddot{x}_n(t) + x_n^3(t) + x_n^2(t) \ddot{x}_n(t)] \end{aligned} \quad (19)$$

where L is the Laplace transform operator.

Assume the initial approximation be

$$x_0 = A \cos \omega \tau \quad (20)$$

Then we have

$$\begin{aligned} L[x_1(t)] &= L[A \cos \omega \tau] - \frac{1}{\omega} L[\sin \omega \tau] \bullet L[-A\omega^2 \cos \omega t + A^3 \cos^3 \omega t - A^3 \omega^2 \cos^3 \omega \tau] \\ &= L[A \cos \omega \tau] - \frac{1}{\omega} L[\sin \omega \tau] L \left[\left(\frac{3(A^3 - A^3 \omega^2)}{4} - A\omega^2 \right) \cos \omega t + \frac{A^3 - A^3 \omega^2}{4} \cos 3\omega \tau \right] \end{aligned} \quad (21)$$

Imposing the inverse Laplace transform on equality (21), there holds the first-order approximant

$$x_1(t) = A \cos \omega \tau - \frac{t}{4\omega} \left(3(A^3 - A^3 \omega^2) - 4A\omega^2 \right) \sin \omega \tau + \frac{A^3 \omega^2 - A^3}{32\omega^2} (\cos \omega \tau - \cos 3\omega \tau) \quad (22)$$

No secular term in the equality (22) requires that

$$3(A^3 - A^3 \omega^2) - 4A\omega^2 = 0 \quad (23)$$

which yields this

$$\omega^2 = \frac{3A^2}{4 + 3A^2} \quad (24)$$

Eq. (24) satisfies two scale extremals.

Discussion and Conclusions

Eq. (1) can be solved effectively by the homotopy perturbation method [24-28] and the variational approach [2, 29]. By the semi-inverse method [30, 31], we can establish a variational formulation for Eq. (1), which is

$$J(u) = \int_0^{T/4} \left\{ \frac{1}{2} \dot{x}^2 - \frac{1}{2} x^2 + \frac{1}{2} \ln(1 + x^2) \right\} dt \quad (25)$$

By comparison also with those in open literatures, we conclude that He's frequency is the simplest while its accuracy is also extremely high. Considering the simplest calculation, He's frequency formulation greatly promote the development of the nonlinear science, especially the nonlinear vibration.

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